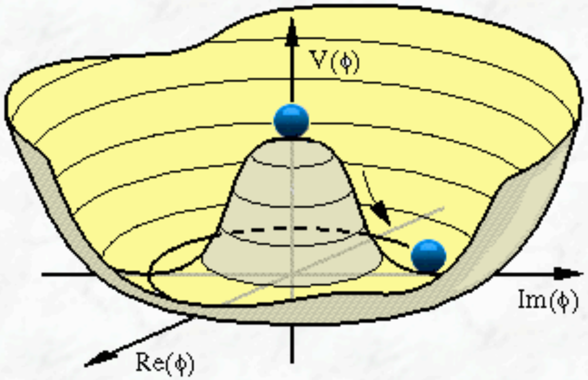
7. Physics of the Higgs Boson



7.1 The Higgs boson in the Standard Model

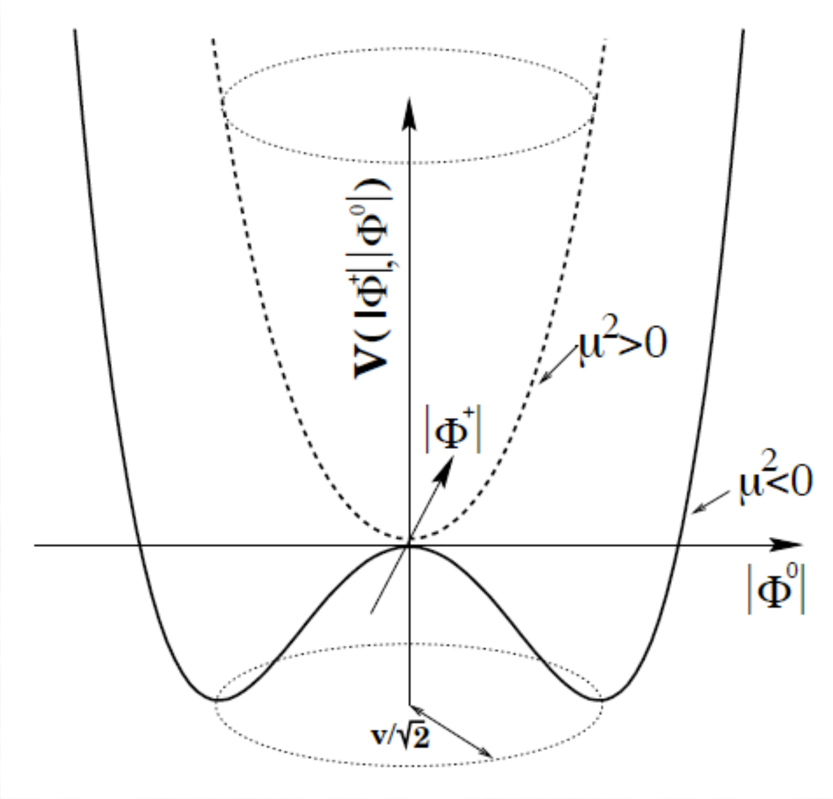
7.2 Properties of the Higgs boson

7.3 Higgs boson production at hadron colliders

7.4 The discovery of a Higgs boson at the LHC

7.5 What are its properties? Is it the Higgs boson of the Standard Model?

7.1 The Higgs boson in the Standard Model



The structure of the Standard Model

Fundamental principle: **Local gauge invariance**
Prototype: **Quantum Electrodynamics (QED)**

Free Dirac equation: $$i\gamma^{\mu}\partial_{\mu}\psi - m\psi = 0$$

Lagrangian formalism: $$L = i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi$$

Local gauge transformation: $$\psi(x) \rightarrow e^{i\alpha(x)}\psi(x)$$

(derivative: $$\partial_{\mu}\psi \rightarrow e^{i\alpha(x)}\partial_{\mu}\psi + ie^{i\alpha(x)}\psi\partial_{\mu}\alpha \; ,$$
$\delta_{\mu}\alpha$ term breaks the invariance of L)

Invariance of L under local gauge transformations can be accomplished by introducing a gauge field A_{μ}, which transforms as:

$$A_{\mu} \rightarrow A_{\mu} + \frac{1}{e}\partial_{\mu}\alpha$$ where $e = g_e/4\pi$ = coupling strength

Can be formally achieved by the construction of a “modified” derivative

$$\partial_{\mu} \rightarrow D_{\mu} = \partial_{\mu} - ieA_{\mu}$$ (covariant derivative)

→ Lagrangian of QED:

$$L = i\overline{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\overline{\psi}\psi + e\overline{\psi}\gamma^{\mu}A_{\mu}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

interaction term

where $F_{\mu\nu}$ is the usual field strength tensor: $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$

Note:

(i) Imposing local gauge invariance leads to the interacting field theory of QED

(ii) A mass term ($\frac{1}{2}m^2 A_{\mu}A^{\mu}$) for the gauge field A_{μ} would violate gauge invariance

Similar for the Standard Model interactions:

Quantum Chromodynamics (QCD):

$$D_\mu = \partial_\mu + igT_a G_\mu^a$$

$$G_\mu^a \rightarrow G_\mu^a - \frac{1}{g}\partial_\mu \alpha_a - f_{abc}\alpha_b G_\mu^c$$

SU(3) transformations, 8 gauge fields,
8 massless gluons, Gluon self-coupling

- T_a (a = 1,…,8) generators of the SU(3) group
  (independent traceless 3x3 matrices)
- G_μ gluon fields
- g = coupling constant

Electroweak Interaction (Glashow, Salam, Weinberg):

$SU(2)_L \times U(1)_Y$ transformations,
4 gauge fields, (W_μ^1, W_μ^2, W_μ^3, B_μ)

Physical states:

$$W_\mu^\pm = \frac{1}{\sqrt{2}}\left(W_\mu^1 \mp iW_\mu^2\right)$$

$$Z_\mu = -\sin\theta_W B_\mu + \cos\theta_W W_\mu^3$$

$$A_\mu = \cos\theta_W B_\mu + \sin\theta_W W_\mu^3$$

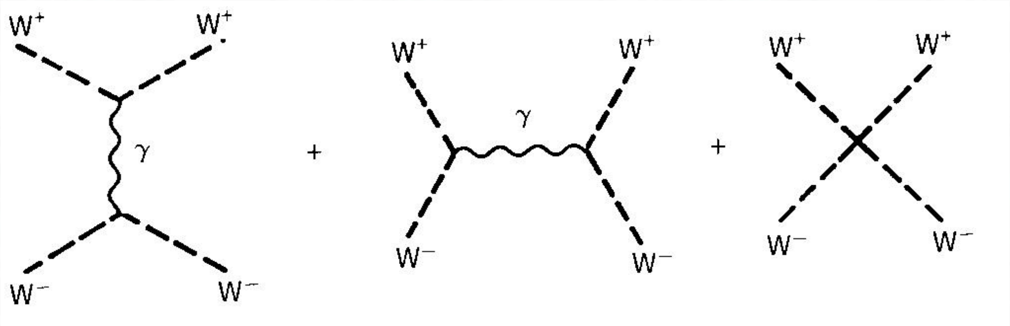
Problems at that stage:

- **Masses of the vector bosons W and Z:**

Experimental results: $m_W = 80.385 \pm 0.015$ GeV / c^2
$m_Z = 91.1875 \pm 0.0021$ GeV / c^2

A local gauge invariant theory requires massless gauge fields

- **Divergences in the theory (scattering of W bosons)**



$$-iM(W^+W^- \rightarrow W^+W^-) \sim \frac{s}{M_W^2} \quad \text{for} \quad s \rightarrow \infty$$

Solution to **both** problems:

- create mass via spontaneous breaking of electroweak symmetry

- introduce a scalar particle that regulates the WW scattering amplitude

➔ Higgs Mechanism

Der Higgs-Mechanismus in einer Toy-Theorie (I)

Mathematische Formulierung des Standard Model beruht auf zwei Grundprinzipien:

- **Lokale Eichinvarianz**: Lagrangedichte ist invariant unter bestimmten ortsabhängigen Eichtransformationen, z.B.: $\boldsymbol{\psi \rightarrow e^{i\theta(x)}\psi}$
- **Spontane Symmetriebrechung**

Klein-Gordon Gleichung für ein skalares (Spin 0) Feld:

$$\mathcal{L} = \frac{1}{2}\left(\partial_\mu \phi\right)\left(\partial^\mu \phi\right) - \frac{1}{2} m^2 \phi^2$$

Was ist die Masse von φ hier?

$$\mathcal{L} = \frac{1}{2}\left(\partial_\mu \phi\right)\left(\partial^\mu \phi\right) + \frac{1}{2}\mu^2 \phi^2 - \frac{1}{4}\lambda^2 \phi^4$$

- *Falsches Vorzeichen!*
- *φ = 0 ist nicht der Grundzustand!*

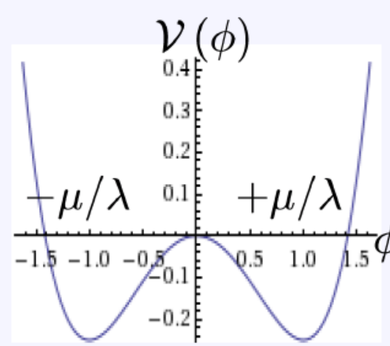
$$\mathcal{L} = \mathcal{T} - \mathcal{V}$$

Also ist:

$$\mathcal{V} = -\frac{1}{2}\mu^2 \phi^2 + \frac{1}{4}\lambda^2 \phi^4$$

Mit dem Minimum bei:

$$\phi = \pm\mu/\lambda$$



Der Higgs-Mechanismus in einer Toy-Theorie (II)

Definiere eine neue Feldvariable um ein Minimum herum:

$$\eta = \phi \pm \frac{\mu}{\lambda}$$

Damit wird der Lagrangian zu

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \eta)(\partial^\mu \eta) - \mu^2\eta^2 \pm \mu\lambda\eta^3 - \frac{1}{4}\lambda^2\eta^4 + \frac{1}{4}\left(\mu^2/\lambda\right)^2$$

und die Masse von η ist

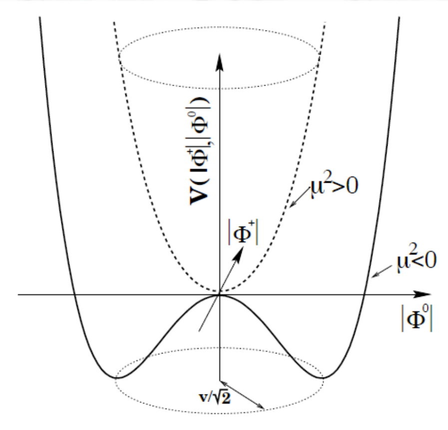
$$m = \sqrt{2}\mu$$

- Die ursprüngliche **Symmetrie wurde gebrochen**
- Genauer: Die **Symmetrie wurde versteckt**

- Brout-Englert-Higgs-Mechanismus: Ersetze ϕ durch komplexes skalares Doublet ϕ

The Higgs mechanism

Spontaneous breaking of the SU(2) x U(1) gauge symmetry



- Scalar fields are introduced $\phi = \frac{1}{\sqrt{2}}\begin{pmatrix}\phi_1 + i\phi_2 \\ \phi_3 + i\phi_4\end{pmatrix} = \begin{pmatrix}\phi^+ \\ \phi^0\end{pmatrix}$

  Potential : $V(\phi) = \mu^2(\phi^*\phi) + \lambda(\phi^*\phi)^2$

- Lagrangian for the scalar fields:
  g, g ' = SU(2), U(1) gauge couplings

$$L_2 = \left|\left(i\partial_\mu - g\mathbf{T}\cdot\mathbf{W}_\mu - g'\frac{Y}{2}B_\mu\right)\phi\right|^2 - V(\phi)$$

- For $\mu^2 < 0$, $\lambda > 0$, minimum of potential:

$$\phi_1^2 + \phi_2^2 + \phi_3^2 + \phi_4^2 = v^2 \qquad v^2 = -\mu^2/\lambda$$

- Perturbation theory around ground state:

$$\phi_0(x) = \frac{1}{\sqrt{2}}\begin{pmatrix}0 \\ v + h(x)\end{pmatrix} \Rightarrow$$

Particle content and masses

- Mass terms for the $W^{\pm}$ bosons:

$$m_{W^{\pm}} = \frac{1}{2} v g$$

- Remaining terms off-diagonal in W_μ^3 and B_μ:

$$\frac{1}{8} v^2 (W_\mu^3, B_\mu) \begin{pmatrix} g^2 & -gg' \\ -gg' & g'^2 \end{pmatrix} \begin{pmatrix} W^{3\mu} \\ B^\mu \end{pmatrix} = \frac{1}{8} v^2 \left[g W_\mu^3 - g' B_\mu \right]^2$$

$$+ 0 \left[g' W_\mu^3 + g B_\mu \right]^2$$

- Massless photon:

$$A_\mu = \frac{g' W_\mu^3 + g B_\mu}{\sqrt{g^2 + g'^2}} \quad with \quad m_A = 0$$

- Massive neutral vector boson:

$$Z_\mu = \frac{g W_\mu^3 - g' B_\mu}{\sqrt{g^2 + g'^2}} \quad with \quad m_Z = \frac{1}{2} v \sqrt{g^2 + g'^2}$$

Masses of the gauge bosons:

$$\left|\left(-ig\frac{\boldsymbol{\tau}}{2}\cdot\mathbf{w}_\mu - i\frac{g'}{2}B\right)\phi\right|^2$$

$$=\frac{1}{8}\left|\begin{pmatrix} gW_\mu^3 + g'B_\mu & g(W_\mu^1 - iW_\mu^2) \\ g(W_\mu^1 + iW_\mu^2) & -gW_\mu^3 + g'B_\mu \end{pmatrix}\begin{pmatrix} 0 \\ v \end{pmatrix}\right|^2$$

$$=\frac{1}{8}v^2g^2\left[(W_\mu^1)^2 + (W_\mu^2)^2\right] + \frac{1}{8}v^2(g'B_\mu - gW_\mu^3)(g'B^\mu - gW^{3\mu})$$

$$=\left(\frac{1}{2}vg\right)^2 W_\mu^+W^{-\mu} + \frac{1}{8}v^2(W_\mu^3, B_\mu)\begin{pmatrix} g^2 & -gg' \\ -gg' & g'^2 \end{pmatrix}\begin{pmatrix} W^{3\mu} \\ B^\mu \end{pmatrix}$$

Important relations in the Glashow-Salam-Weinberg model:

- Relation between the gauge couplings: $$\frac{g'}{g} = \tan\theta_W$$

→ Important prediction of the GSW with a Higgs doublet: $$\frac{m_W}{m_Z} = \cos\theta_W$$

or expressed in terms of the ρ parameter: $$\rho \equiv \frac{m_W^2}{m_Z^2 \cos^2\theta_W} = 1$$

- From the M_W relation the value of the vacuum expectation value of the Higgs field can be calculated:

$$\frac{1}{2v^2} = \frac{g^2}{8m_W^2} = \frac{G_F}{\sqrt{2}} \qquad \rightarrow \quad v = 246\ GeV$$

where G_F = Fermi constant, know from low energy experiments (muon decay)

Masses of the Fermions:

- The same Higgs doublet which generates $W^{\pm}$ and Z masses is sufficient to give masses to the fermions (leptons and quarks):
  e.g. for electrons: use an arbitrary coupling G_e

$$L_3 = -G_e \left[(\bar{\nu}_e, \bar{e})_L \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} e_R + \bar{e}_R (\phi^-, \bar{\phi}^0) \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L \right]$$

- Spontaneous symmetry breaking:

$$\phi = \sqrt{\frac{1}{2}} \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix}$$

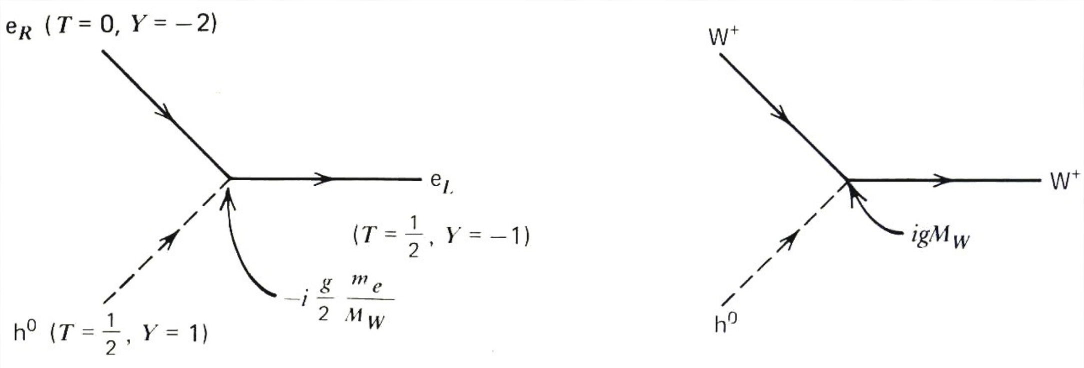
$$L_3 = -\frac{G_e v}{\sqrt{2}} (\bar{e}_L e_R + \bar{e}_R e_L) - \frac{G_e}{\sqrt{2}} (\bar{e}_L e_R + \bar{e}_R e_L) h$$

mass term

interaction term with the Higgs field

- Important relation: coupling of the Higgs boson to fermions is proportional to their mass

$$G_f = \frac{\sqrt{2}\, m_f}{v}$$



and finally...... a massive scalar with self-coupling, the **Higgs boson**:

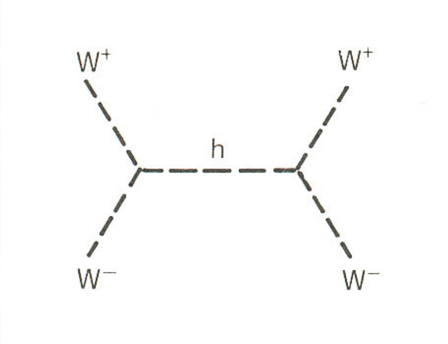
• <u>Mass:</u> $$m_h^2 = 2v^2\lambda$$

(since λ is not predicted by theory, the mass of the Higgs boson is unknown)

• <u>Self-coupling</u>: $$-\lambda v h^3 - \frac{1}{4}\lambda h^4$$

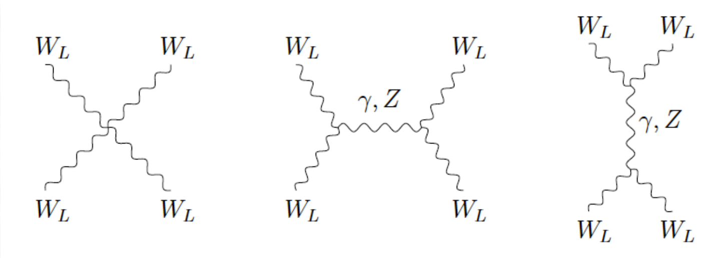
….. and:

• The additional diagram, with Higgs boson exchange, regulates the divergences in the longitudinal WW scattering



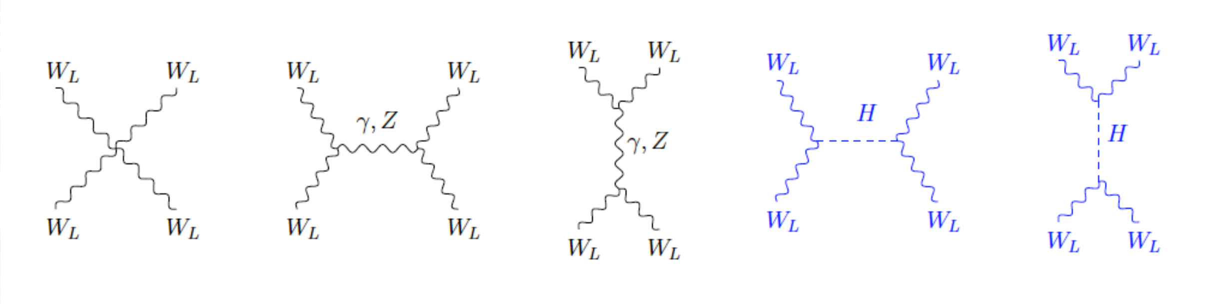
The Higgs boson as a UV regulator

Scattering of longitudinally polarized W bosons



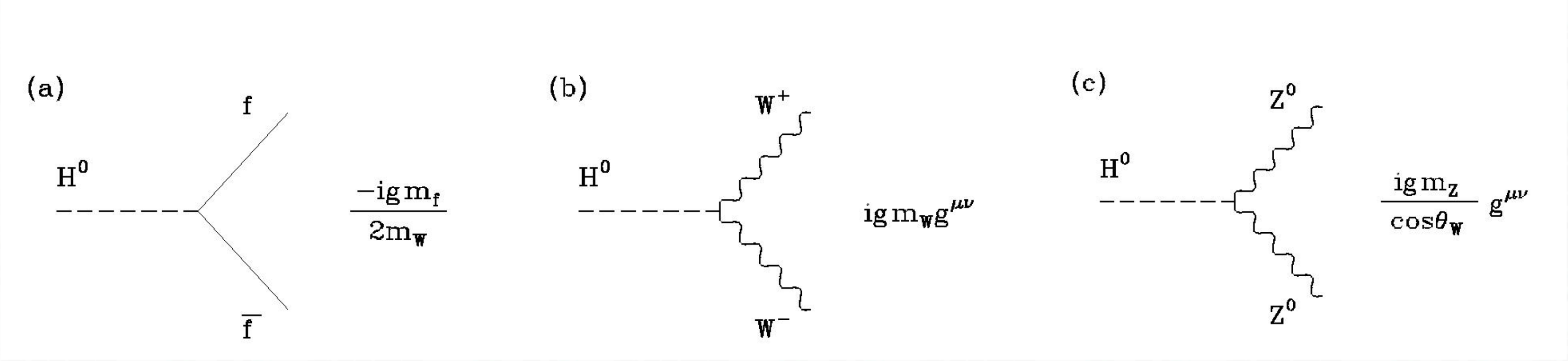
$$-iM(W^+W^- \rightarrow W^+W^-) \sim \frac{s}{m_W^2} \quad \text{for} \quad s \rightarrow \infty$$

