

## 6. Physics of W and Z bosons

- 6.1 The theory of electroweak interactions (a short repetition)
- 6.2 Summary of precision tests at LEP
- 6.3 W and Z boson production in hadron colliders
- 6.4 Test of QCD in W/Z (+jet) production
- 6.5 Di-boson measurements
- 6.6 W mass measurement

## Weak Isospin and Hypercharge Quantum

Lepton	$T$	$T^3$	$Q$	$Y$
$\nu_e$	$\frac{1}{2}$	$\frac{1}{2}$	0	-1
$e_L^-$	$\frac{1}{2}$	$-\frac{1}{2}$	-1	-1
$e_R^-$	0	0	-1	-2

## Numbers of Leptons and Quarks

Quark	$T$	$T^3$	$Q$	$Y$
$u_L$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{1}{3}$
$d_L$	$\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{3}$
$u_R$	0	0	$\frac{2}{3}$	$\frac{4}{3}$
$d_R$	0	0	$-\frac{1}{3}$	$-\frac{2}{3}$

# Important Milestones towards Electroweak Unification

1961	S. Glashow proposes an electroweak gauge theory, Introduction of massive $W^\pm$ and $Z^0$ bosons, to explain the large difference in strength of electromagnetic and weak interactions. Key question: how acquire W and Z bosons mass?
1964	R. Brout, F. Englert and P. Higgs demonstrate that mass terms for gauge bosons can be introduced in local gauge invariant theories via spontaneous symmetry breaking
1967	S. Weinberg and A. Salam use Brout-Englert-Higgs mechanism to introduce mass terms for W and Z bosons in Glashow's theory → GSW theory (Glashow, Salam, Weinberg) → mass terms for W, Z bosons, $\gamma$ remains massless → Higgs particle (see chapter 7)
1973	G. t'Hooft and M. Veltman show that GSW theory is renormalizable
1979	Nobel price for S. Glashow, A. Salam and S. Weinberg
1983	Experimental discovery of the W and Z bosons by UA1 and UA2 experiments at the CERN SpbarS collider ( $\sqrt{s} = 540$ GeV)
1990-2000	Precise test of the electroweak theory at LEP
1999	Nobel price for G. t'Hooft and M. Veltman
2012	Discovery of a Higgs particle by the ATLAS and CMS experiments at the LHC

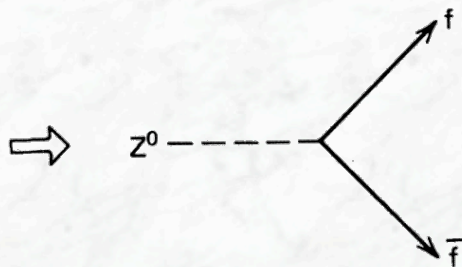
# W and Z vertex factors

$$\left. \begin{aligned} & -i \frac{g}{\sqrt{2}} (\bar{\chi}_L \gamma^\mu \tau_+ \chi_L) W_\mu^+ \\ & = -i \frac{g}{\sqrt{2}} (\bar{\nu}_L \gamma^\mu e_L) W_\mu^+ \end{aligned} \right\} W^+ \rightarrow$$

$$\boxed{-i \frac{g}{\sqrt{2}} \gamma^\mu \frac{1}{2} (1 - \gamma^5)}$$

$$\left. \begin{aligned} & -i \frac{g}{\sqrt{2}} (\bar{\chi}_L \gamma^\mu \tau_- \chi_L) W_\mu^- \\ & = -i \frac{g}{\sqrt{2}} (\bar{e}_L \gamma^\mu \nu_L) W_\mu^- \end{aligned} \right\} W^- \rightarrow$$

$$\boxed{-i \frac{g}{\cos \theta_W} \gamma^\mu \frac{1}{2} (c_V^f - c_A^f \gamma^5)}$$