Higgs boson guarantees unitarity (if its mass is < ~1 TeV)



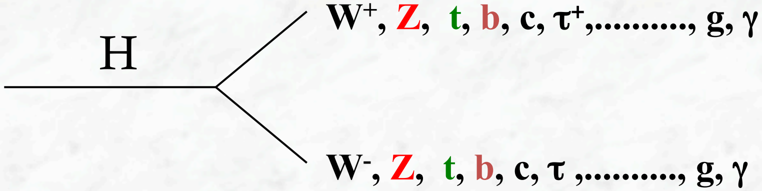
$$-iM(W^+W^- \rightarrow W^+W^-) \sim m_H^2 \quad \text{for} \quad s \rightarrow \infty$$

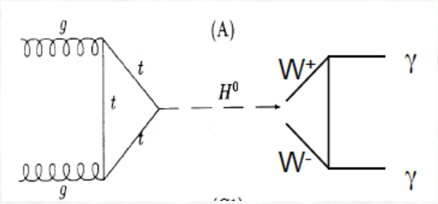
7.2 Higgs boson properties



Higgs Boson Decays

The decay properties of the Higgs boson are fixed, if the **mass** is known:





$$\Gamma(H \to f\bar{f}) = N_C \frac{G_F}{4\sqrt{2}\pi} m_f^2(m_H^2) m_H$$

$$\Gamma(H \to VV) = \delta_V \frac{G_F}{16\sqrt{2}\pi} m_H^3 (1 - 4x + 12x^2)\beta_V$$

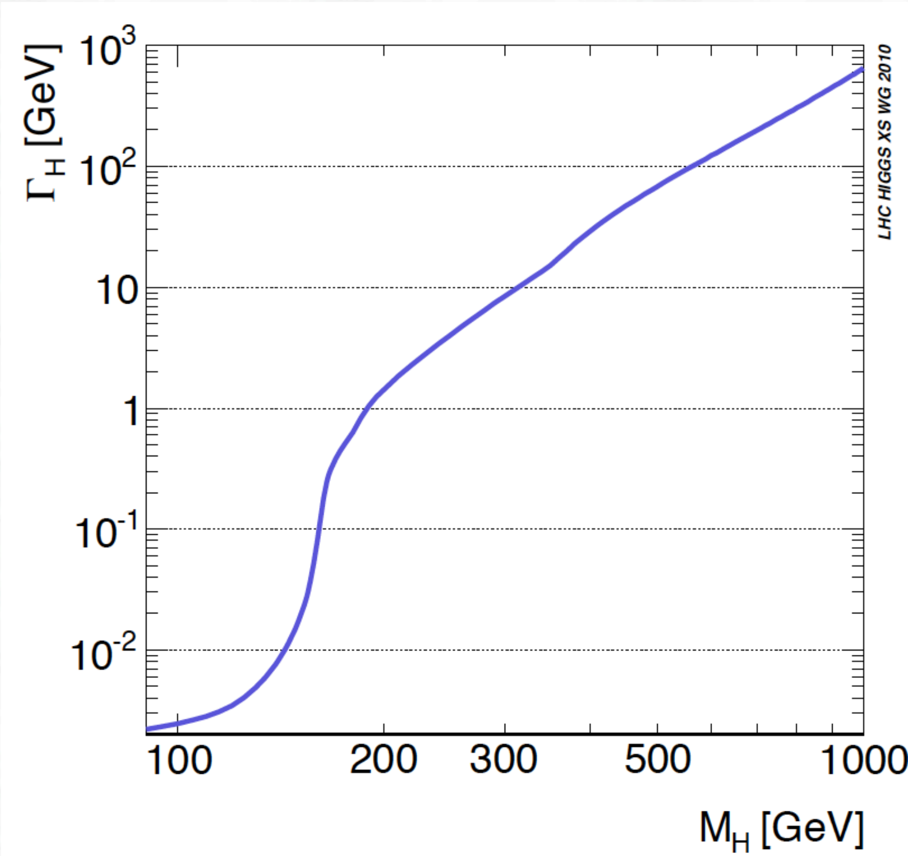
where: $\delta_Z = 1,\ \delta_W = 2,\ x = m_V^2 / m_H^2, \quad \beta = \text{velocity}$

(+ W-loop contributions)

$$\Gamma(\mathrm{H} \to \mathrm{gg}) = \frac{G_F \alpha_a^2(m_H^2)}{36\sqrt{2}\pi^3} m_H^3 \left[1 + \left(\frac{95}{4} - \frac{7N_f}{6}\right)\frac{\alpha_a}{\pi}\right]$$

$$\Gamma(H \to \gamma\gamma) = \frac{G_F \alpha_a^2}{128\sqrt{2}\pi^3} m_H^3 \left[\frac{4}{3} N_C e_t^2 - 7\right]^2$$

Total width



Higgs Boson Decays

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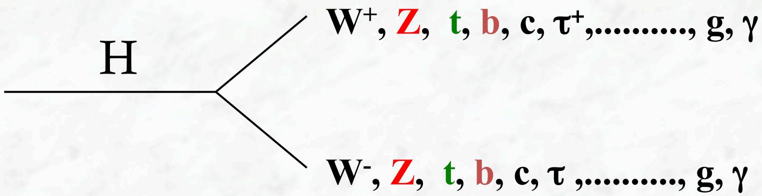
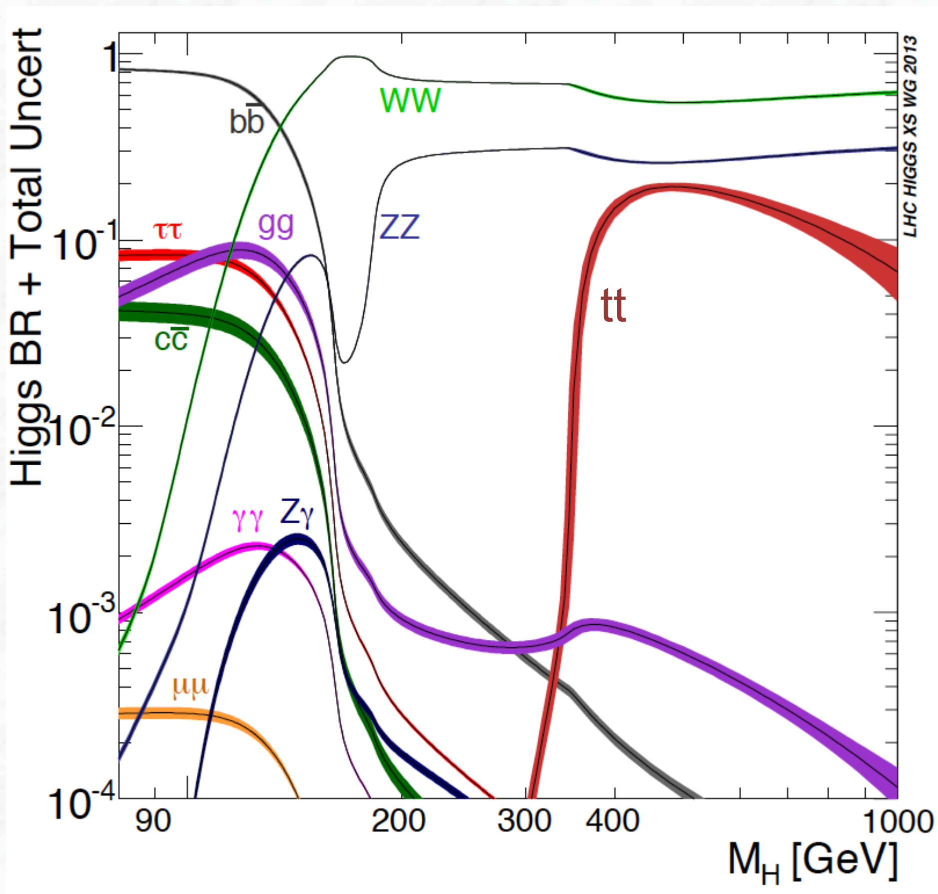




Constraints on the Higgs boson mass

1. Constraints from theory

2. Indirect limits from electroweak precision data (theory and experiment)

3. Limits from Direct Searches (LEP)

(i) Theory Constraints on the Higgs boson mass

- Unitarity limit:

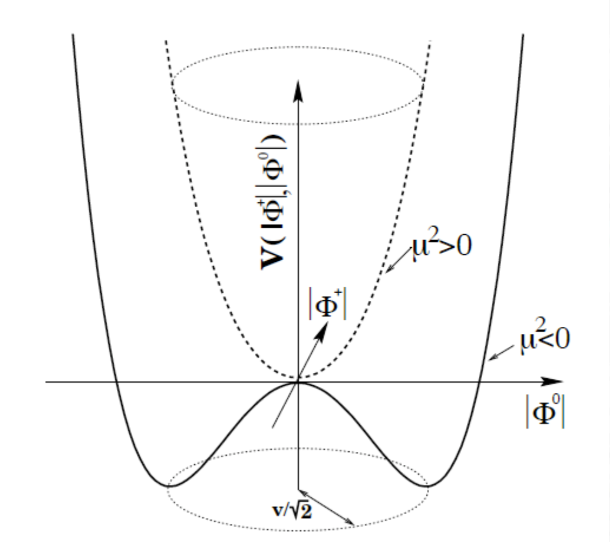
  If Higgs boson too heavy, the regulation of the WW cross section is less effective and unitarity is violated again

  → $m_H <$ ~1 TeV (as just discussed)

- Stricter limits from the energy dependence of the Higgs boson self coupling λ

  – Stability of the vacuum
  - Diverging coupling $\lambda(Q^2)$

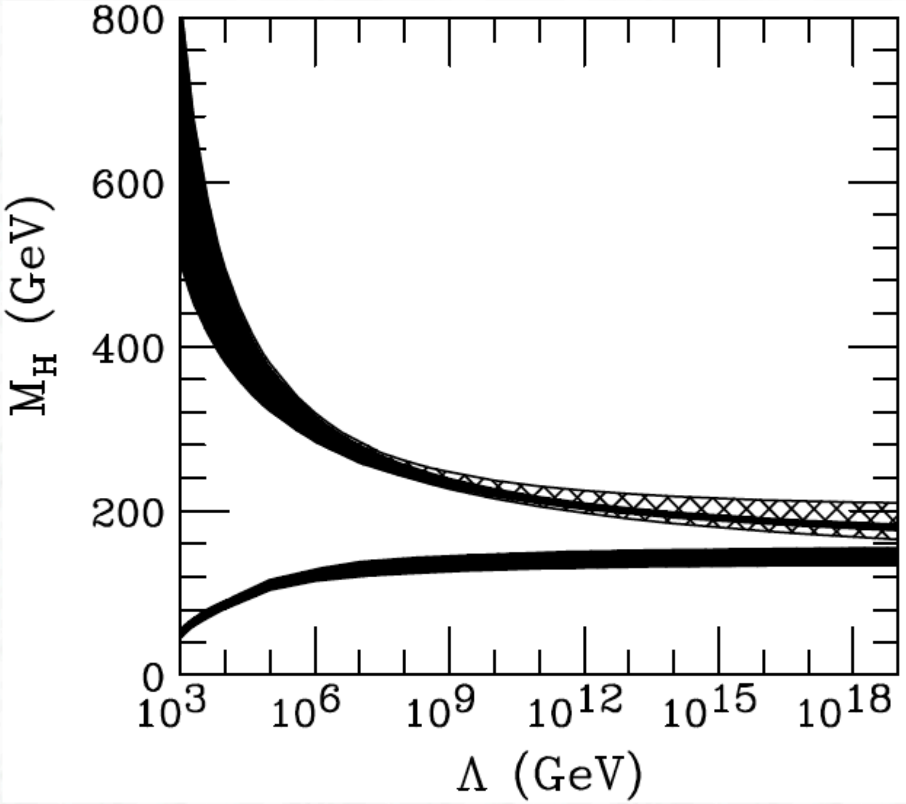
  → next slides



Tighter Higgs mass constraints:

Stronger bounds on the Higgs-boson mass result from the energy dependence of the Higgs coupling $\lambda\,(Q^2)$
(if the Standard Model is assumed to be valid up to some scale Λ)

$$\lambda(Q^2) = \lambda_0 \left\{ 1 + \frac{3\lambda_0}{2\pi^2} \log\left(2\frac{Q^2}{v^2}\right) + \ldots - \frac{3g_t^4}{32\pi^2} \log\left(2\frac{Q^2}{v^2}\right) + \ldots \right\} \qquad \text{where} \quad \lambda_0 = \frac{m_h^2}{v^2}$$



Hambye, Riesselmann (1996)

Upper bound: diverging coupling (Landau Pole)

Lower bound: stability of the vacuum (negative contribution from top quark dominates)

Mass bounds depend on scale Λ up to which the Standard Model should be valid

(ii) Indirect limits from electroweak precision data (m$_W$ and m$_t$)

Motivation:

W mass and top quark mass are fundamental parameters of the Standard Model;
The standard theory provides well defined relations between m_W, m_t and m_H

Electromagnetic constant
measured in atomic transitions,
e^+e^- machines, etc.

$$m_W = \left(\frac{\pi \, \alpha_{EM}}{\sqrt{2} \, G_F} \right)^{1/2} \frac{1}{\sin\theta_W \sqrt{1-\Delta r}}$$

Fermi constant
measured in muon
decay

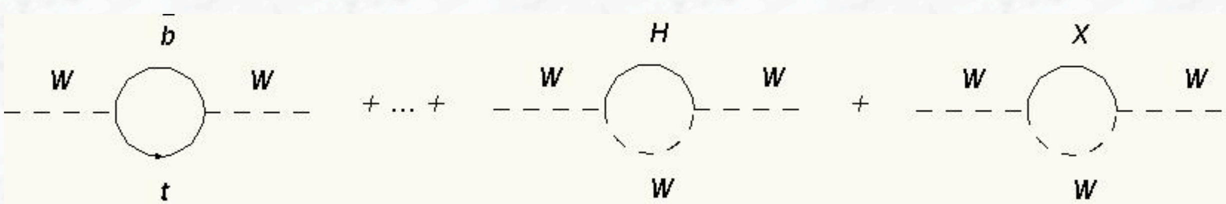
weak mixing angle
measured at
LEP/SLC

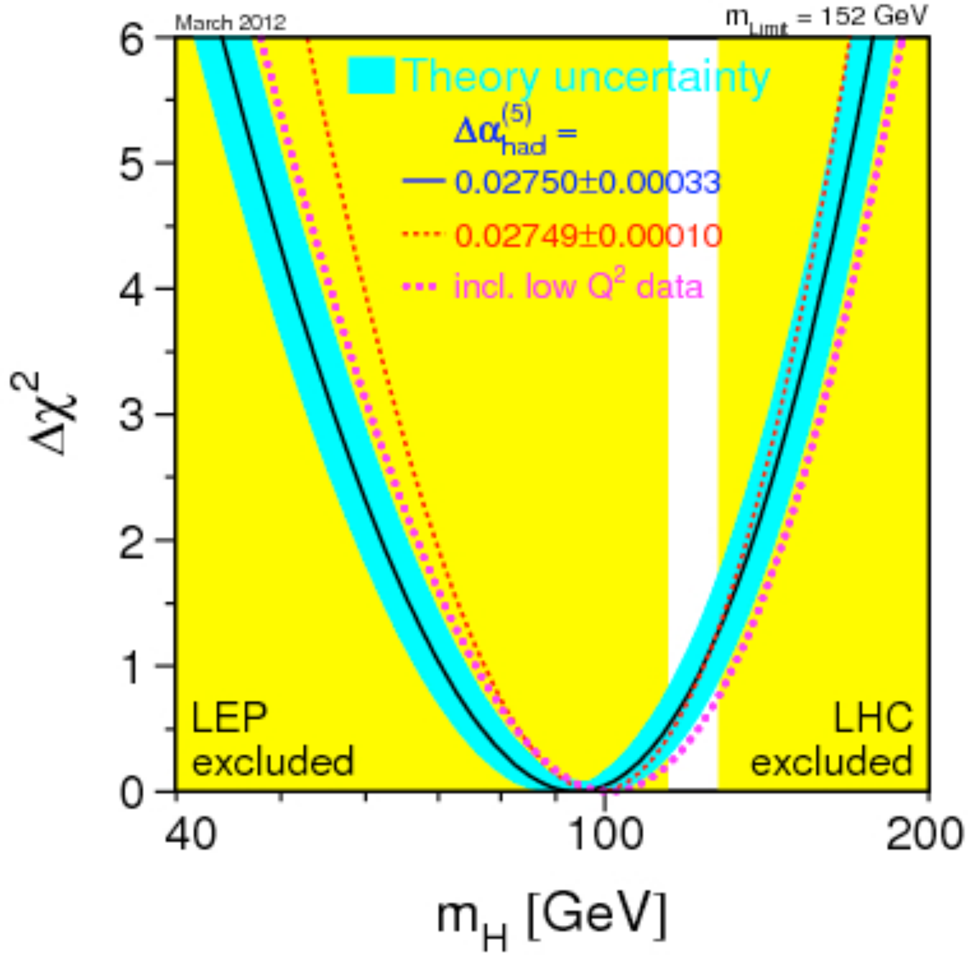
radiative corrections
$\Delta r \sim f(m_t^2, \log m_H)$
$\Delta r \approx 3\%$

G_F, α_{EM}, $\sin\theta_W$

are known with high precision

Precise measurements of the W mass and the top-quark mass constrain the Higgs-boson mass
(and/or the theory, radiative corrections)





Results of the precision el.weak measurements: (LEWWG-2012):

$$m_H = 94^{+29}_{-24}\,GeV$$

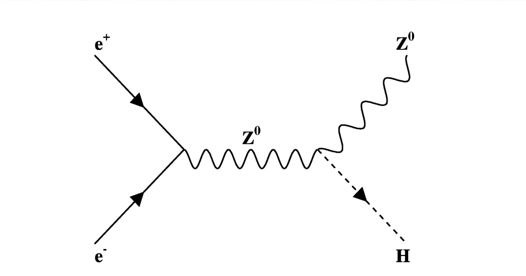
(iii) Constraints from

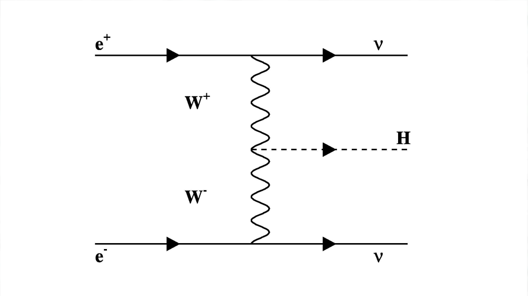
direct searches at LEP

Higgs bosons searches at LEP

Higgs-Strahlung: e+ e- → Z H

WW-Fusion: e+ e- → νν H





Higgs decay branching ratios for m_H=115 GeV/c^2:
BR (H → bb) = 74%, BR (H→ ττ, WW, gg) = 7% each, BR(H → cc) = 4%

Decay modes searched for:

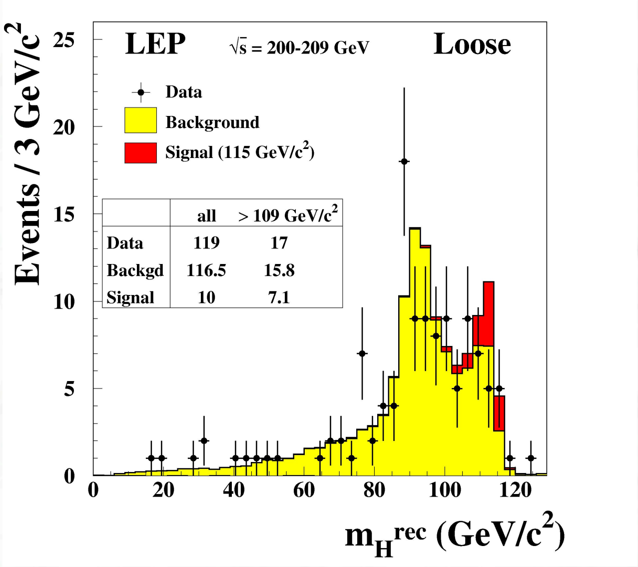
- Four Jet channel: HZ → bb qq
- Missing energy channel: → bb νν
- Leptonic channel: → bb ee, bb μμ
- Tau channels: → bb ττ, and ττ qq

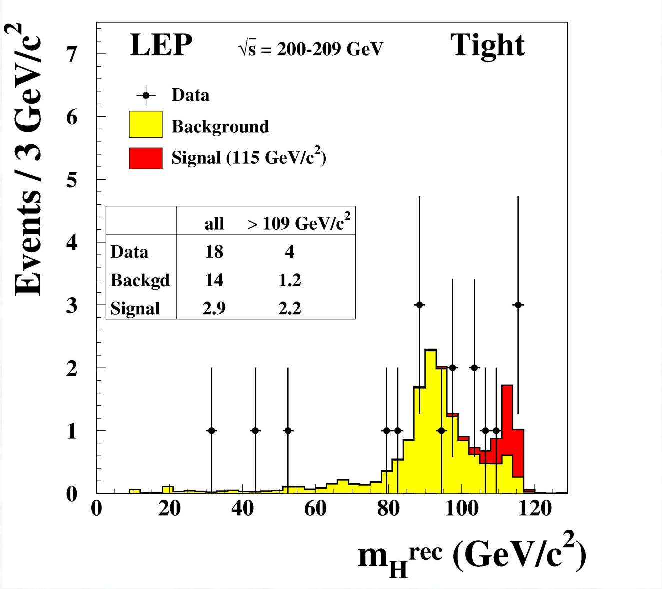
Results of the final LEP analysis:

Final results have been published: CERN-EP / 2003-011:

Based on final calibrations of the detectors, LEP-beam energies, final Monte Carlo simulations and analysis procedures.