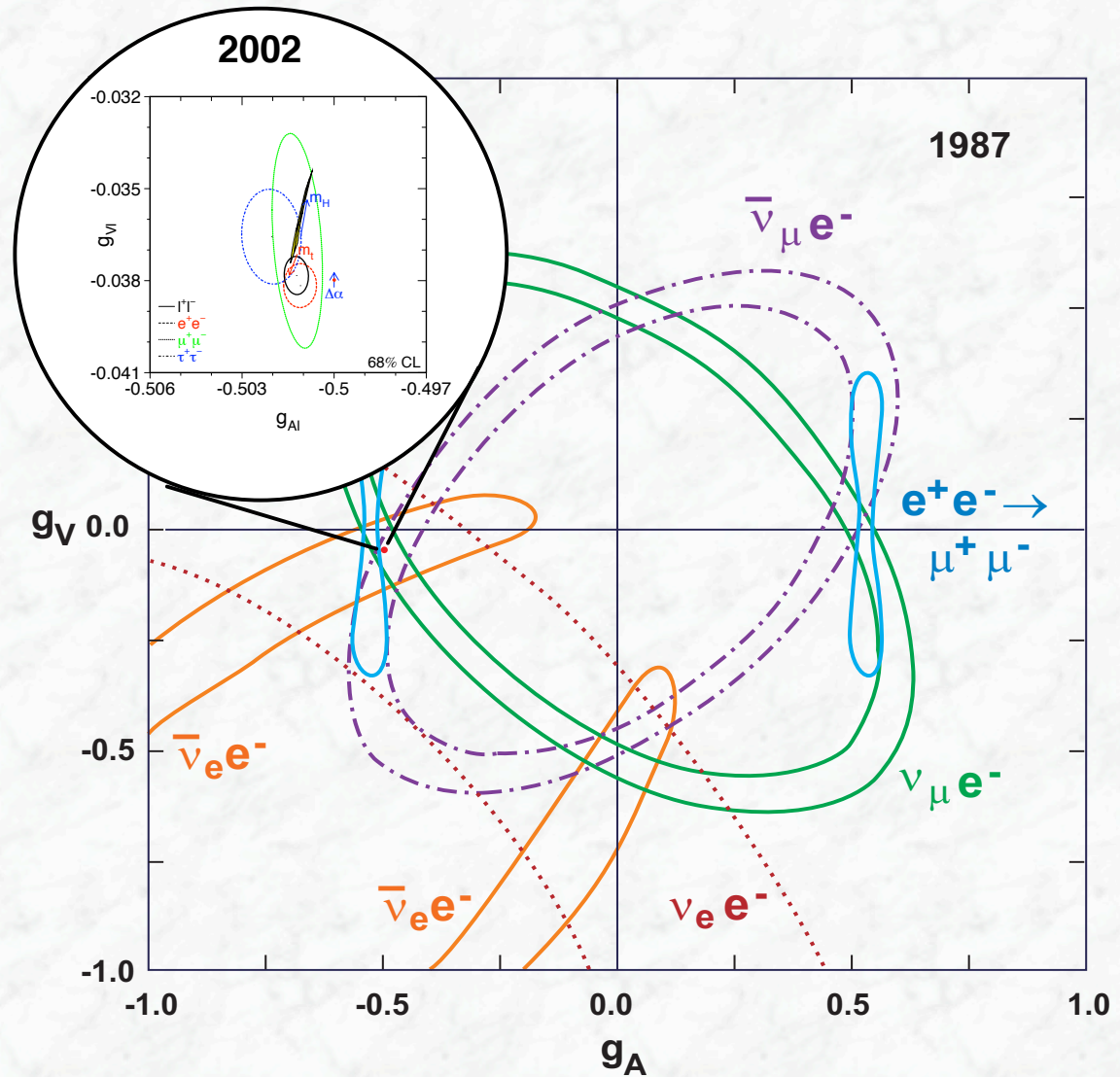
The  $Z \rightarrow ff$  vertex factors in the Standard Model  
 ( $\sin^2 \theta_W$  is assumed to be 0.234)

$f$	$Q_f$	$c_A^f$	$c_V^f$
$\nu_e, \nu_\mu, \dots$	0	$\frac{1}{2}$	$\frac{1}{2}$
$e^-, \mu^-, \dots$	-1	$-\frac{1}{2}$	$-\frac{1}{2} + 2 \sin^2 \theta_W \approx -0.03$
$u, c, \dots$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{2} - \frac{4}{3} \sin^2 \theta_W \approx 0.19$
$d, s, \dots$	$-\frac{1}{3}$	$-\frac{1}{2}$	$-\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W \approx -0.34$

[Halzen&Martin]

## 6.2 Summary of electroweak precision tests at LEP

- Results of 30 years of experimental and theoretical progress
- The electroweak theory is tested at the level of  $10^{-4}$



# LEP am CERN / Genf

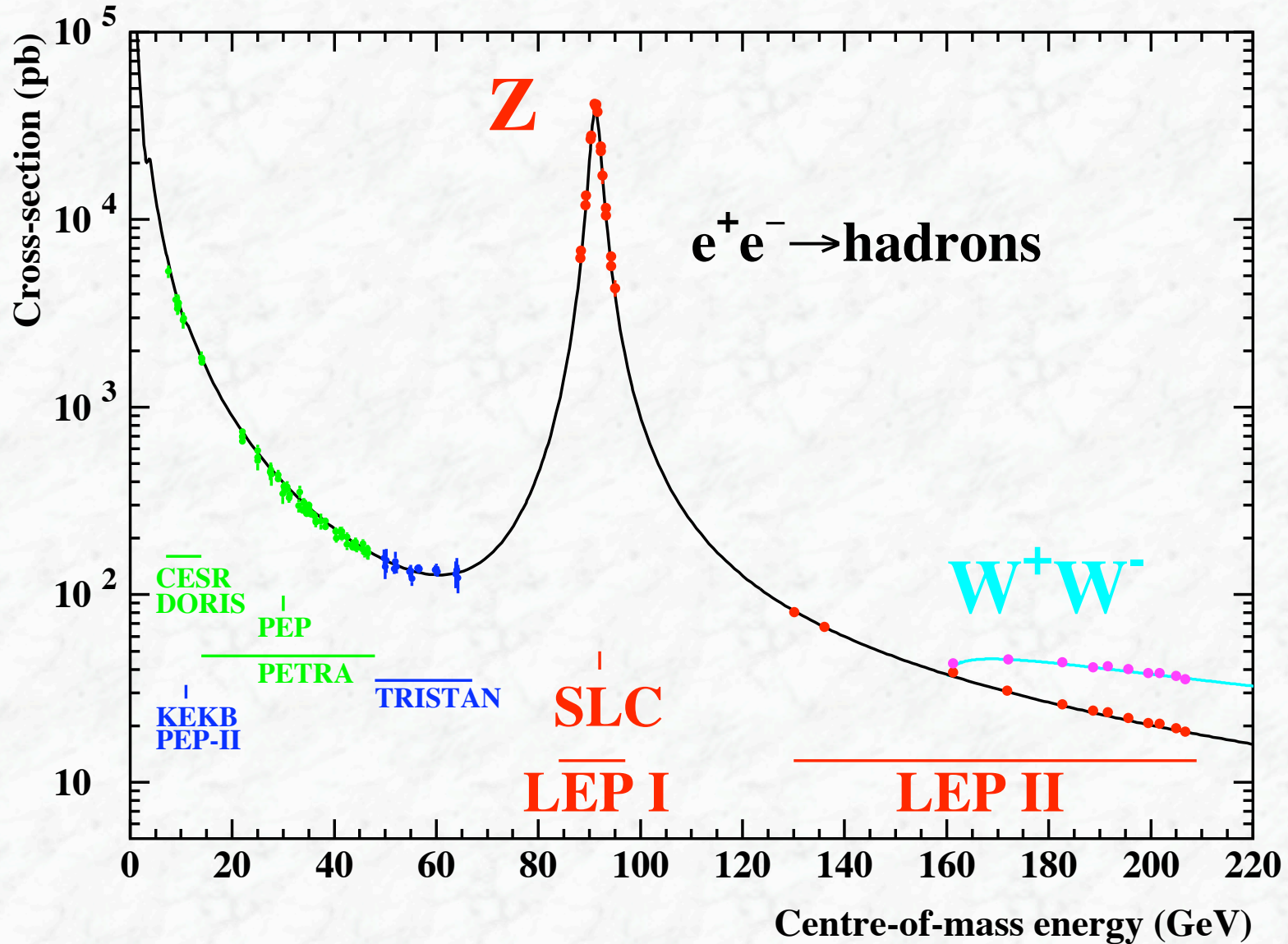


$e^+e^-$ -Beschleuniger, 27 km Umfang

Schwerpunktenergie: LEP-I (1989-1995) 91 GeV  
LEP-II (1996-2000) → 208 GeV

4 Experimente: ALEPH, DELPHI, L3, OPAL

# Cross sections for W and Z boson production



Precision tests  
of the Z sector

Tests of the  
W sector



# Cross section for $e^+e^- \rightarrow \mu^+\mu^-$ at LEP I

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left[ F_\gamma(\cos\theta) + F_{\gamma Z}(\cos\theta) \frac{s(s-M_Z^2)}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} + F_Z(\cos\theta) \frac{s^2}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} \right]$$

$\gamma$

$\gamma/Z$  interference

$Z$

vanishes at  $\sqrt{s} \approx M_Z$

$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^e)^2 + (g_A^e)^2] (g_V^\mu)^2 + (g_A^\mu)^2 (1 + \cos^2\theta) + 8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

$\alpha = \alpha(m_Z)$ : running el. magnetic coupling [ $\alpha(M_Z) = \alpha / (1 - \Delta\alpha)$  mit  $\Delta\alpha \approx 0.06$ ]

$g_V, g_A = c_V, c_A$ : effective coupling constants (vector and axial vector)

# Cross section for $e^+e^- \rightarrow ff$ at LEP I

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left[ F_\gamma(\cos\theta) + F_{\gamma Z}(\cos\theta) \frac{s(s-M_Z^2)}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} + F_Z(\cos\theta) \frac{s^2}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} \right]$$

$\gamma$

$\gamma/Z$  interference

$Z$

vanishes at  $\sqrt{s} \approx M_Z$

$\times N_C^f$   
number of colour degrees of freedom for fermion f

$\times (1 + \delta_{\text{QCD}})$   
QCD correction term

$$F_\gamma(\cos\theta) = Q_e^2 Q_f^2 (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_f}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^f \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{f^2} + g_A^{f^2})(1 + \cos^2\theta) + 8g_V^e g_A^e g_V^f g_A^f \cos\theta]$$

# Cross section for $e^+e^- \rightarrow ff$ on resonance ( $\sqrt{s} = m_Z$ )

- On resonance,  $\sqrt{s} = m_Z$ :
  - $\gamma^*/Z$  interference terms vanishes
  - $\gamma$  term contributes  $\sim 1\%$
  - **Z contribution dominates !**
- Contribution of the  $\gamma^*/Z$  interference term at  $s = (M_Z - 3 \text{ GeV})^2$  :  $\sim 0.2\%$