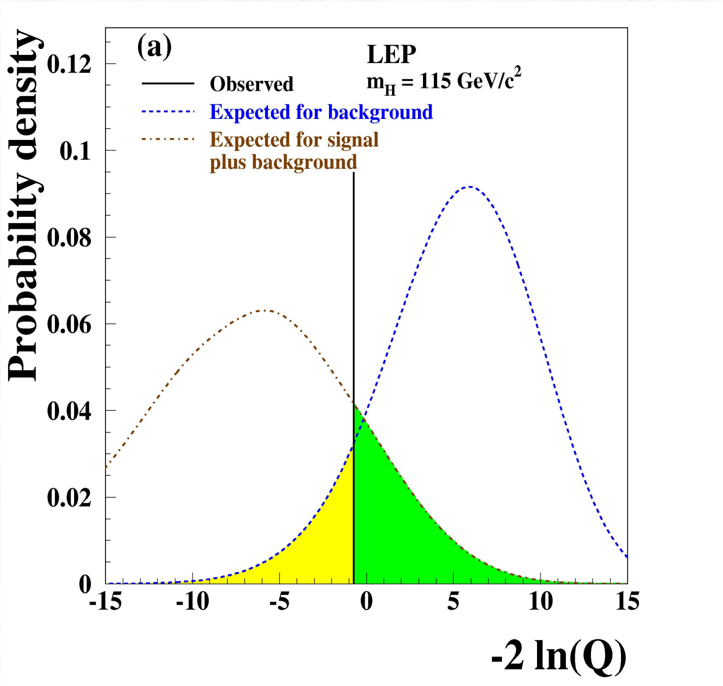
The reconstructed bb mass for two levels of signal purity (loose and tight cuts):





Clear peak in the background prediction in the vicinity of m_Z due to the $e^+e^- \rightarrow ZZ$ background, which is consistent with the data.

Final combined LEP result



	$1 - CL_b$	CL_{s+b}
LEP	0.09	0.15
ALEPH	3.3×10^{-3}	0.87
DELPHI	0.79	0.03
L3	0.33	0.30
OPAL	0.50	0.14
Four-jet	0.05	0.44
All but four-jet	0.37	0.10

$1- CL_B = 0.09$ ↔

Signal significance = 1.7 σ

Likelihood ratio Q : = L_{S+B} / L_B
Test statistics: - 2 ln Q

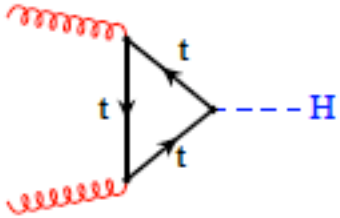
M_H > 114.4 GeV/c^2 (95% CL)

expected mass limit: 115.3 GeV/c^2
(sensitivity)

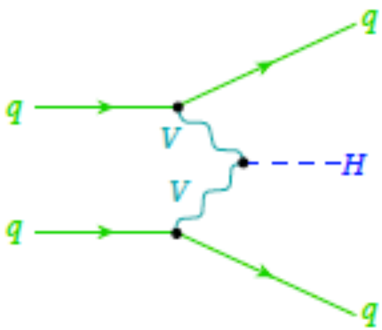
7.3 Higgs boson production at Hadron Colliders

Higgs Boson production processes at Hadron Colliders

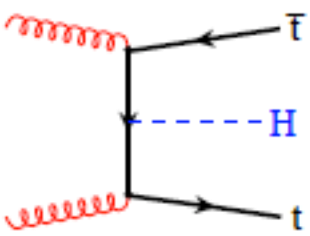
Relative importance of the various processes is different at the LHC and at the Tevatron



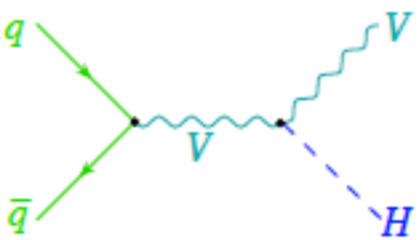
Gluon Fusion



Vector boson fusion

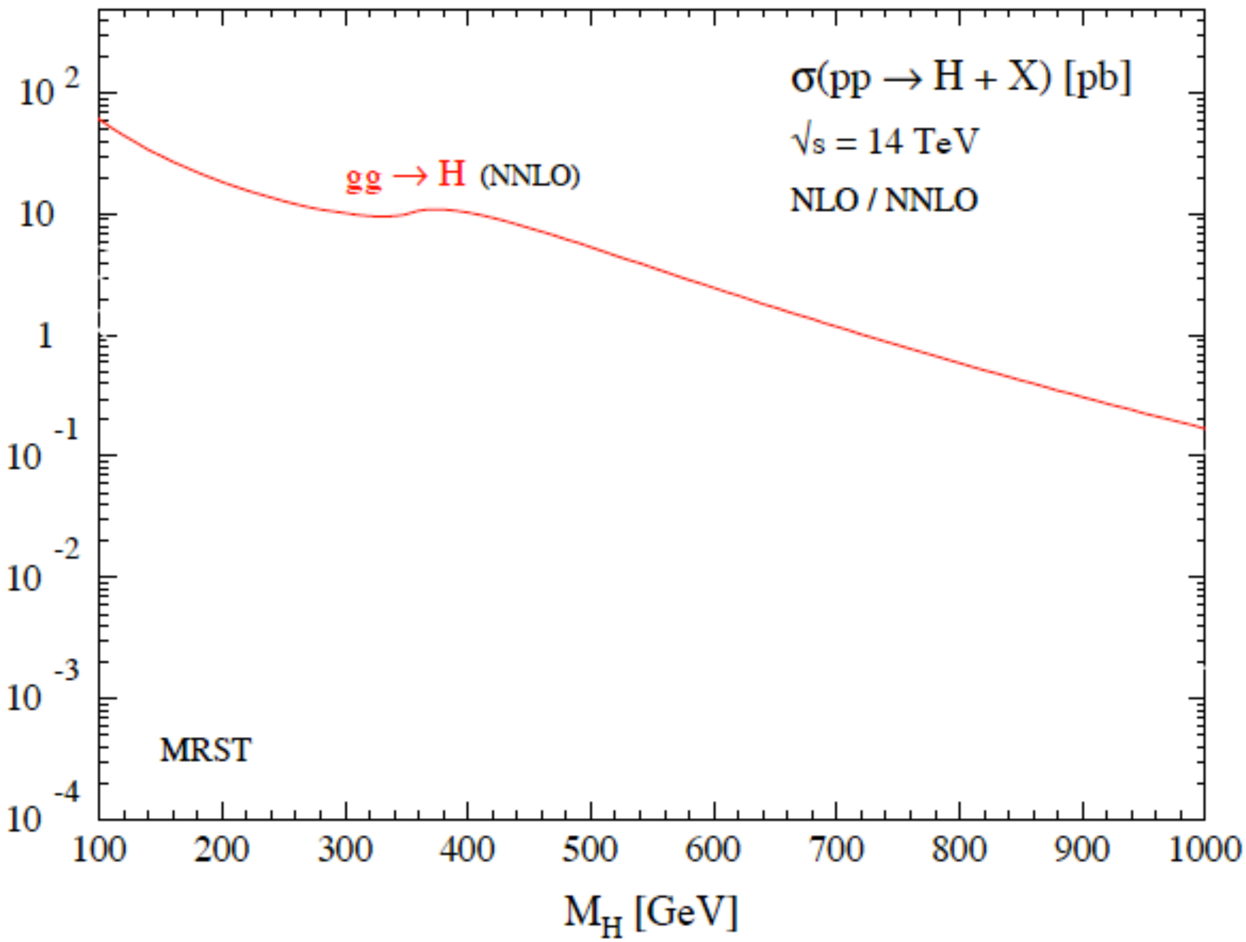


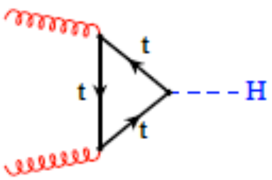
tt associated production



WH/ZH associated production

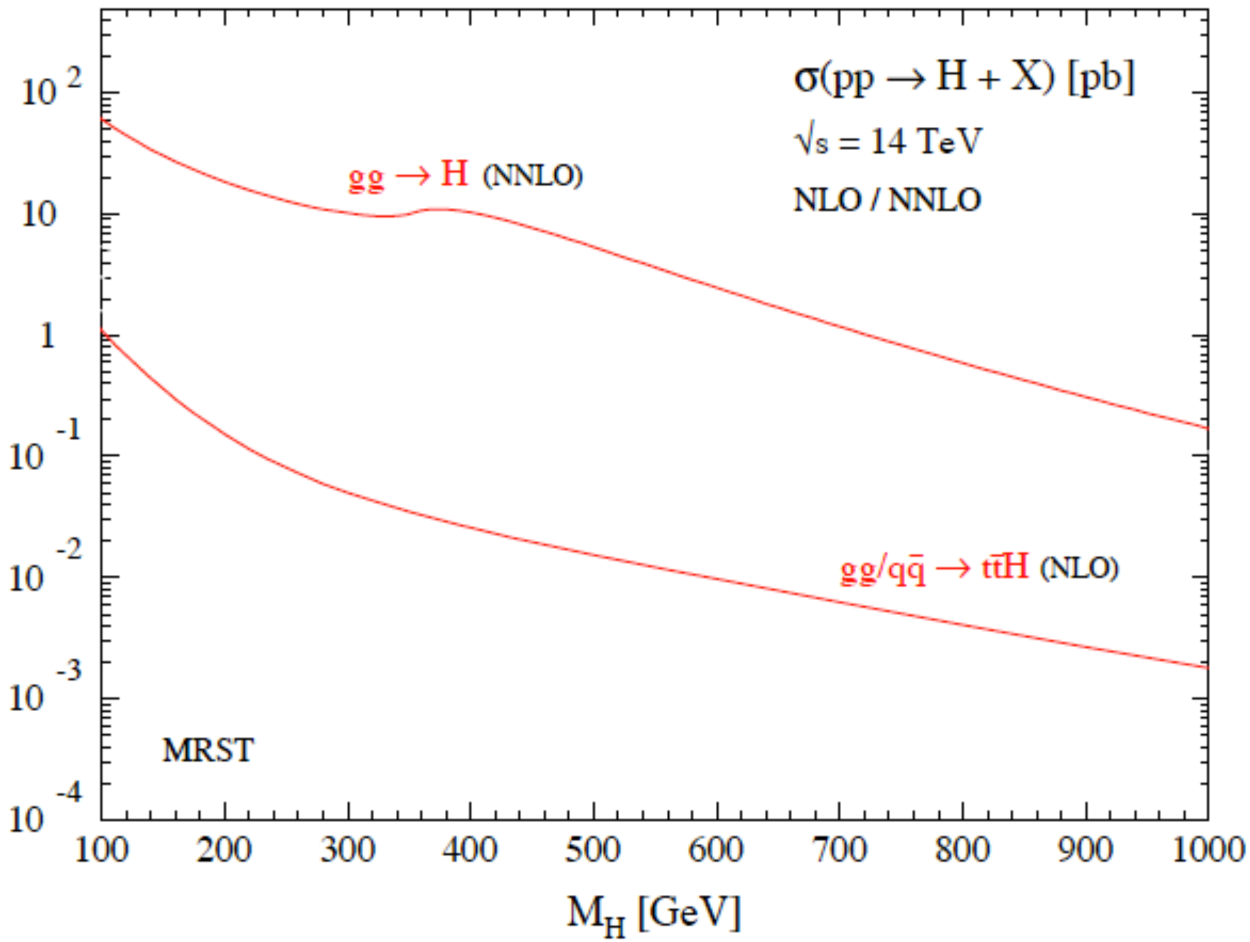
Production cross sections at the LHC

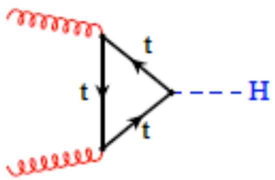


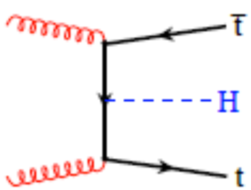


(for $\sqrt{s}$ = 14 TeV)

Production cross sections at the LHC

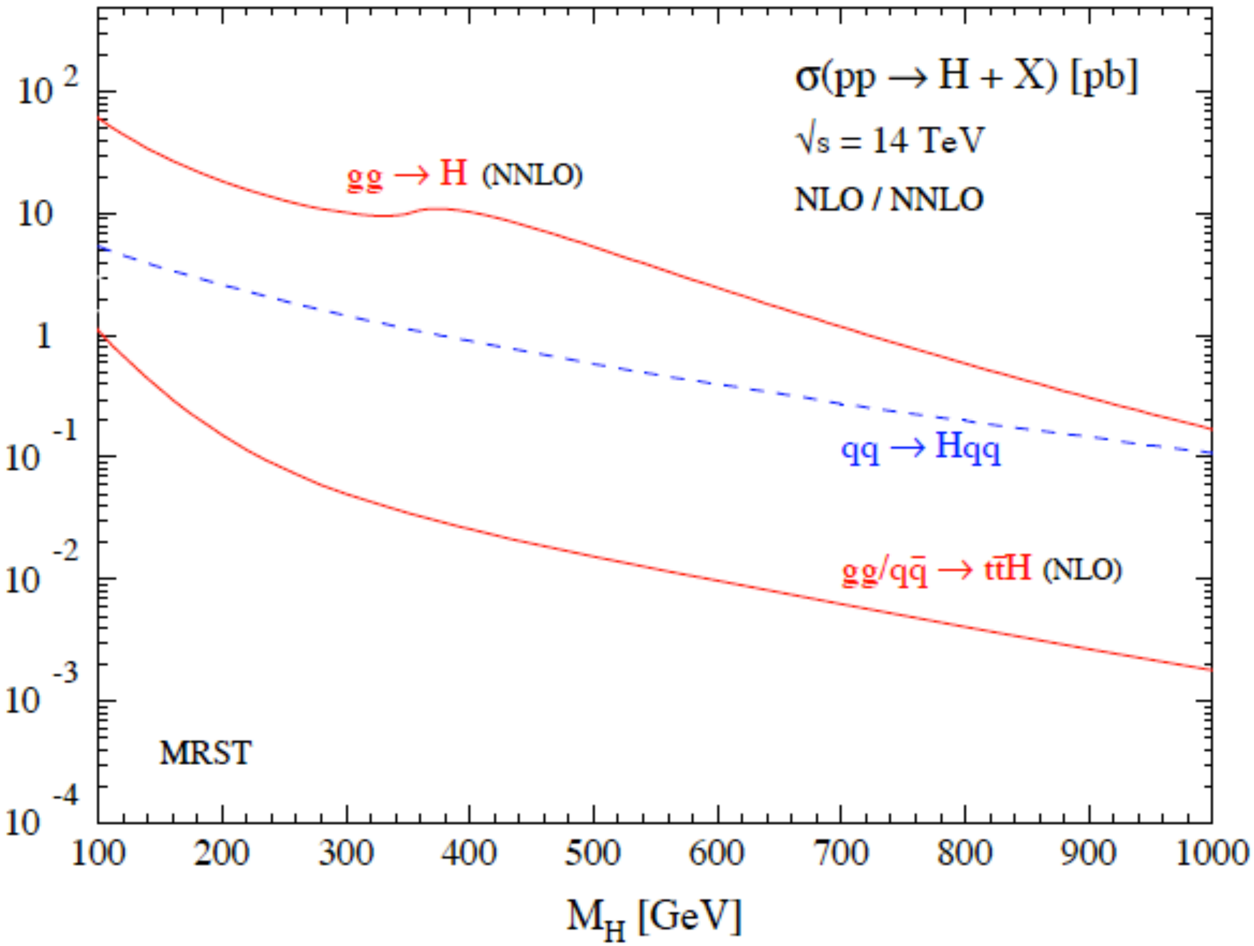


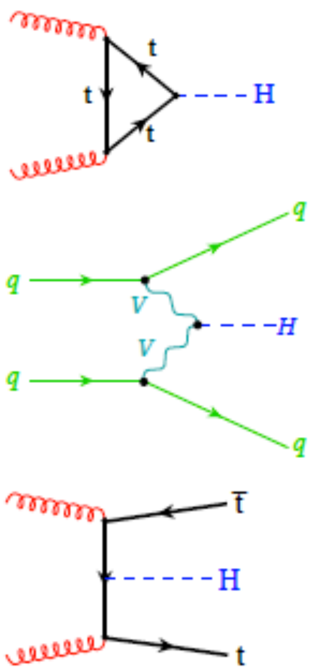




(for $\sqrt{s}$ = 14 TeV)