Total cross section for  $e^+e^- \rightarrow \mu^+\mu^-$  (integration over  $\cos \theta$ )

$$\sigma_{\text{tot}} \approx \sigma_Z = \frac{4\pi}{3s} \frac{\alpha^2}{16 \sin^4 \theta_W \cos^4 \theta_W} \cdot [(g_V^e)^2 + (g_A^e)^2][(g_V^\mu)^2 + (g_A^\mu)^2] \cdot \frac{s^2}{(s - M_Z^2)^2 + (M_Z \Gamma_Z)^2}$$

$$\sigma_Z(\sqrt{s} = M_Z) = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2} \quad \text{Peak cross section}$$

$$\Gamma_f = \frac{\alpha M_Z}{12 \sin^2 \theta_W \cos^2 \theta_W} \cdot [(g_V^f)^2 + (g_A^f)^2]$$

Partial width

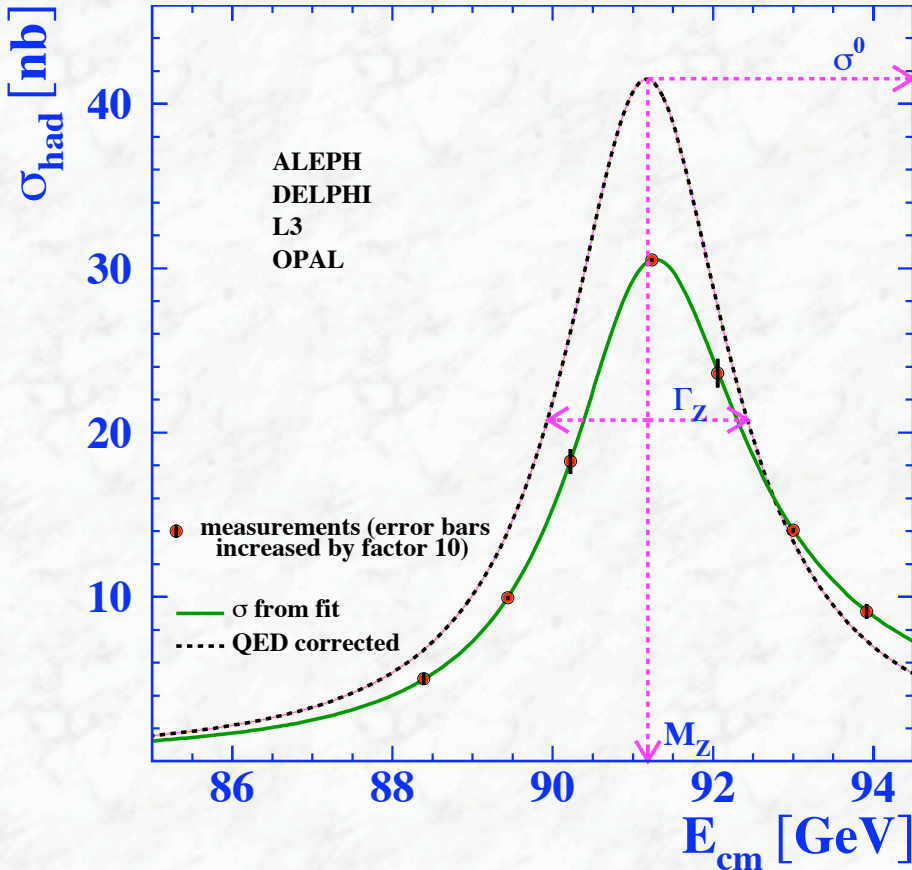
$$\Gamma_Z = \sum_i \Gamma_i \quad \text{Total width}$$

From the energy dependence of the total cross section (for various fermions f) the parameters

$M_Z, \Gamma_Z, \Gamma_f$

can be determined.

# Measurement of the Z line-shape



Line shape (resonance curve):

$$\sigma(s) = 12\pi \frac{\Gamma_e \Gamma_\mu}{M_Z^2} \cdot \frac{s}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

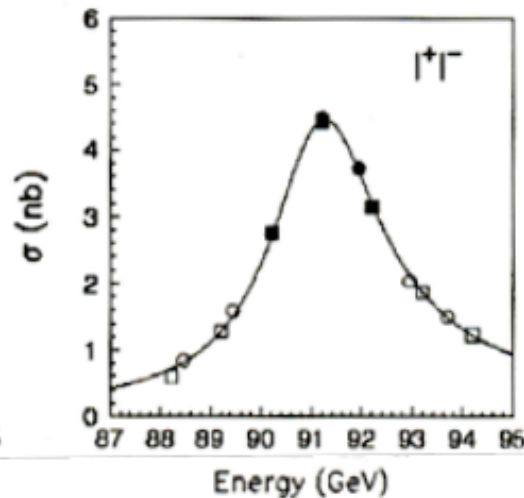
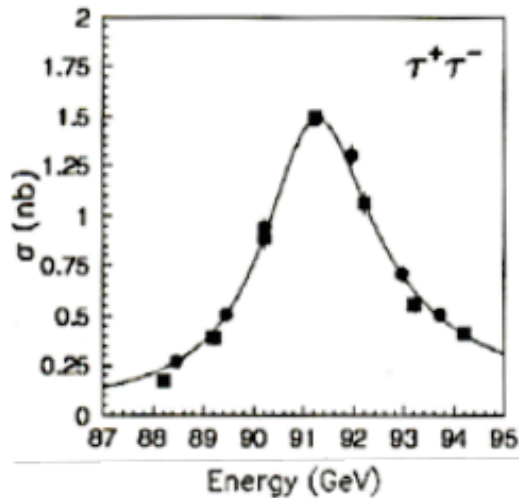
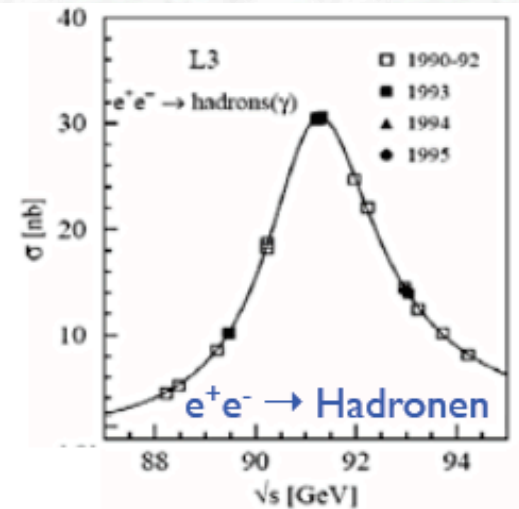
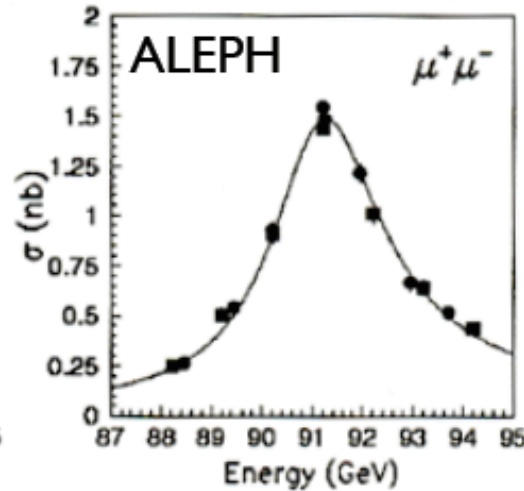
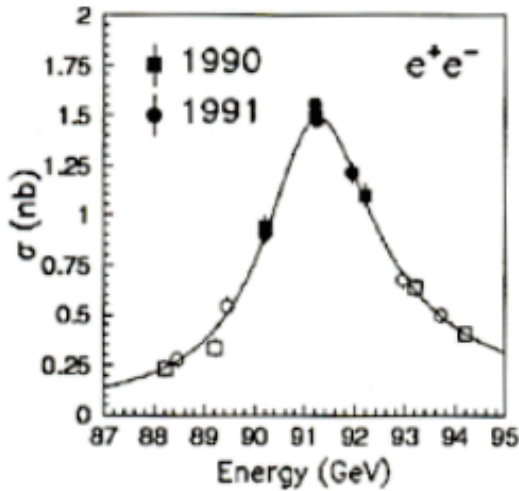
Peak: 
$$\sigma_0 = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$$

- Position of maximum  $\rightarrow M_Z$
- Full width at half maximum  $\rightarrow \Gamma_Z$
- Peak cross section  $\sigma_0 \rightarrow \Gamma_e \Gamma_\mu$

Radiative corrections (photon radiation)  
important

- with ISR (initial state radiation)
- - - without ISR

# Measurement of the Z line-shape (cont.)



Quark-Flavor i.a. nicht exp. trennbar  
(Ausnahme:  $c, b \rightarrow$  Lebensdauer)  
 $\Rightarrow$  had. Breite:  $\Gamma_{\text{had}} = \Gamma_u + \Gamma_d + \Gamma_s + \Gamma_c + \Gamma_b$

Messe Verhältnisse der Pol-WQ:

$$R_l^0 \equiv \frac{\Gamma_{\text{had}}}{\Gamma_{ll}} \quad l = e, \mu, \tau$$

$$R_q^0 \equiv \frac{\Gamma_{qq}}{\Gamma_{\text{had}}} \quad q = b, c$$

- Keine Unterschiede für verschiedene Leptonarten  $\Rightarrow$  **Leptonuniversalität**
- Form der Resonanzenkurve für alle Endzustände gleich (gleicher Propagator!)