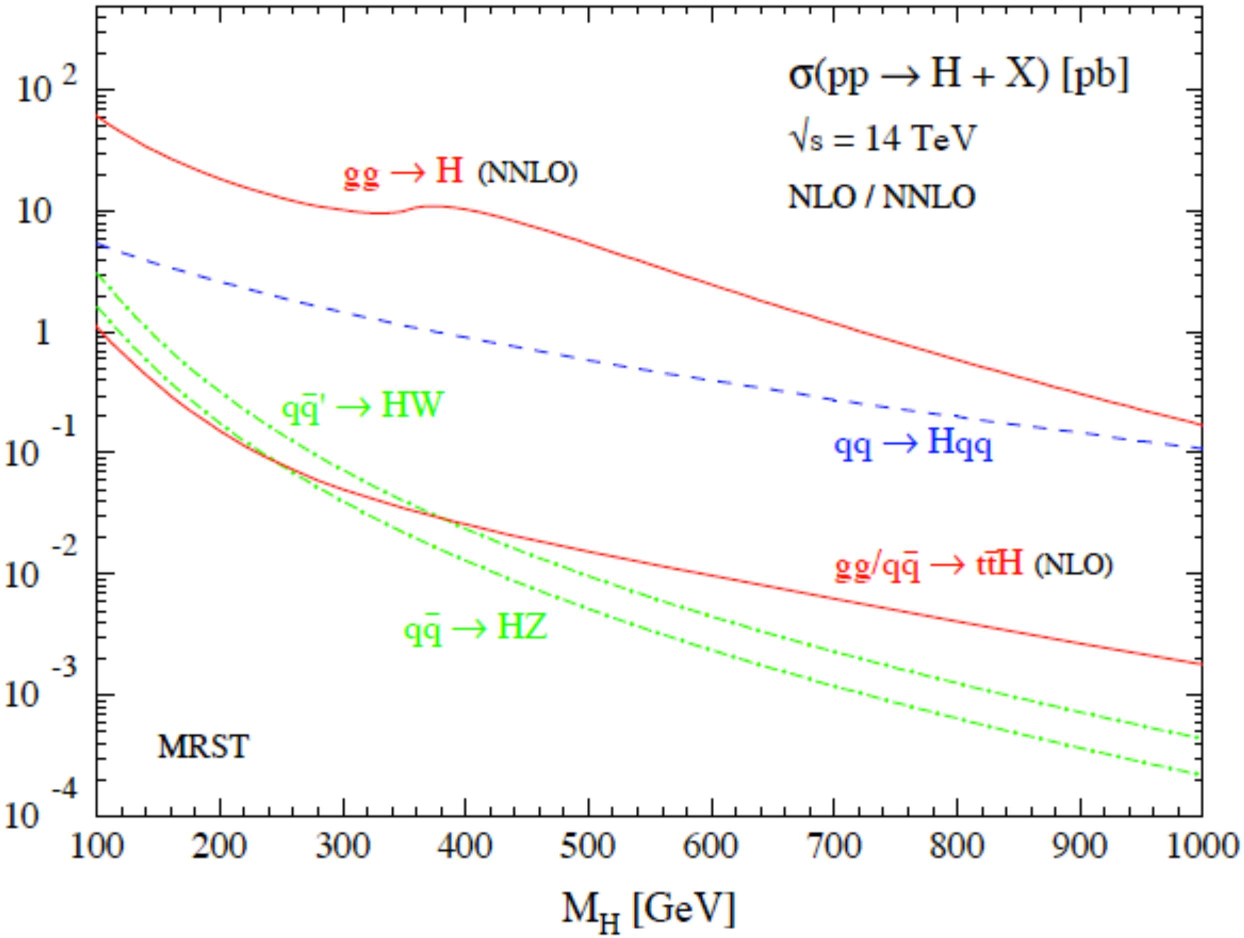
Production cross sections at the LHC

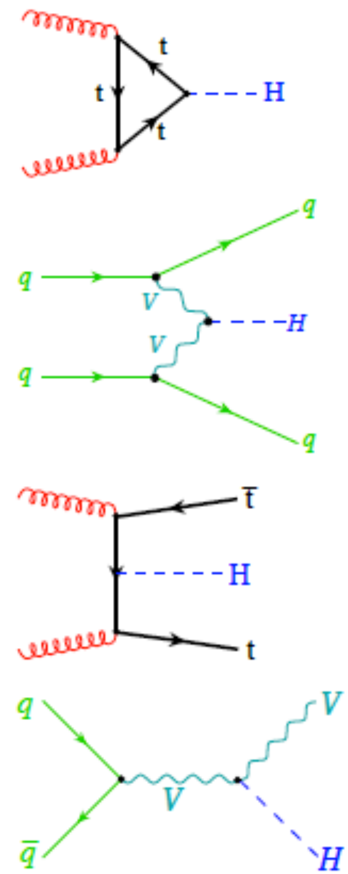




(for √s = 14 TeV)

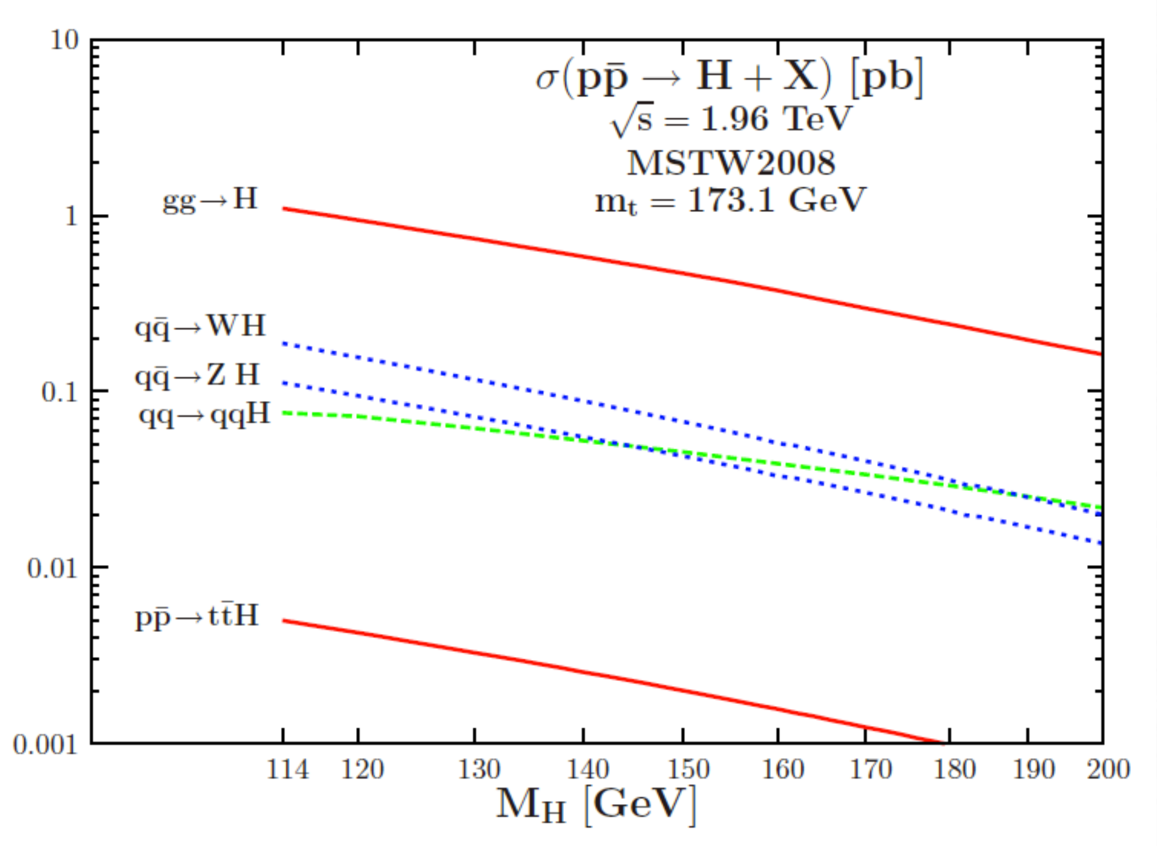
Production cross sections at the LHC





(for √s = 14 TeV)

Production cross sections at the Tevatron



J. Baglio, A. Djouadi, arXiv:1003.4266

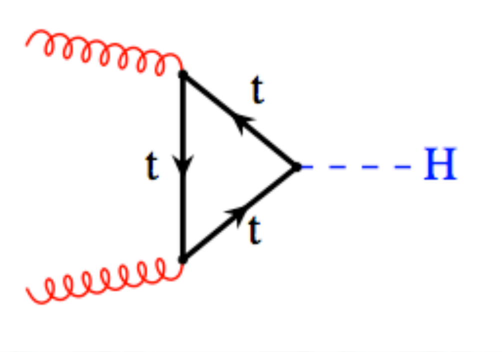
1. Gluon fusion

2./3. W/Z H associated production
Vector boson fusion

4. ttH (very small cross section)

qq → W/Z + H cross sections ~10 x larger at the LHC ($\sqrt{s}$ = 14 TeV)
gg → H ~70-80 x larger at the LHC ($\sqrt{s}$ = 14 TeV)

Gluon fusion:

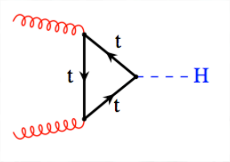


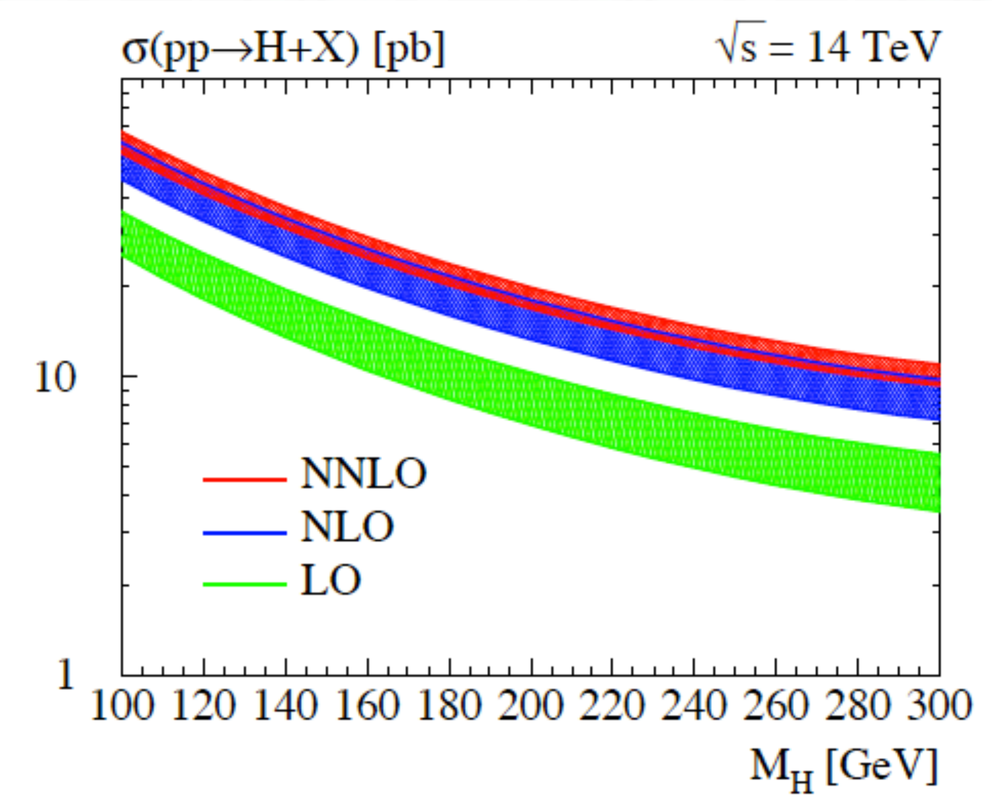
- Dominant production mode

- Sensitive to heavy particle spectrum …
  (e.g. 4^{th} generation quarks)
  …and the corresponding Yukawa couplings
  (important for coupling measurements, top Yukawa coupling)

- Large K-factors (NLO, NNLO corrections)
  - Difficult to calculate, loop already at leading order
    (calculation with infinite top mass is used as an approximation, however, this seems to be a good approximation)
  - Nicely converging perturbative series

Higher order corrections:



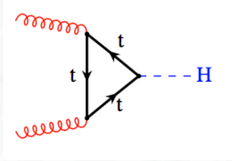


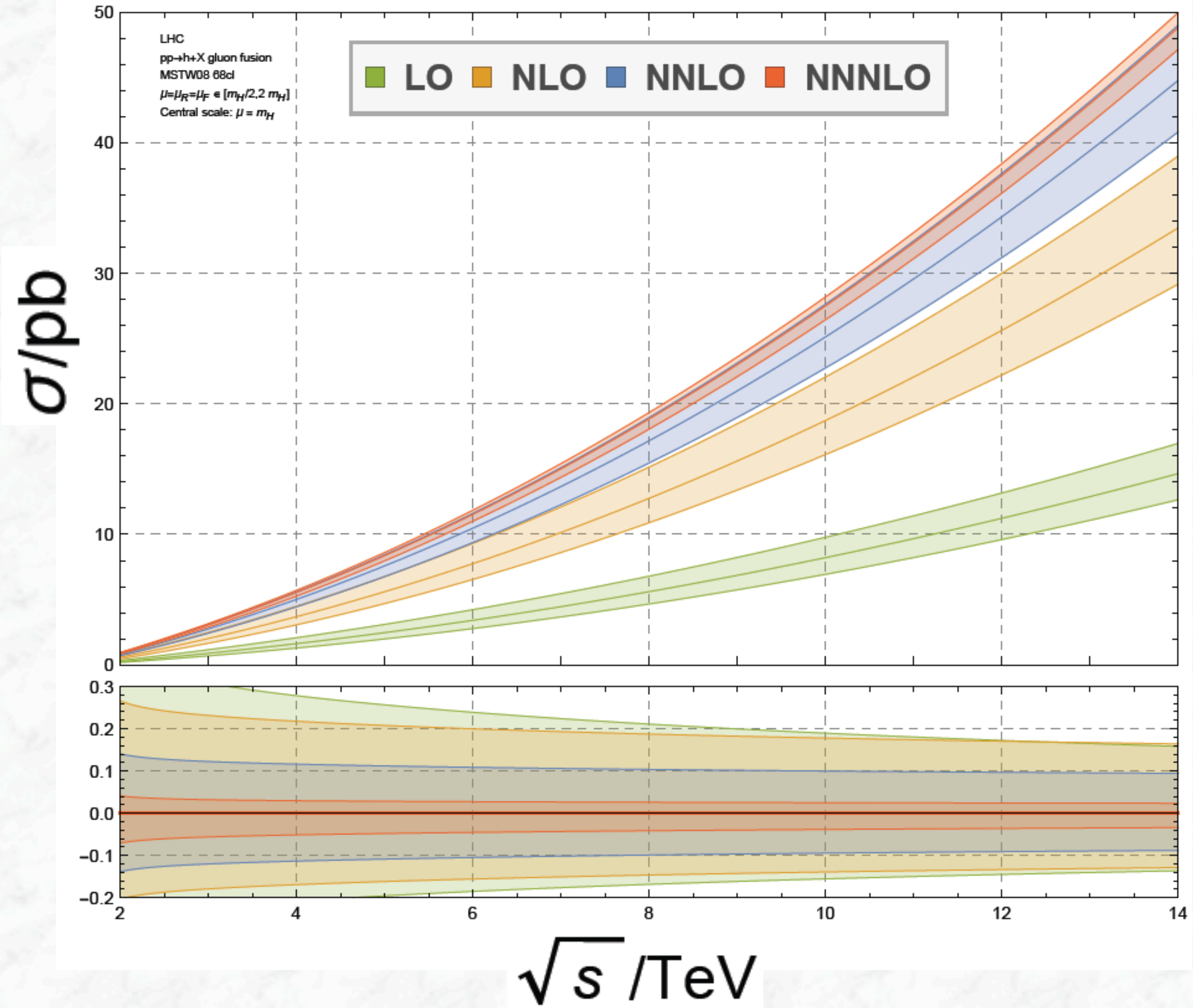
- Spira, Djouadi, Graudenz, Zerwas (1991)
- Dawson (1991)

- Harlander, Kilgore (2002)
- Anastasiou, Melnikov (2002)
- Ravindran, Smith, van Neerven (2003)

Independent variation of renormalization and factorization scales (with $0.5\ m_H < \mu_F, \mu_R < 2\ m_H$)

Higher order corrections:

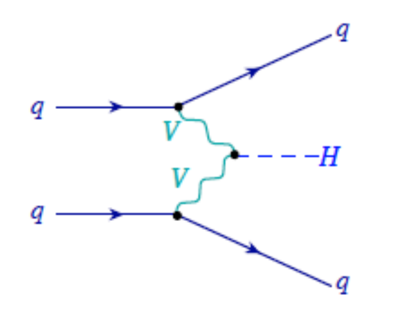




-Mistlberger et al. (2015)

Independent variation of renormalization and factorization scales (with $0.5\ m_H < \mu_F, \mu_R < 2\ m_H$)

Vector boson fusion:



- Second largest production mode,
  Distinctive signature
  (forward jets, little jet activity in the central region)

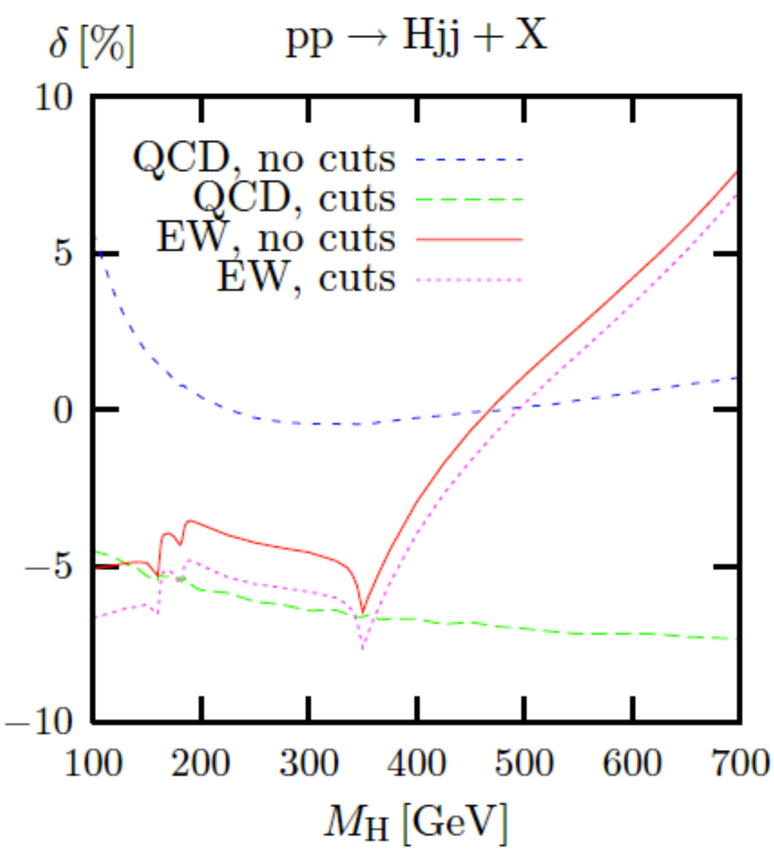
- Sensitivity to W/Z couplings

- Moderate K-factors
  (NLO corrections)

  Both NLO QCD and el.weak have been calculated

- Effective K-factor depends on experimental cuts

  Example: typical VBF cuts
  P_T(jet) > 25 GeV
  $\eta < 4.5$, $\Delta\eta > 4$, $\eta_1 \cdot \eta_2 < 0$



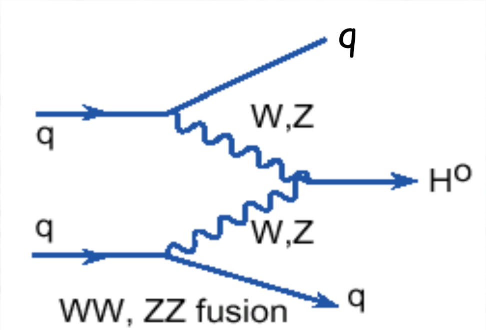
Ciccolini, Denner, Dittmaier (2008)

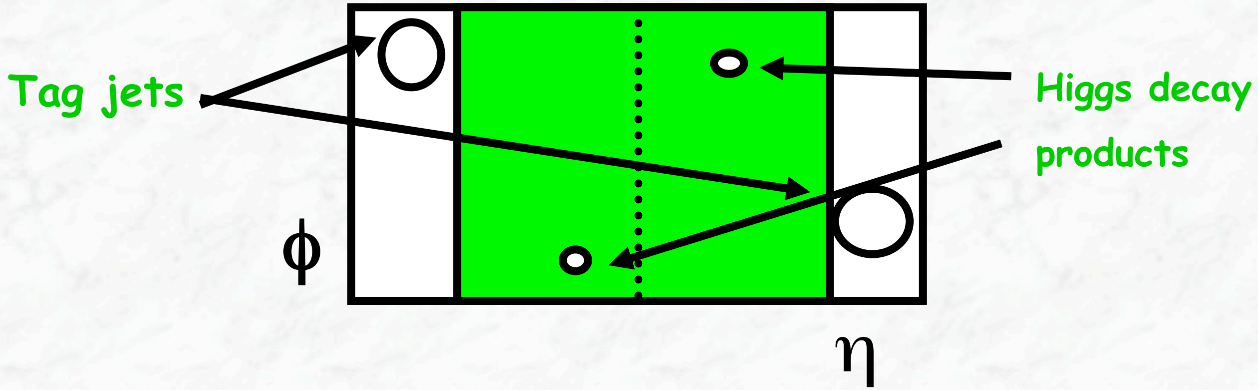
Vector boson fusion:

Motivation: Increase discovery potential at low mass
Improve and extend measurement of Higgs boson parameters
(couplings to bosons, fermions)

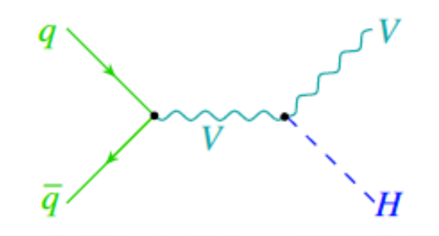
Established (low mass region) by D. Zeppenfeld et al. (1997/98)
Earlier studies: R.Kleiss W.J.Stirling, Phys. Lett. 200 (1988) 193;
Dokshitzer, Khoze, Troyan, Sov.J. Nucl. Phys. 46 (1987) 712;
Dokshitzer, Khoze, Sjöstrand, Phys.Lett., B274 (1992) 116.