# Results on Z line-shape parameters

$$M_Z = 91.1876 \pm 0.0021 \text{ GeV} \quad 23 \text{ ppm} (*)$$

$$\begin{aligned}\Gamma_Z &= 2.4952 \pm 0.0023 \text{ GeV} \\ \Gamma_{\text{had}} &= 1.7458 \pm 0.0027 \text{ GeV} \\ \Gamma_e &= 0.08392 \pm 0.00012 \text{ GeV} \\ \Gamma_\mu &= 0.08399 \pm 0.00018 \text{ GeV} \\ \Gamma_\tau &= 0.08408 \pm 0.00022 \text{ GeV}\end{aligned}$$

3 lepton flavours  
treated independently

Test of lepton  
universality

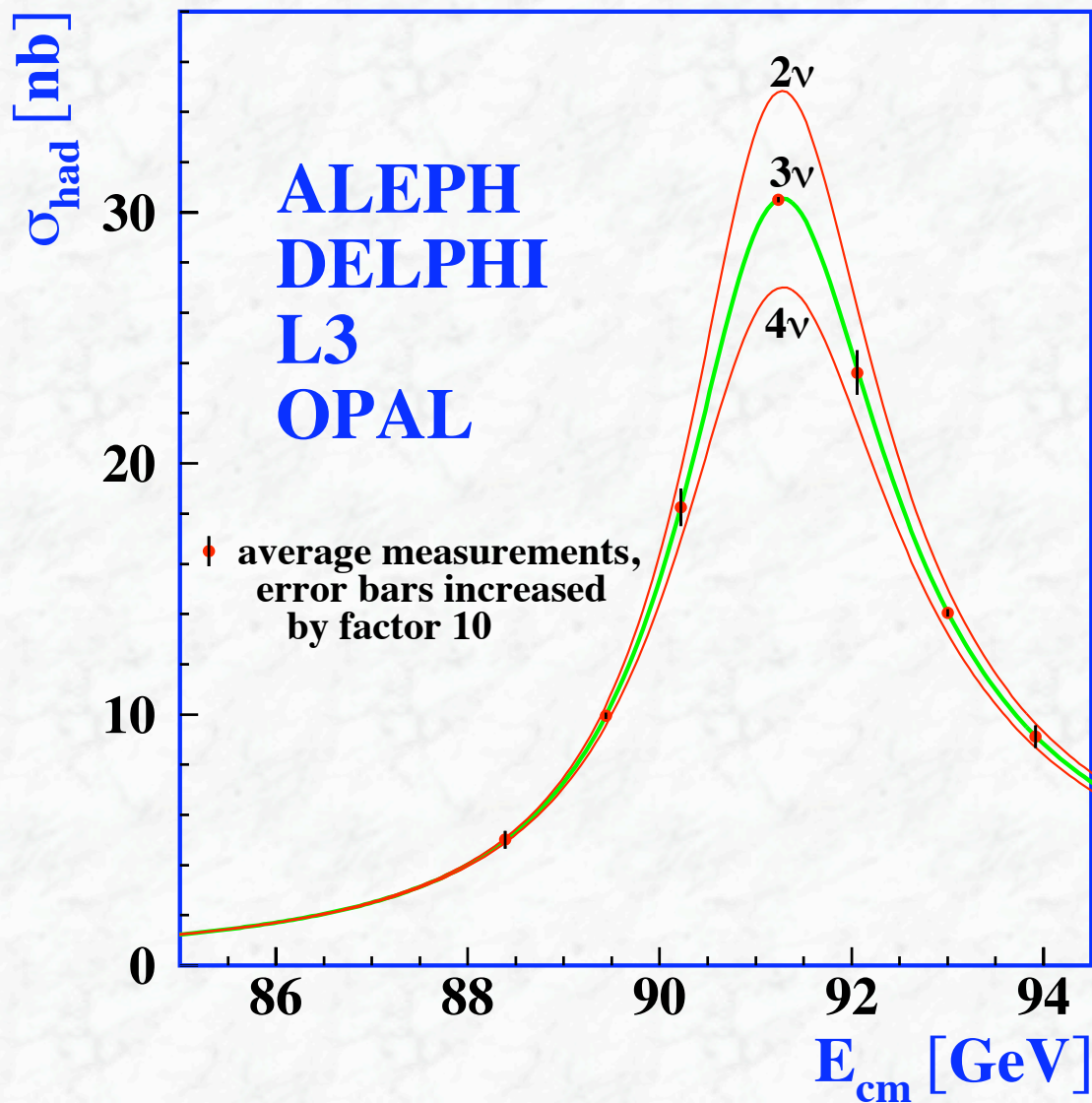
---

$$\begin{aligned}\Gamma_Z &= 2.4952 \pm 0.0023 \text{ GeV} \\ \Gamma_{\text{had}} &= 1.7444 \pm 0.0022 \text{ GeV} \\ \Gamma_e &= 0.083985 \pm 0.000086 \text{ GeV}\end{aligned}$$

lepton universality  
assumed:  
 $\Gamma_e = \Gamma_\mu = \Gamma_\tau$

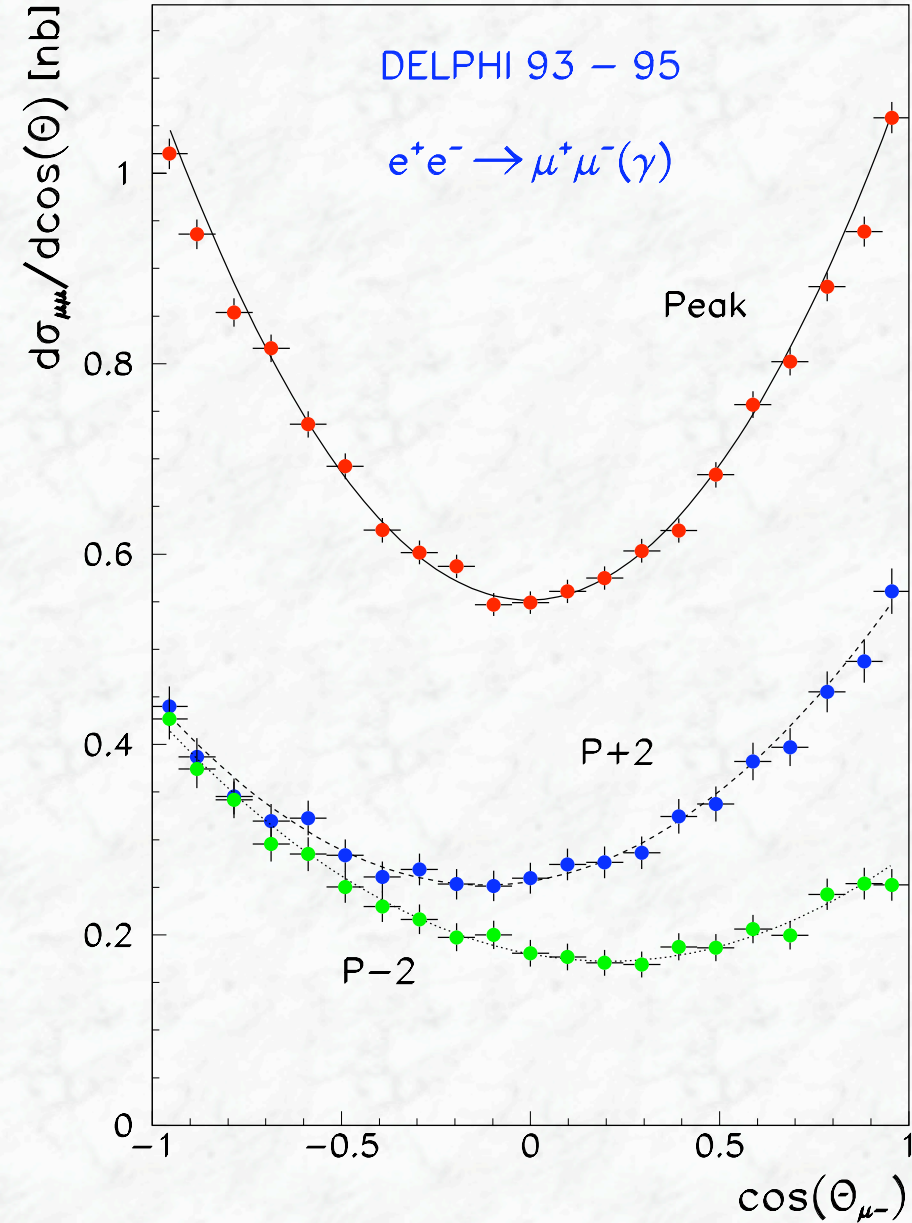
\*) Uncertainty on LEP energy measurement:  $\pm 1.7 \text{ MeV}$  (19 ppm)

# Number of neutrinos



$$N_\nu = 2.9840 \pm 0.0082$$

# Forward-backward asymmetries



$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2})(1 + \cos^2\theta) +$$

$$8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

Terms  $\propto \cos\theta$  in  $d\sigma/d\cos\theta$   
 $\rightarrow$  asymmetry