Distinctive Signature of:
- two high p_T forward jets (tag jets)
- little jet activity in the central region (no colour flow)
  $\Rightarrow$ central jet Veto





WH / ZH associated production:



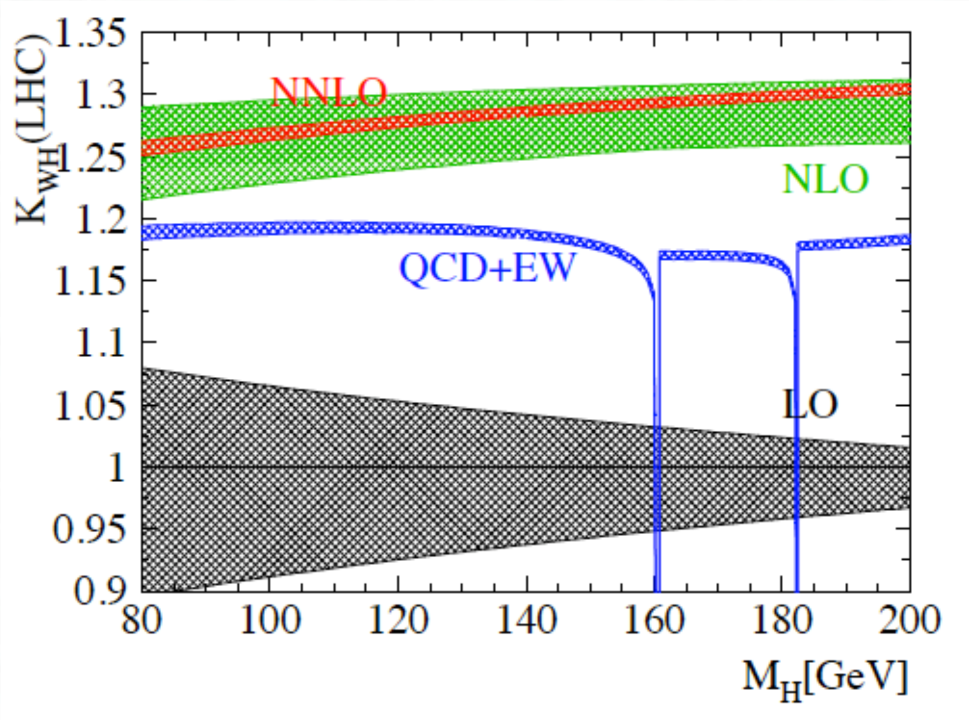
- Weak at the LHC,
  Relatively stronger at the Tevatron

- Allows for a Higgs-decay-independent trigger
  W → lν, Z → ll

- Sensitivity to W/Z couplings

- Moderate K-factors
  (NLO corrections)

  Both NLO QCD and el.weak
  corrections available

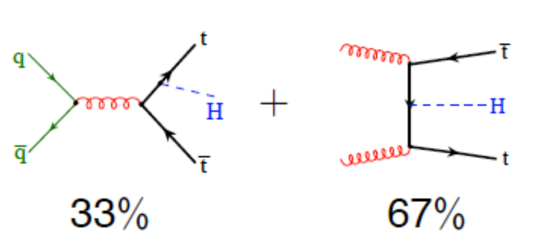


Brein, Djouadi, Harlander, (2003)

Han, Willenbrock (1990)

Ciccolini, Dittmaier, Krämer (2003)

ttH associated production:

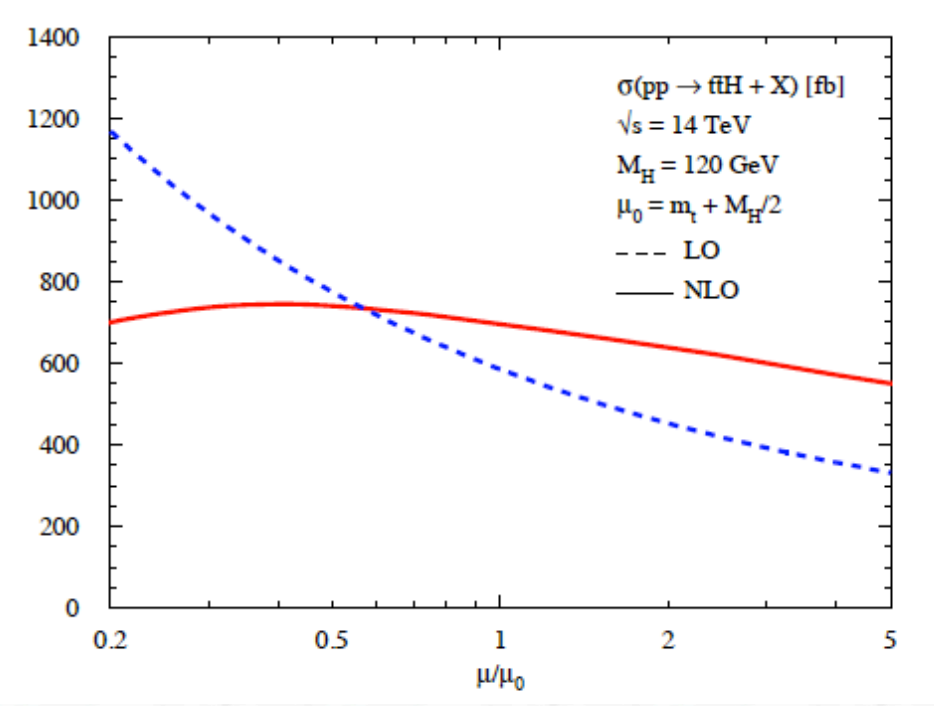


- Weak and difficult at the LHC
- Sensitivity to top-Yukawa coupling
- Moderate K-factors (NLO corrections)

NLO QCD corrections available, scale uncertainty drastically reduced

scale: $\mu_0 = m_t + m_H/2$

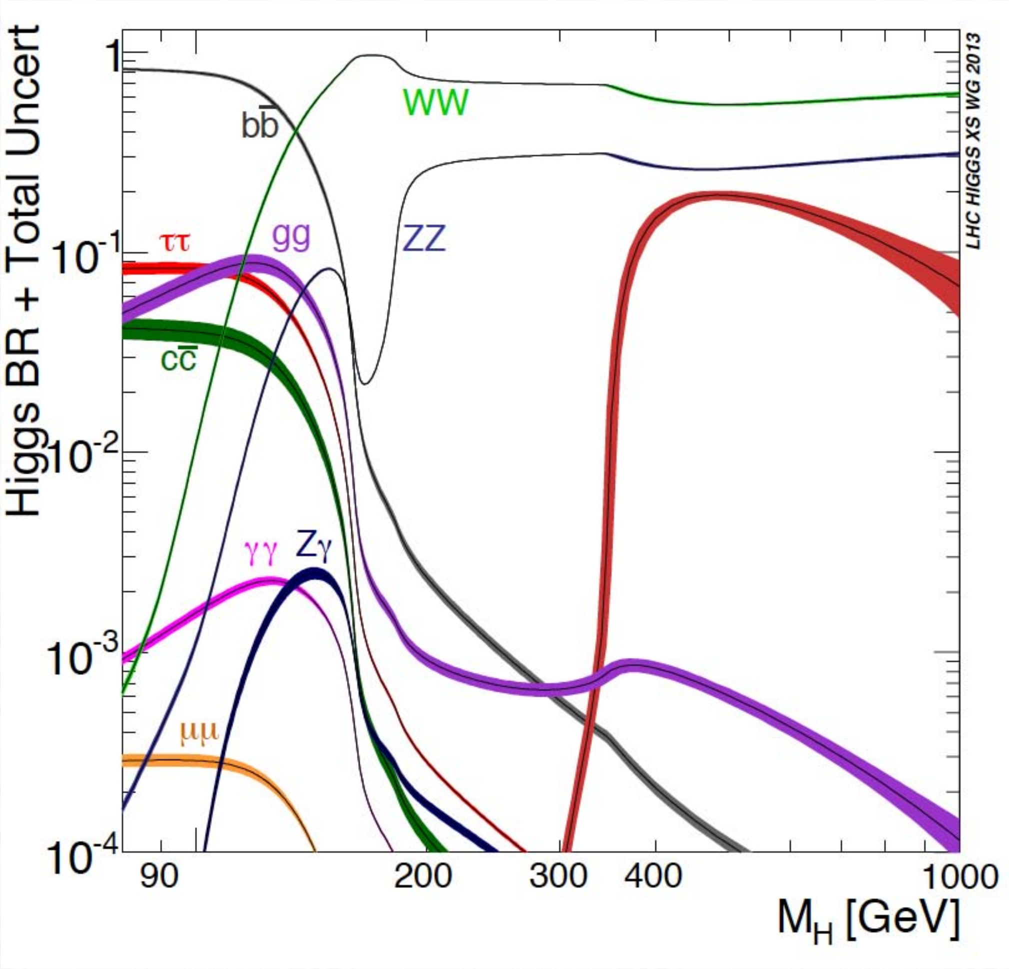
LHC: K ~ 1.2
Tevatron: K ~ 0.8



Beenakker, Dittmaier, Krämer, Plümper, Spira, Zerwas (2001)
Dawson, Reina, Wackeroth, Orr, Jackson (2001, 2003)

7.4 The discovery of a Higgs boson at the LHC

Useful Higgs Boson Decays at Hadron Colliders



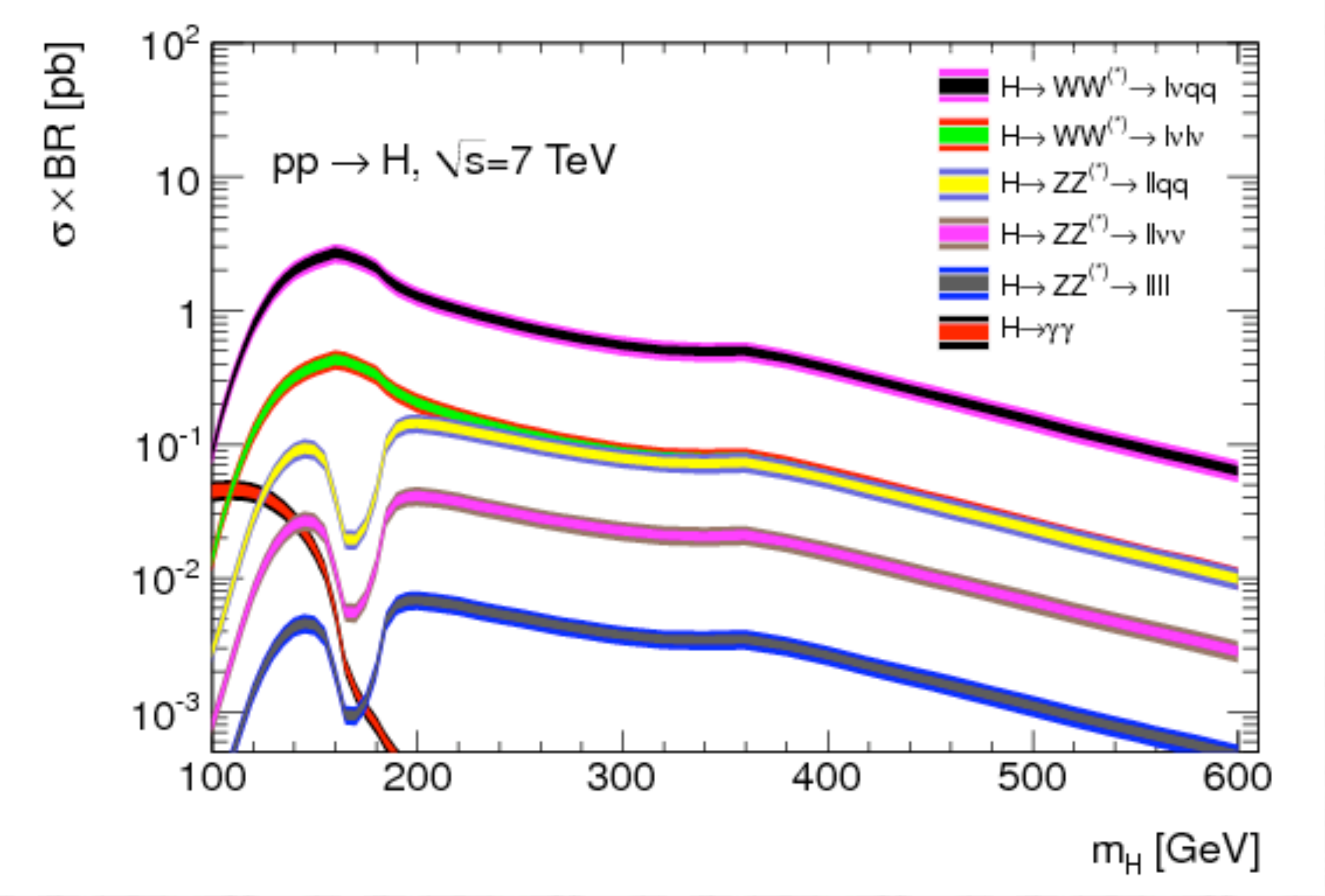
at high mass:
Lepton final states
(via H → WW , ZZ)

at low mass:
Lepton and Photon final states
(via H → WW*, ZZ*)

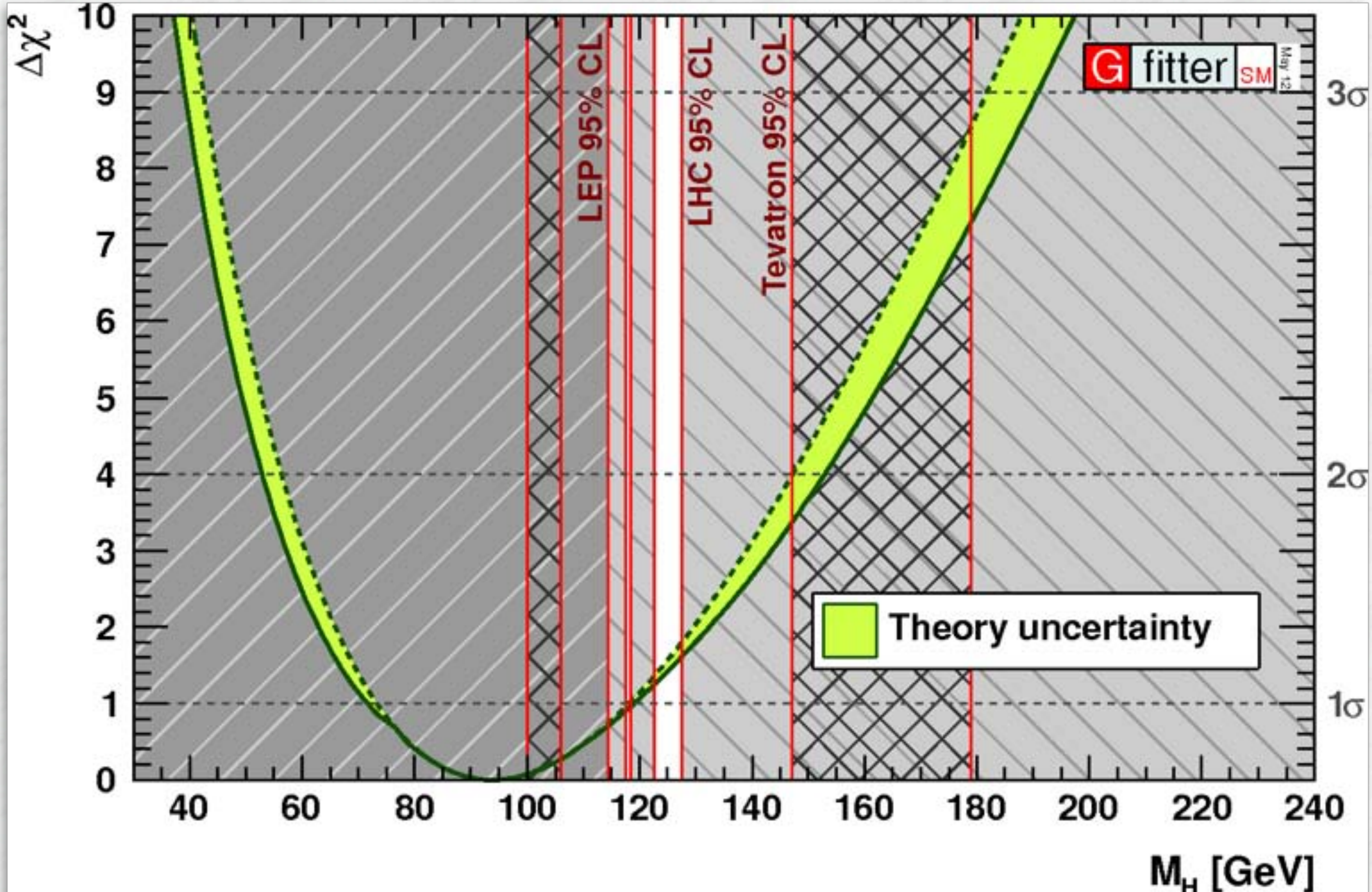
Tau final states

The dominant **bb decay mode** is only useable in the associated production mode (ttH, W/Z H)
(due to the huge QCD jet background, leptons from W/Z or tt decays)

Expected cross sections times production rates at √s = 7 TeV

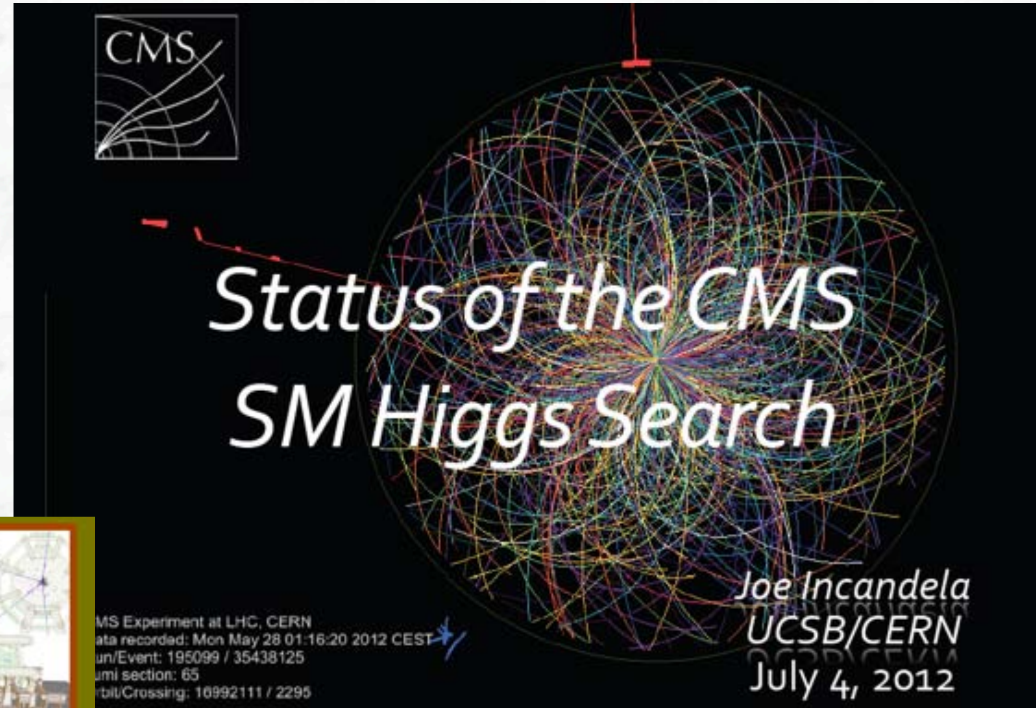


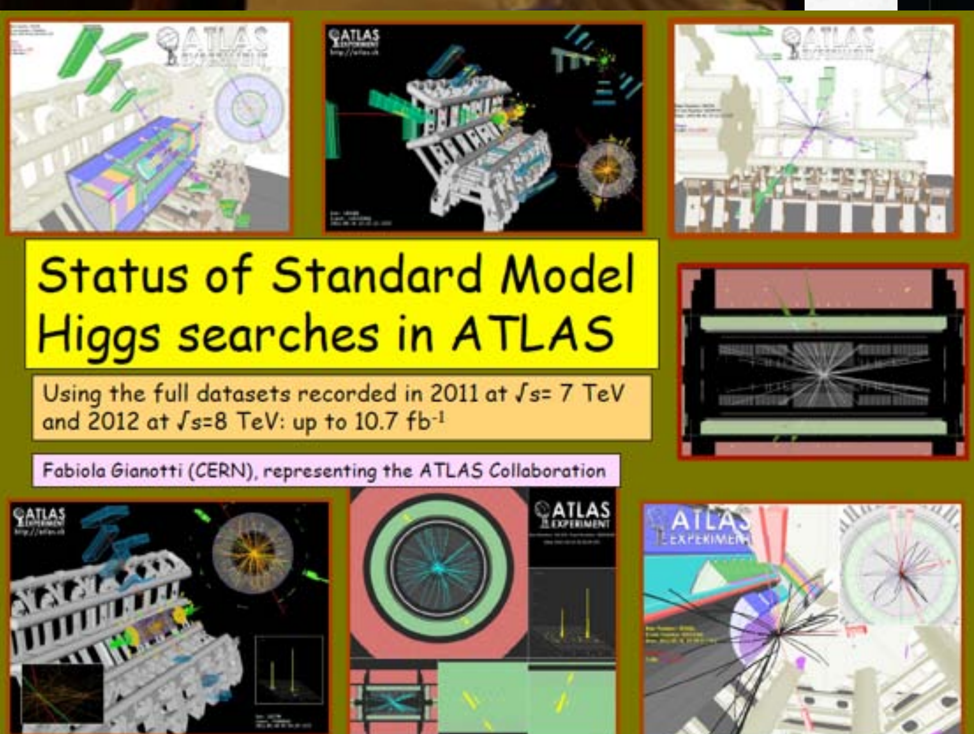
Vor dem 4. Juli 2012



m_H < 152 GeV @ 95% CL (May 2012)

Am 4. Juli 2012



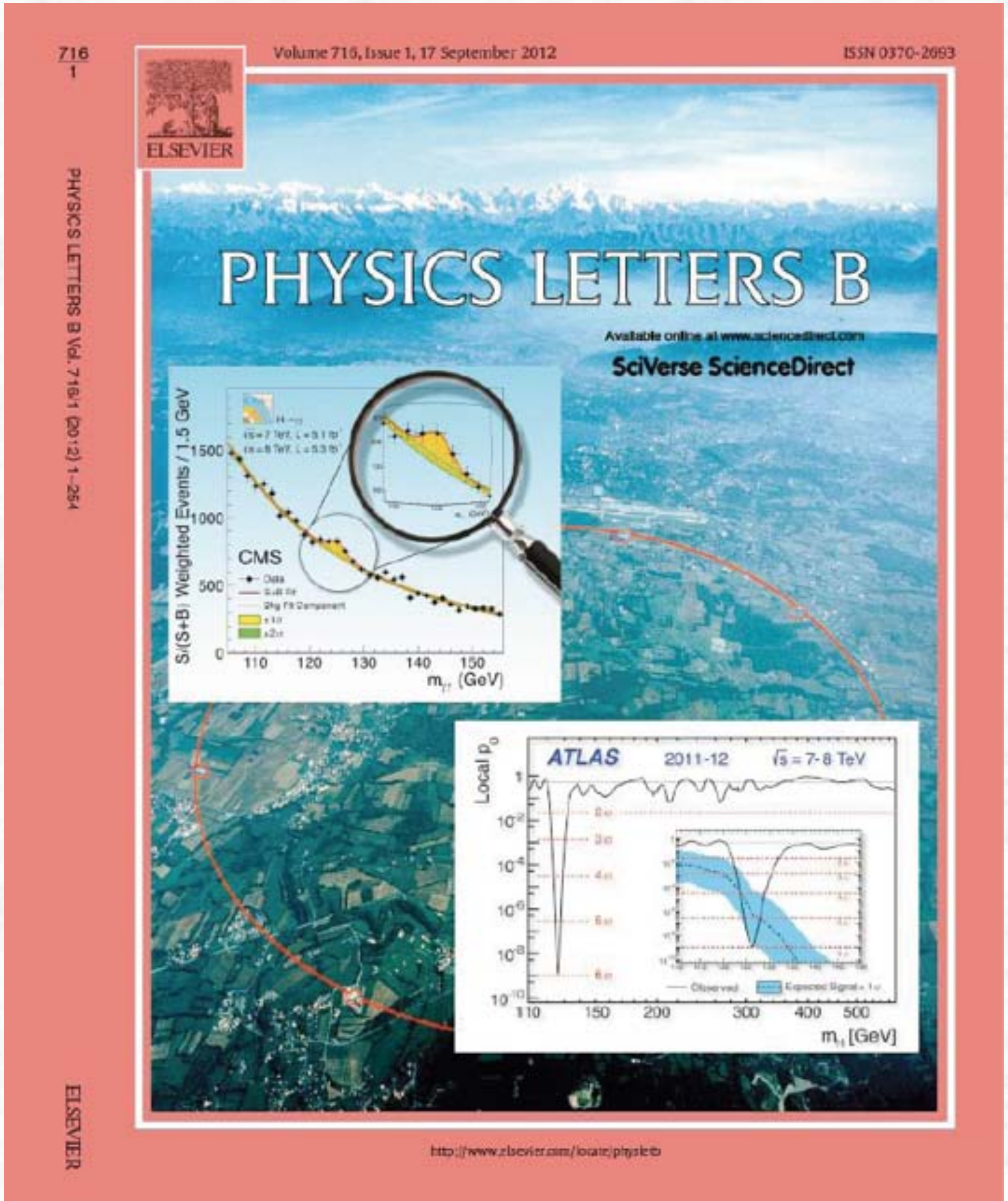


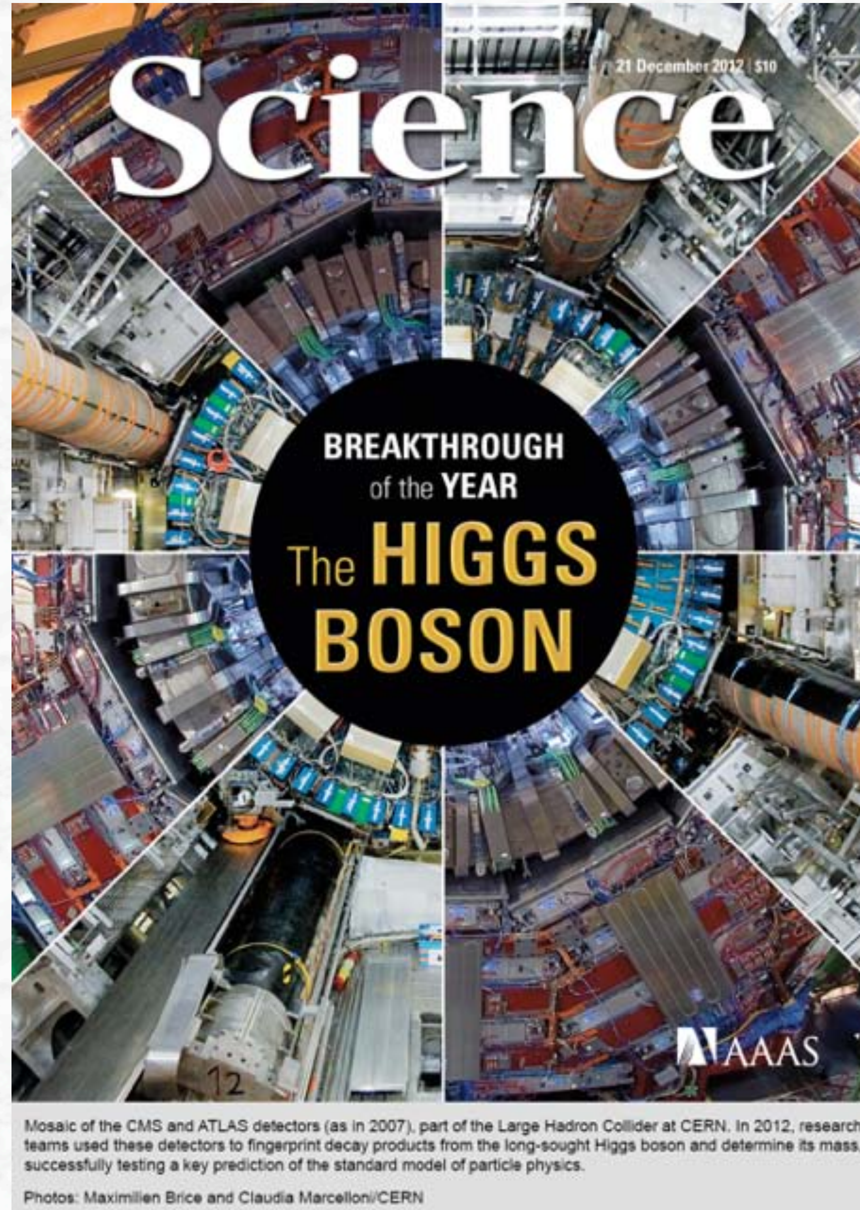
Wikipedia: "In July 2012, when the discovery of the Higgs boson was announced at CERN,[22] the spokesperson of the ATLAS experiment used the font in her presentation of the results."

Am 4. Juli 2012

2013

Der 31. Juli 2012: Die Publikation

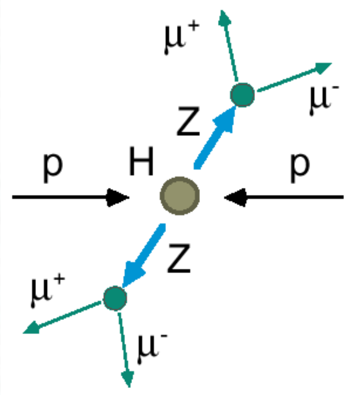




Higgs-boson discovery papers on July 31st 2012, cited over 4400 times each
http://arxiv.org/abs/1207.7214 (ATLAS), http://arxiv.org/abs/1207.7235 (CMS)

$H \rightarrow ZZ^{(*)} \rightarrow \ell\ell\ell\ell$

Signal: σ BR = 22.3 pb * 1.25×10^{-4}
= 2.79 fb
(m_H = 125 GeV, $\sqrt{s}$ = 8 TeV)



P_T (l1) > 20 GeV
P_T (l2) > 15 GeV
P_T (l3) > 10 GeV
P_T (l4) > 6/7 GeV
$|\eta_l| < 2.5$
Isolated leptons
$M(\ell\ell) \sim M_Z$
$M(\ell'\ell') \sim< M_z$

Background: ZZ (irreducible)
Top production (reducible)
$tt \rightarrow Wb\ Wb \rightarrow \ell\nu\ c\ell\nu\ \ell\nu\ c\ell\nu$
σ BR ≈ 1300 fb

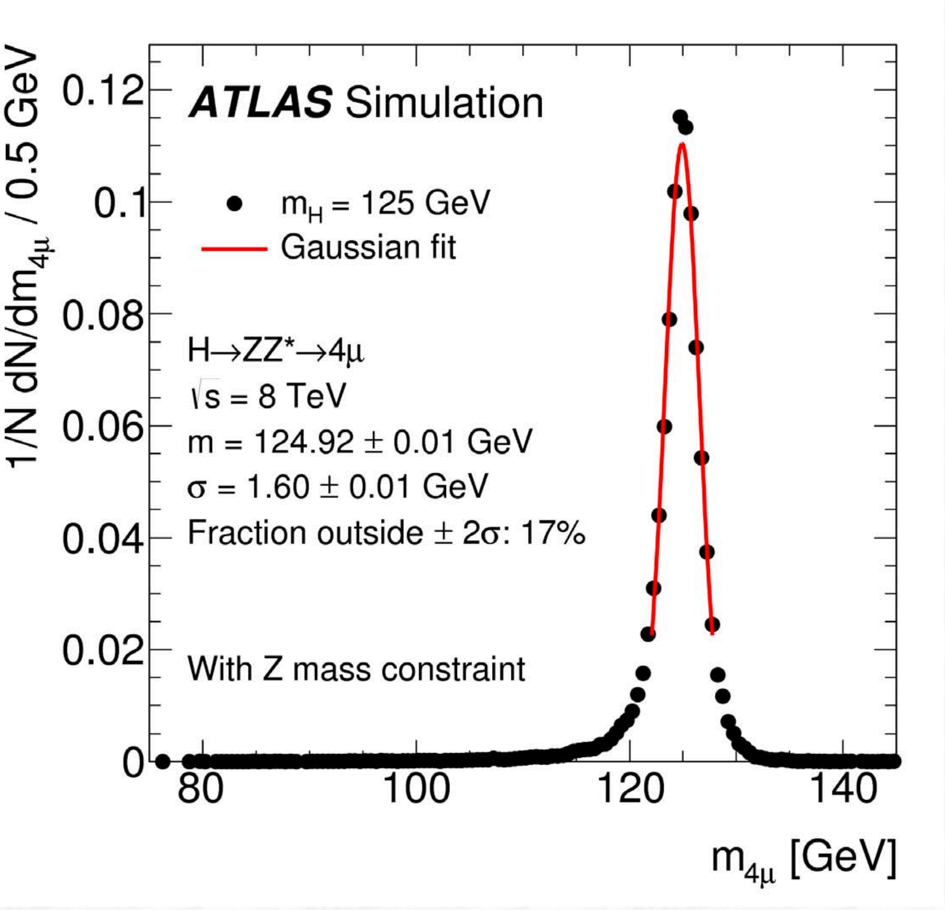
Associated production Z bb (reducible)
$Z\ bb \rightarrow \ell\ell\ c\ell\nu\ c\ell\nu$

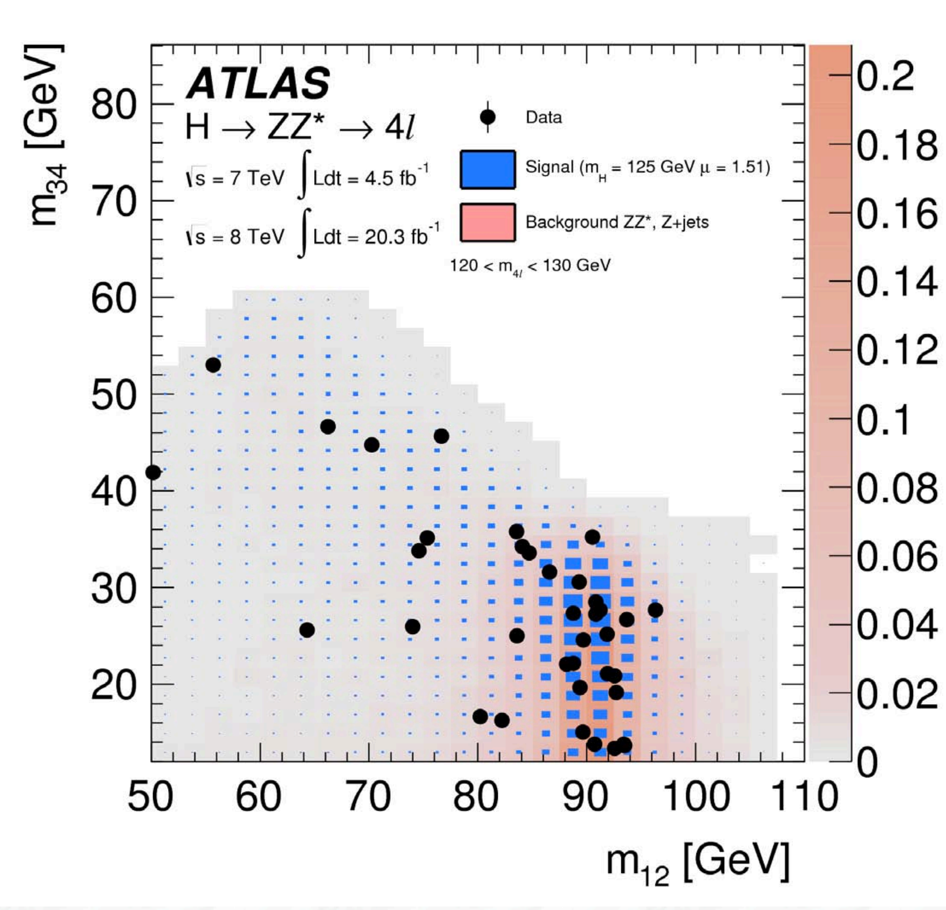
Background rejection: Leptons from b-quark decays
→ non isolated
→ do not originate from primary vertex
(B-meson lifetime: ~ 1.5 ps)

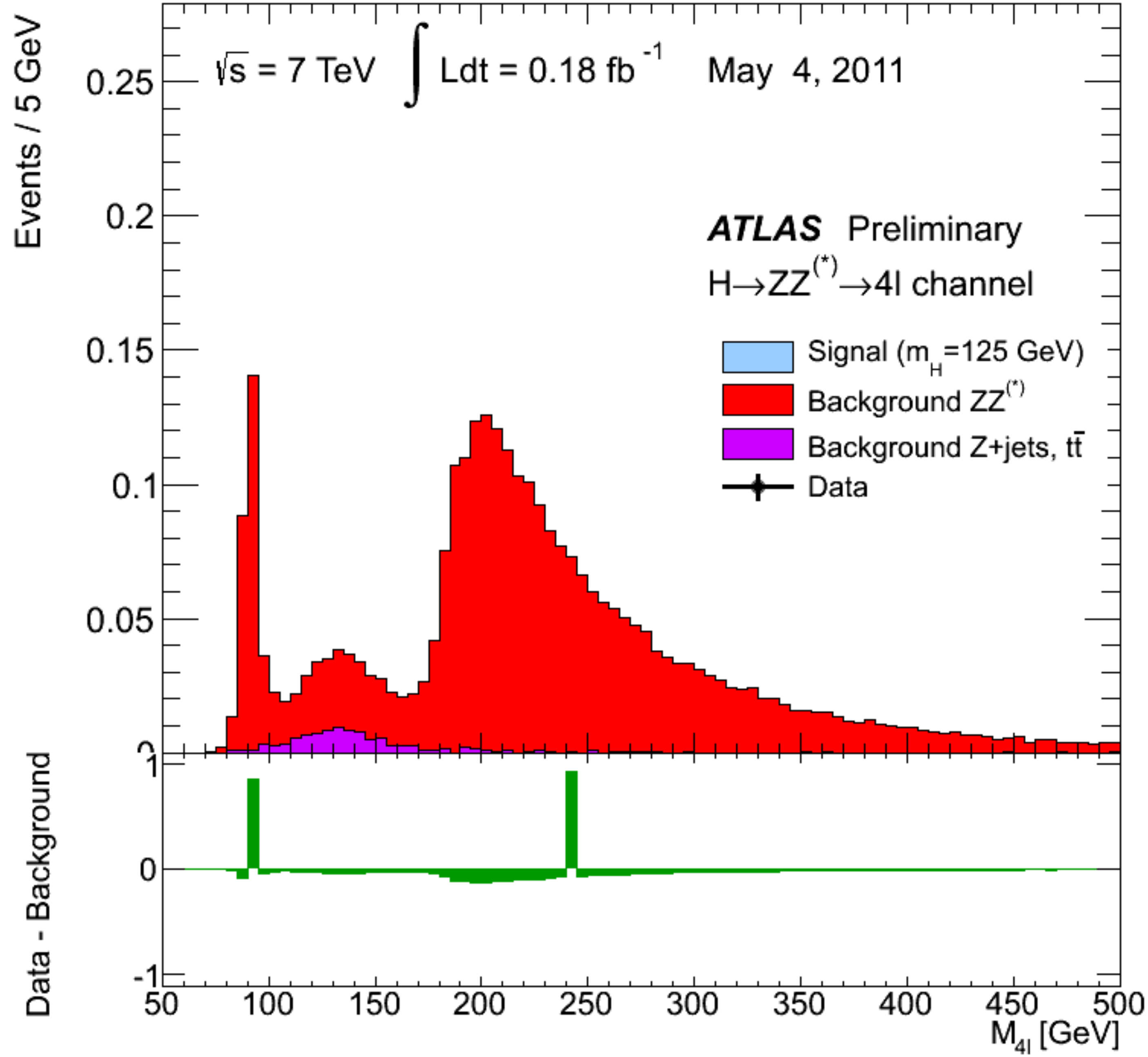
$H \to ZZ^* \to \ell\ell\ \ell\ell$

Excellent m_{4l} invariant-mass resolution

- Excellent lepton momentum resolution
- Photon FSR recovery
- On-shell Z-mass constraint

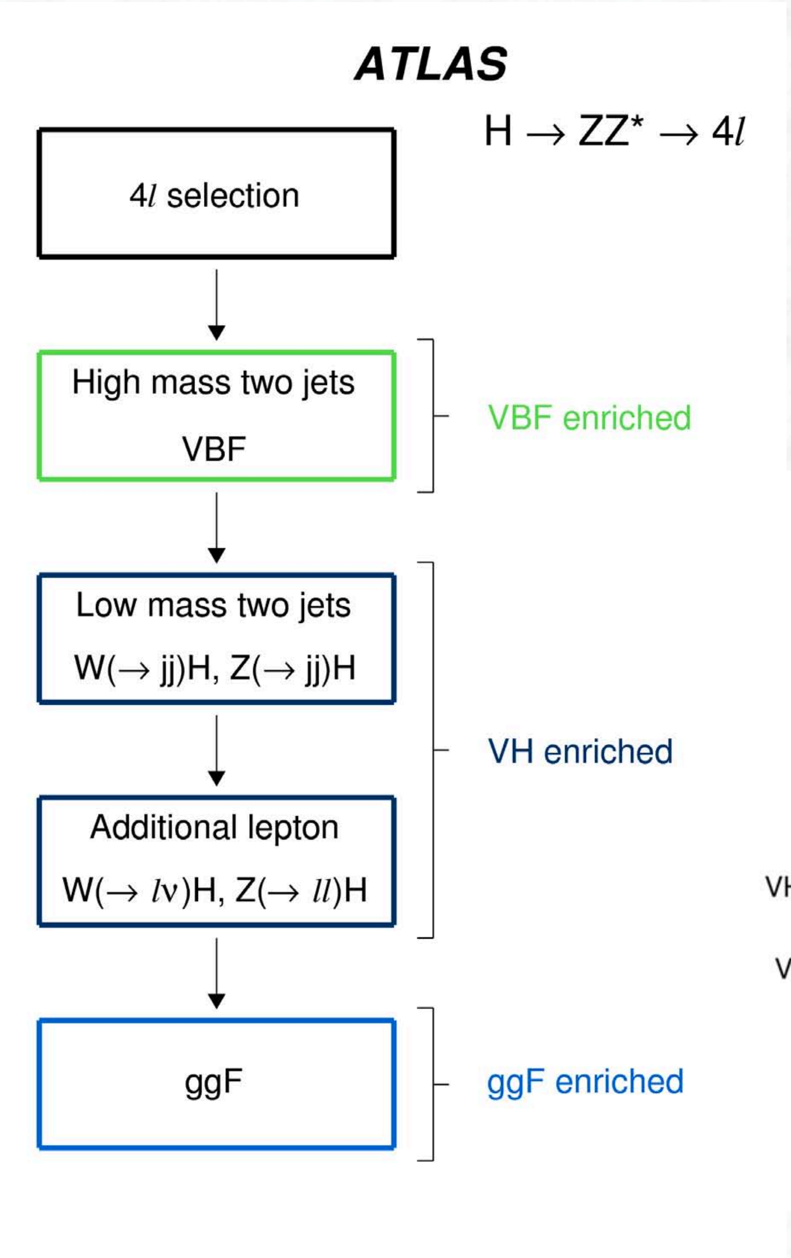






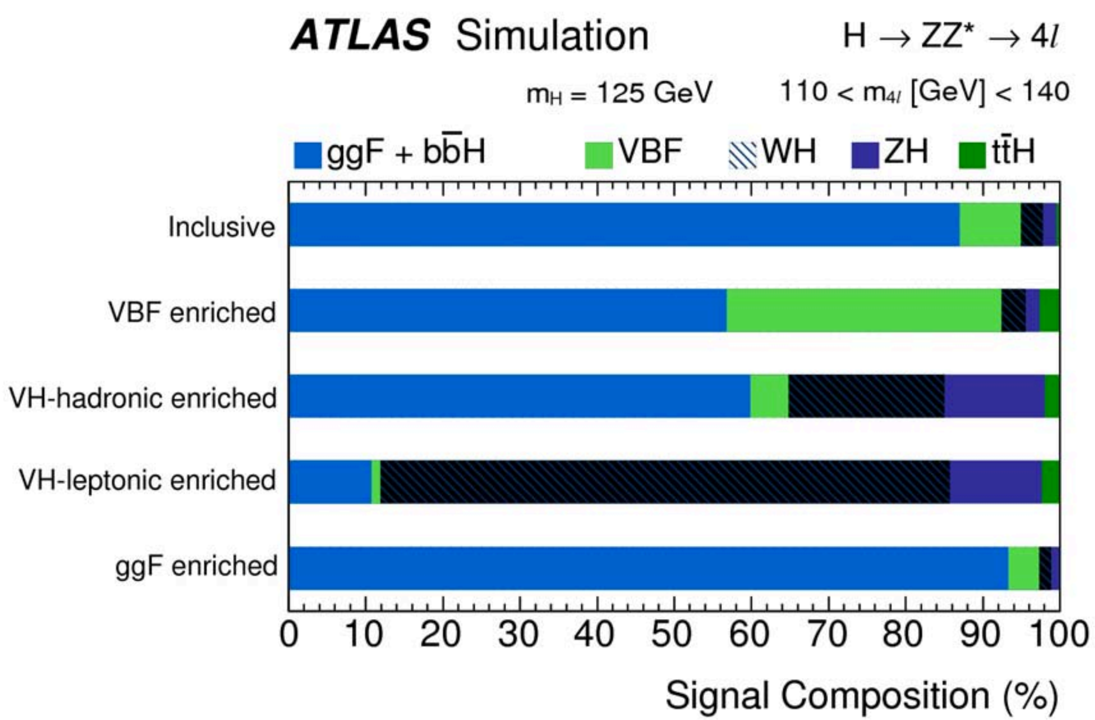
√s = 7 TeV
∫ Ldt = 0.18 fb^-1
May 4, 2011
ATLAS Preliminary
H→ZZ(*)→4l channel
Signal (mH=125 GeV)
Background ZZ(*)
Background Z+jets, tt̄
Data
Events / 5 GeV
0.25
0.2
0.15
0.1
0.05
0
Data - Background
1
0
-1
50
100
150
200
250
300
350
400
450
500
M4l [GeV]

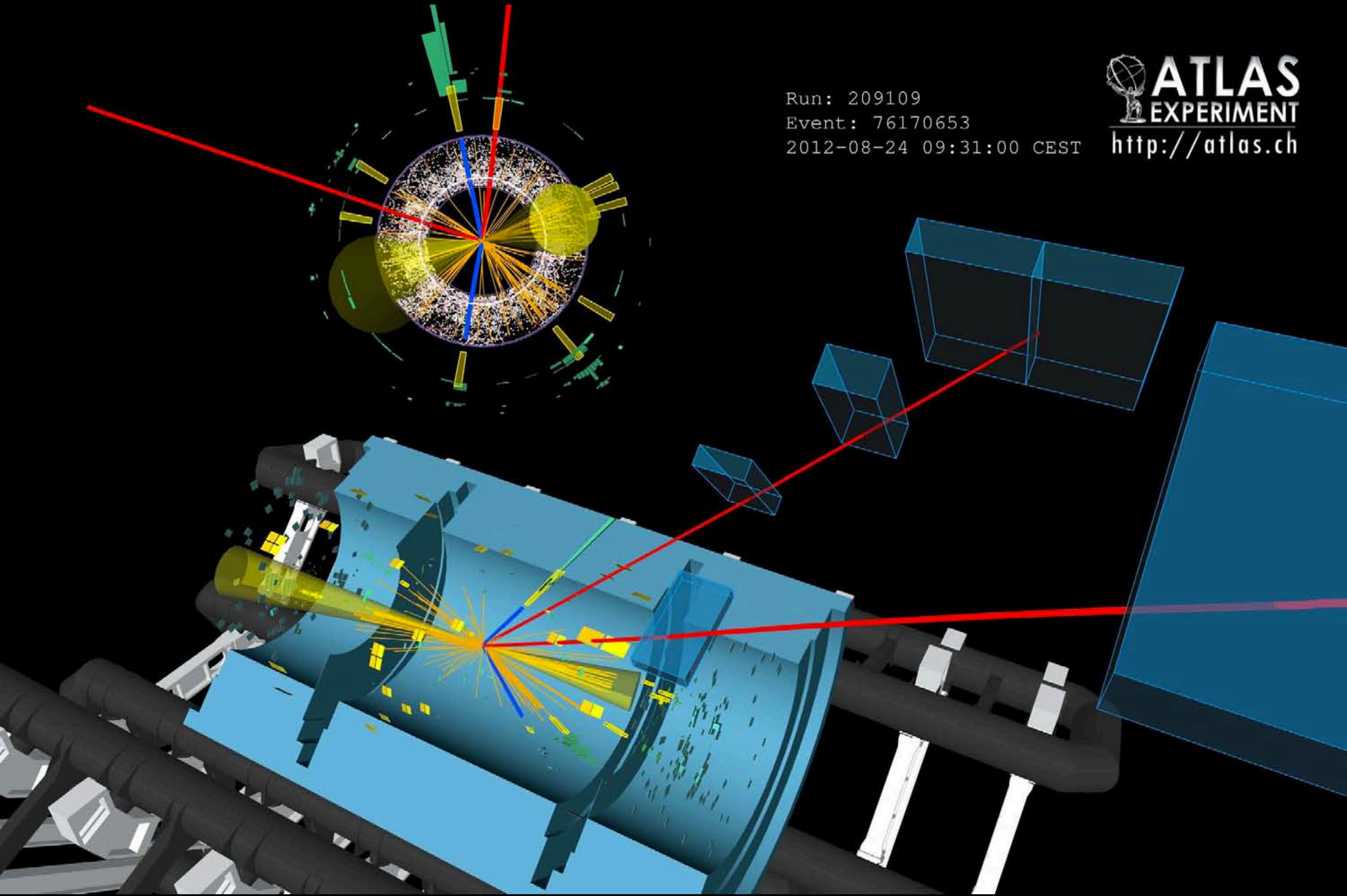
$H \to ZZ^* \to \ell\ell\, \ell\ell$



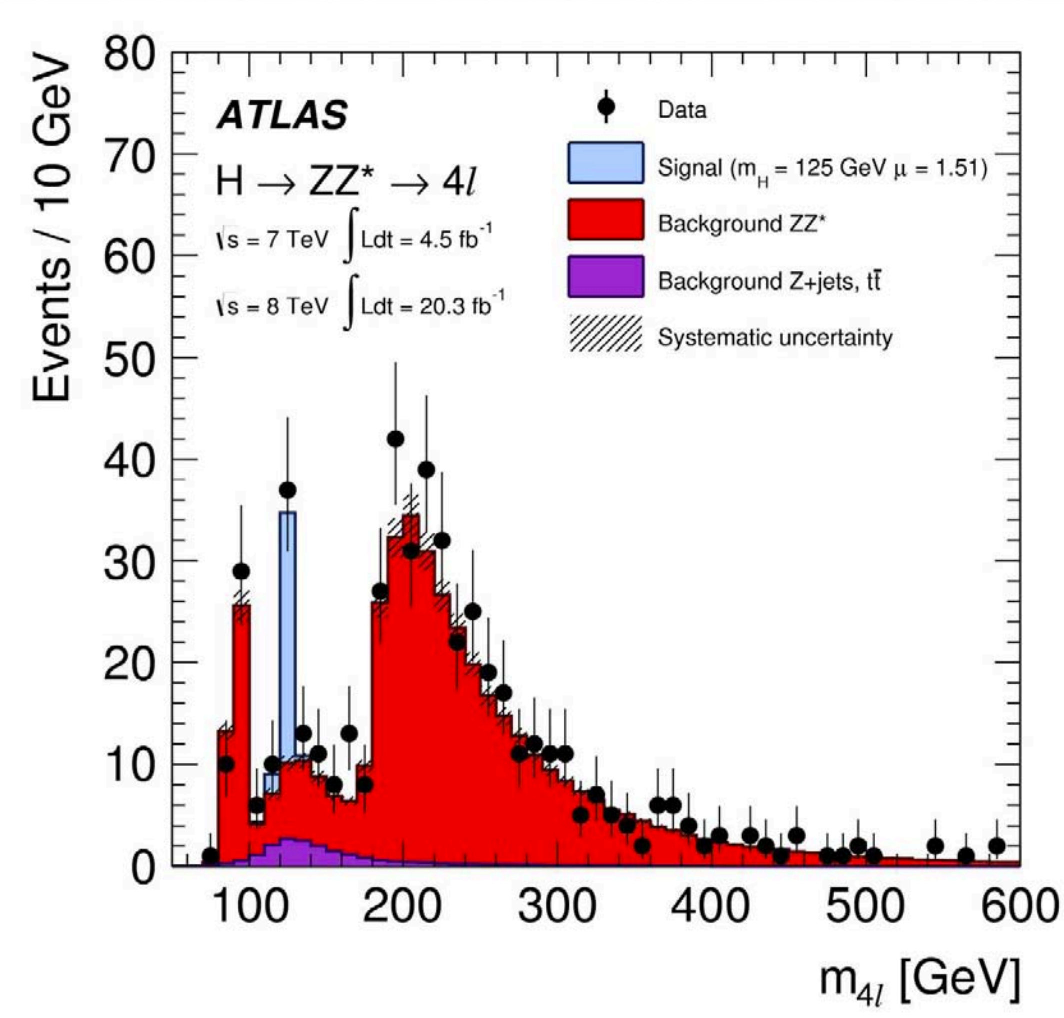
Separate events by production channel

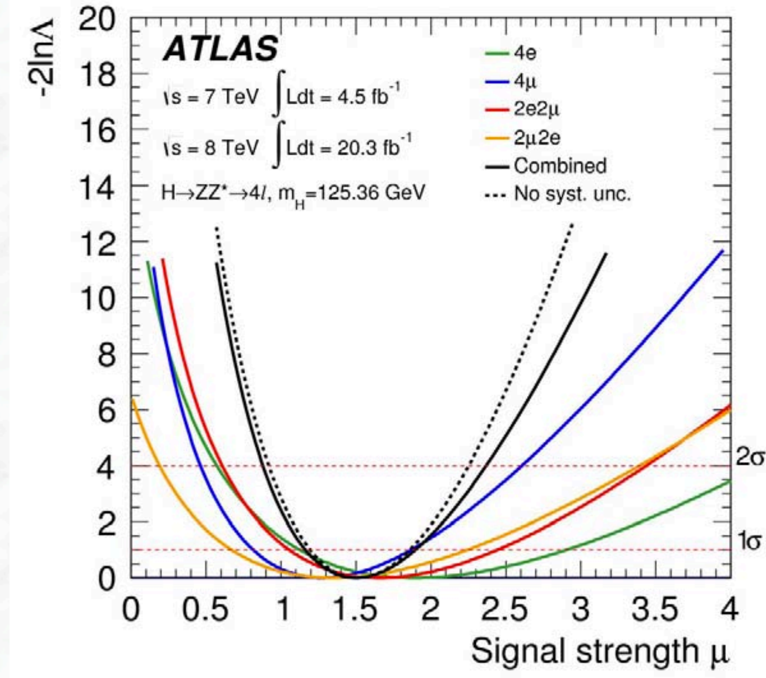
- To improve overall sensitivity
- To allow measurement of
  - individual production modes,
  - and thus individual couplings

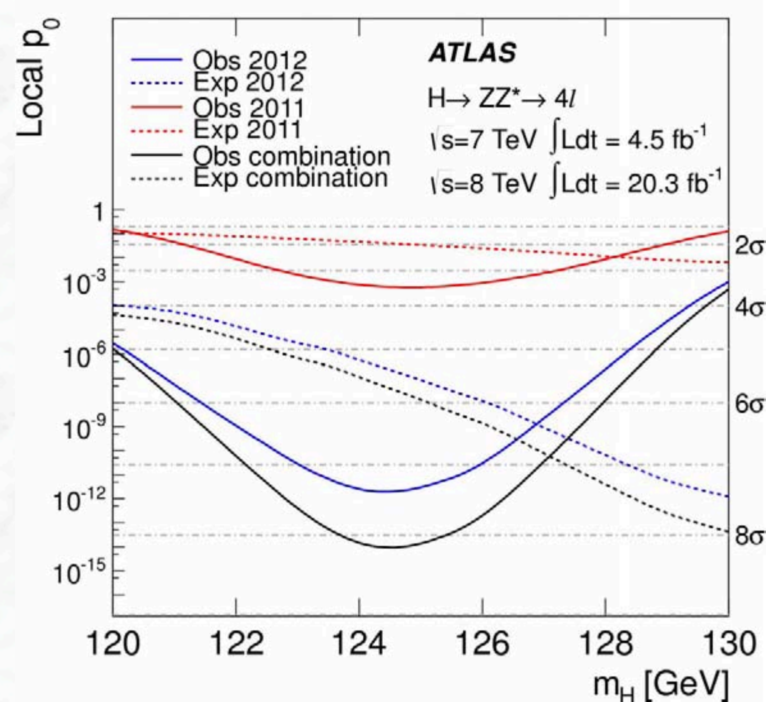



Run: 209109
Event: 76170653
2012-08-24 09:31:00 CEST
ATLAS
EXPERIMENT
http://atlas.ch

H → ZZ* → ℓℓ ℓℓ: Results



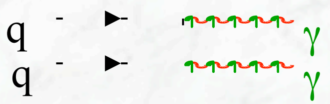




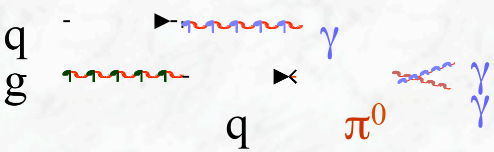
$H \rightarrow \gamma\gamma$

Main backgrounds:

$\gamma\gamma$ irreducible background



γ-jet and jet-jet (reducible)

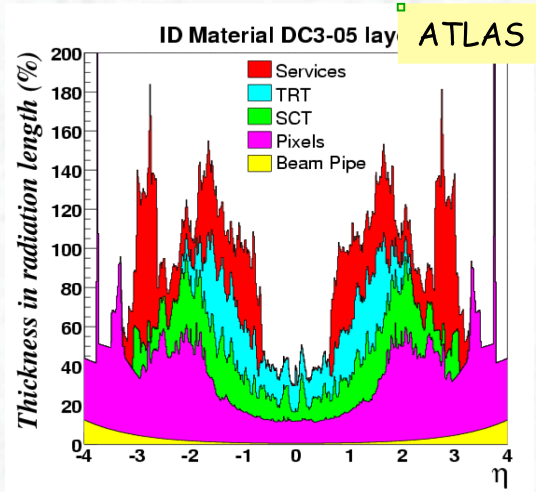


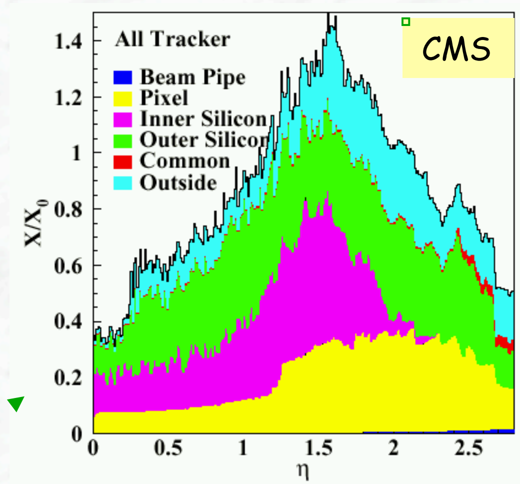
$\sigma_{\gamma j+jj} \sim 10^6\, \sigma_{\gamma\gamma}$ with large uncertainties

→ need $R_j > 10^3$ for $\varepsilon_\gamma \approx 80\%$ to get $\sigma_{\gamma j+jj} \ll \sigma_{\gamma\gamma}$

- Main exp. tools for background suppression:
  - photon identification
  - γ / jet separation (calorimeter + tracker)

  - note: also converted photons need to be reconstructed (large material in LHC silicon trackers)





CMS: fraction of converted γs

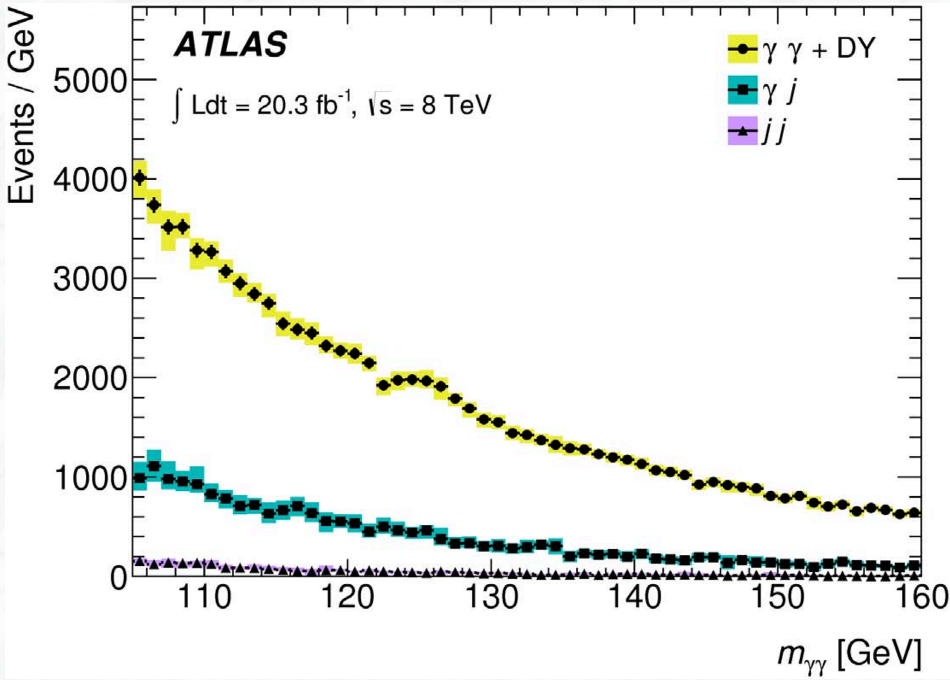
Barrel region: 42.0 %

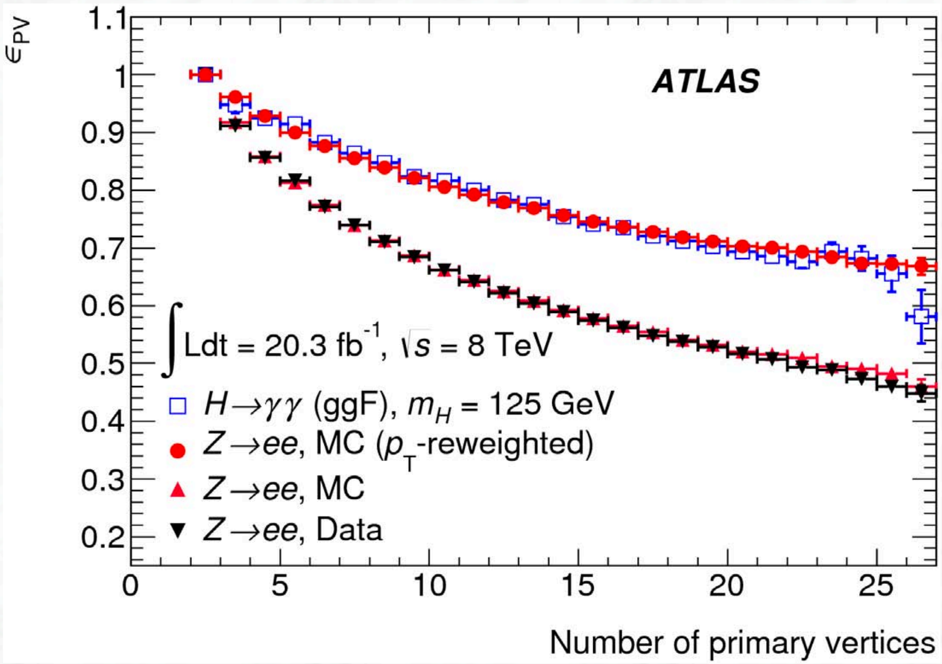
Endcap region: 59.5 %

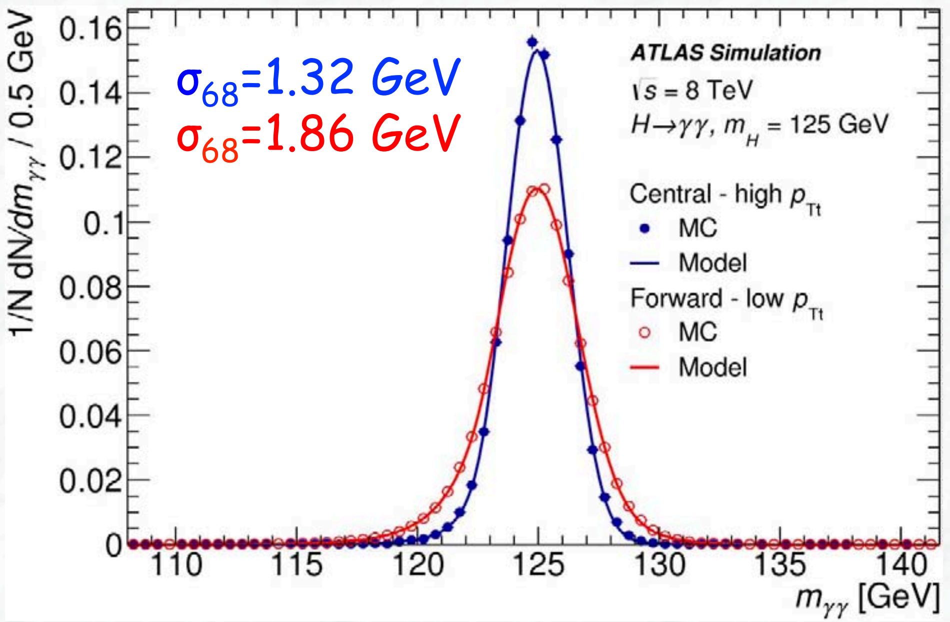
Analysis:

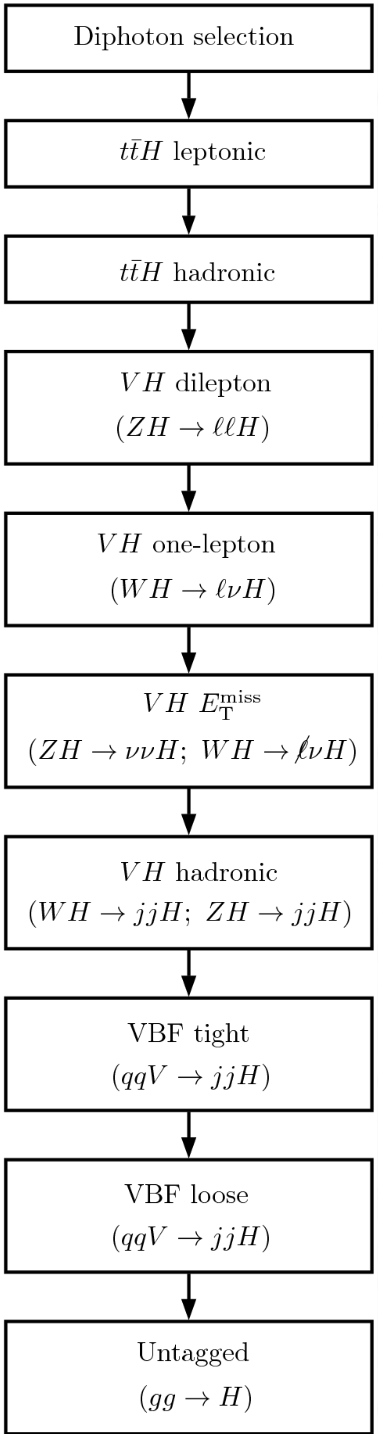
- NLO calculations available
  (Binoth et al., DIPHOX, RESBOS)
- Excellent knowledge of detector material
  - Photon-conversion, hadronic interactions, shower shapes,…
- Split signal sample acc. to $m_{\gamma\gamma}$ resolution
- Improve $m_{\gamma\gamma}$ resolution by neural-network based primary vertex selection

$$m_{\gamma\gamma} = \sqrt{2E_1E_2(1-\cos\,\alpha)}$$



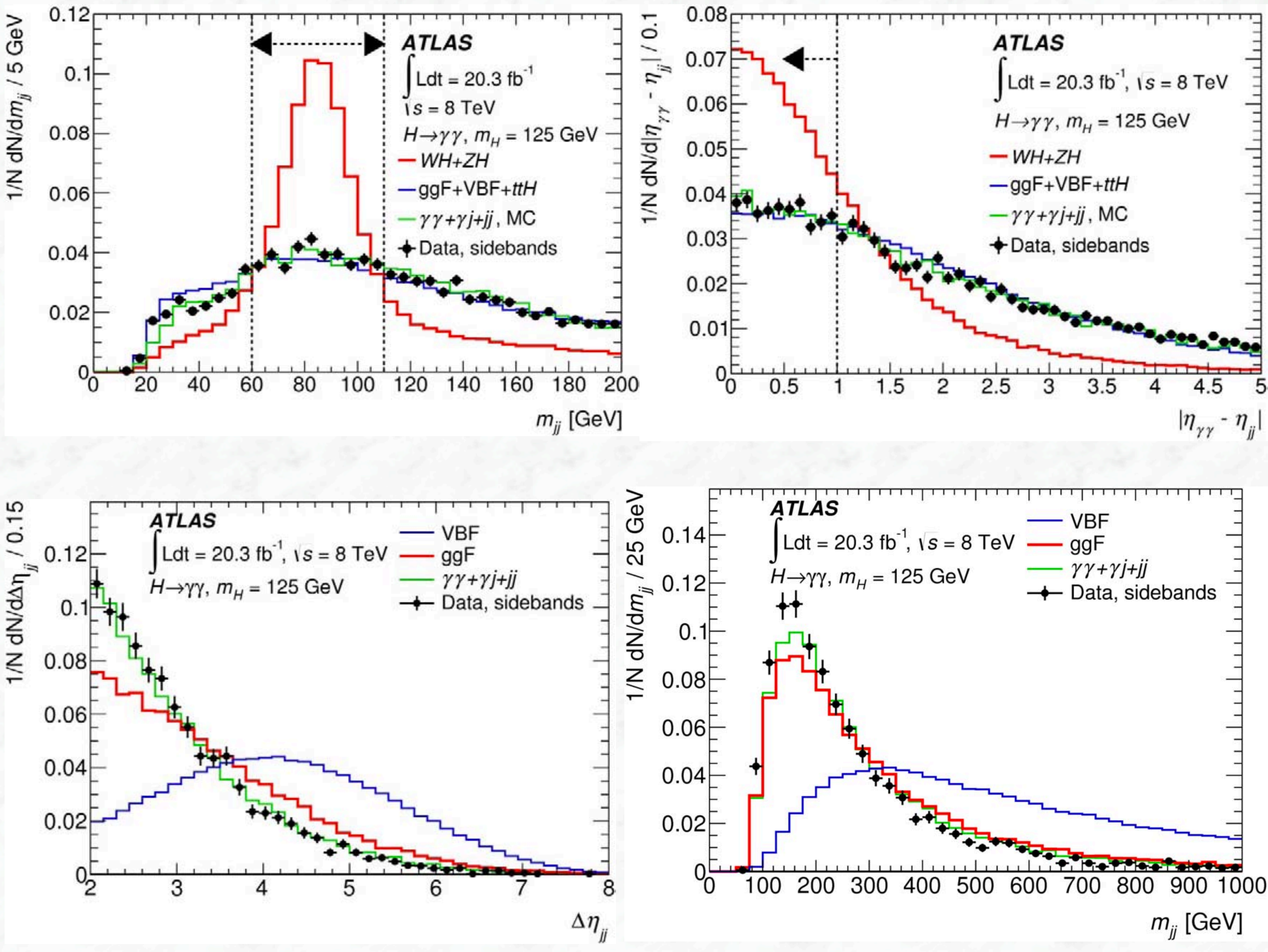






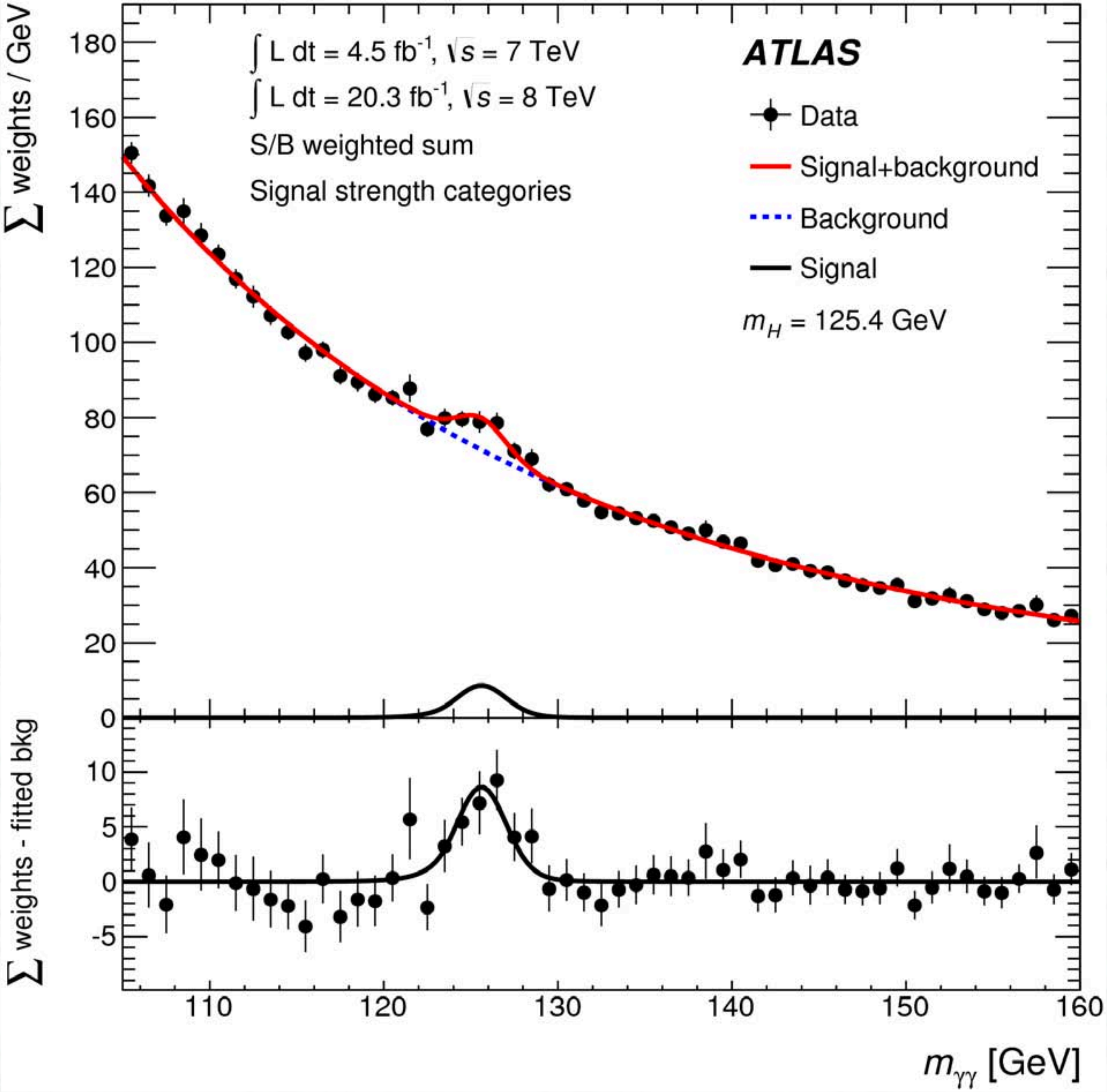
Diphoton selection
$t\bar{t}H$ leptonic
$t\bar{t}H$ hadronic
VH dilepton ($ZH \to \ell\ell H$)
VH one-lepton ($WH \to \ell\nu H$)
VH E_T^{miss} ($ZH \to \nu\nu H$; $WH \to \ell\nu H$)
VH hadronic ($WH \to jjH$; $ZH \to jjH$)
VBF tight ($qqV \to jjH$)
VBF loose ($qqV \to jjH$)
Untagged ($gg \to H$)

H → γγ



ATLAS
∫Ldt = 20.3 fb⁻¹
√s = 8 TeV
H→γγ, m_H = 125 GeV
WH+ZH
ggF+VBF+ttH
γγ+γj+jj, MC
Data, sidebands
1/N dN/dm_jj / 5 GeV
m_jj [GeV]
1/N dN/d|η_γγ - η_jj| / 0.1
|η_γγ - η_jj|
VBF
ggF
γγ+γj+jj
1/N dN/dΔη_jj / 0.15
Δη_jj
1/N dN/dm_jj / 25 GeV
m_jj [GeV]

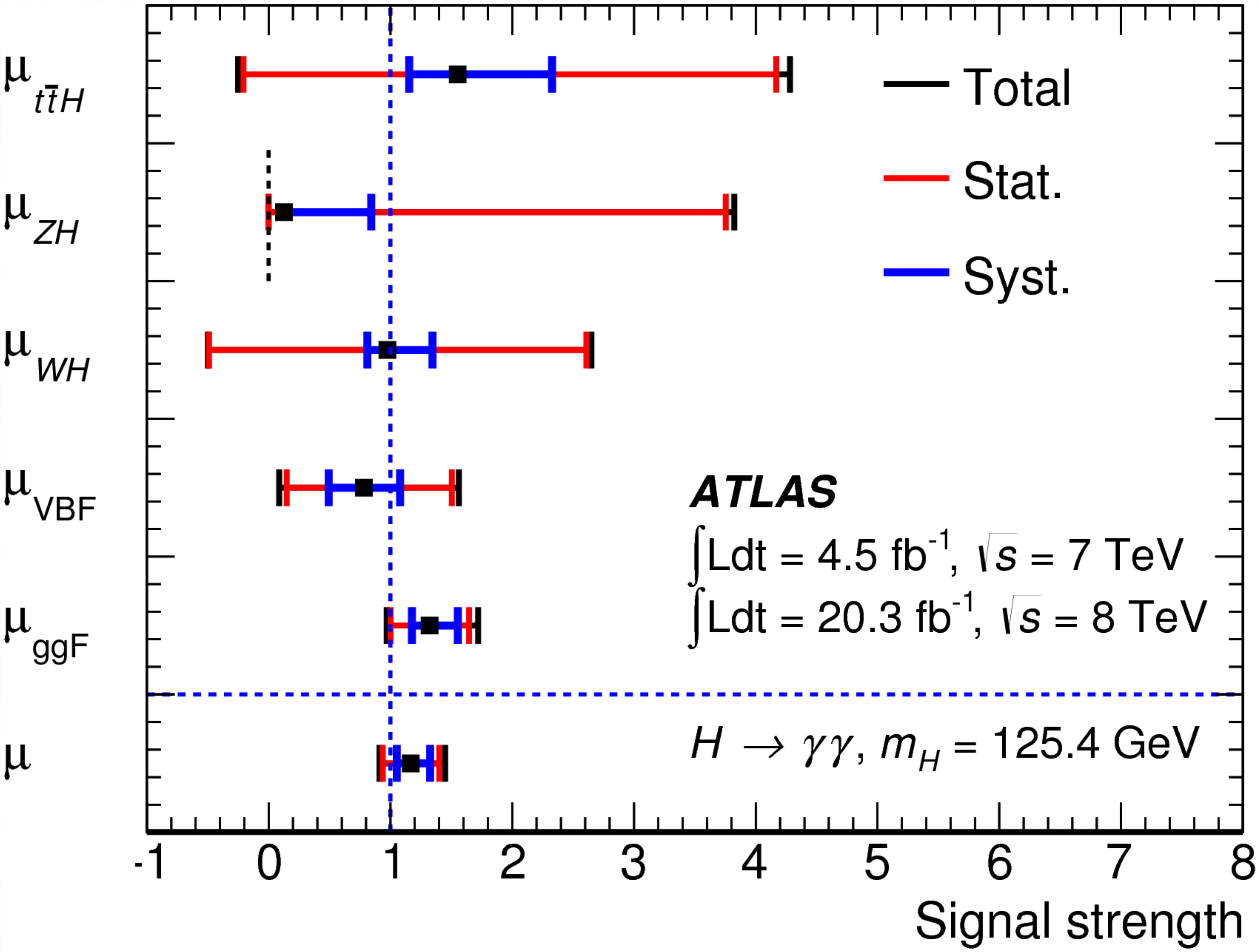
First results on H → γγ searches



Significance of result:

- 5.2σ observed
- 4.6σ expected

H → γγ



Signal strength μ is defined as: $\mu = \frac{\sigma \times \mathrm{BR}}{(\sigma \times \mathrm{BR})_{\mathrm{SM}}}$