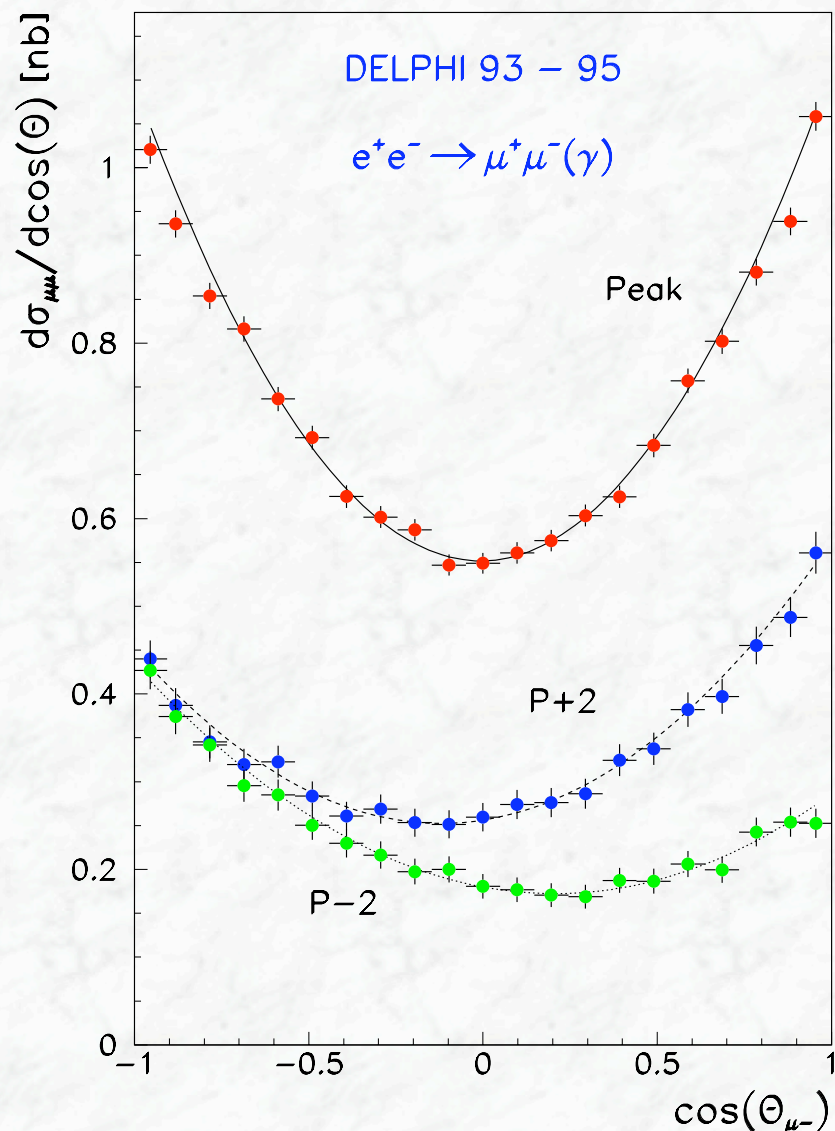
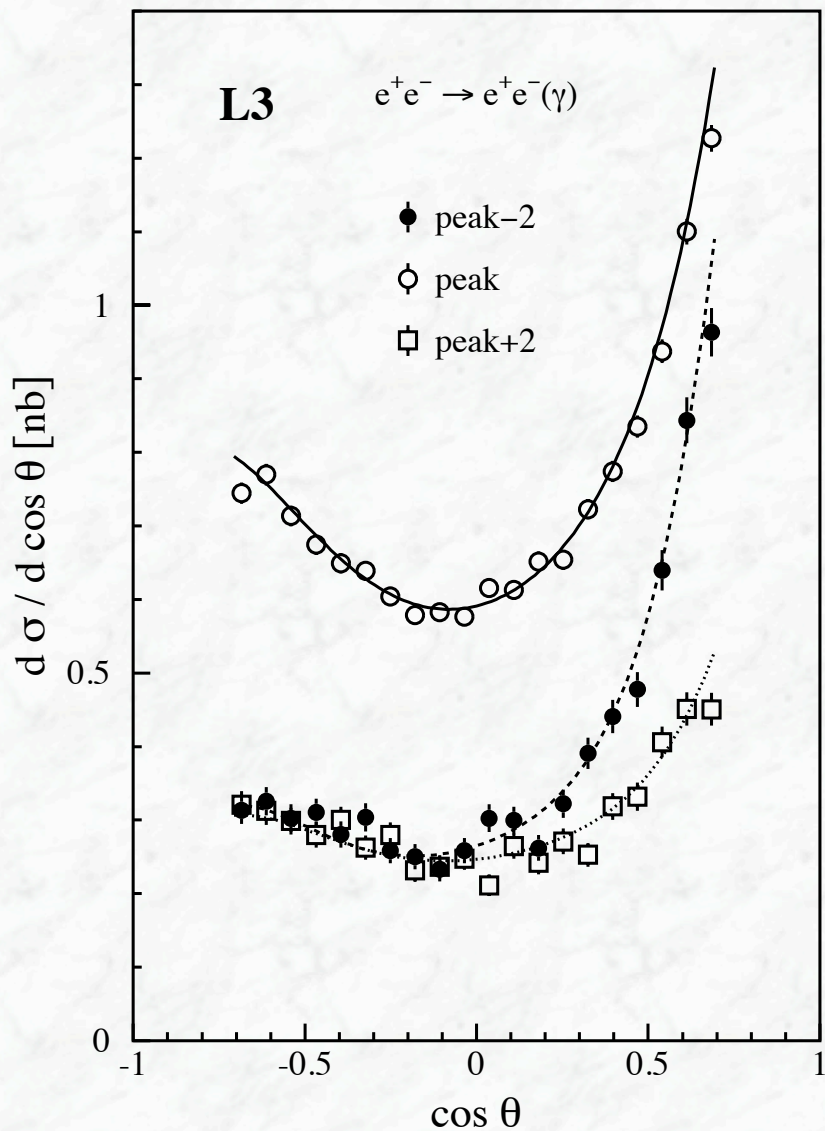
$$\sigma_{F(B)} = \int_{0(-1)}^{1(0)} \frac{d\sigma}{d\cos\theta} d\cos\theta$$

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$



# Forward-backward asymmetries

-comparison between ee and  $\mu\mu$  final states-



# Forward-backward asymmetries and fermion couplings

- Asymmetry at the Z pole (no interference) **is small**

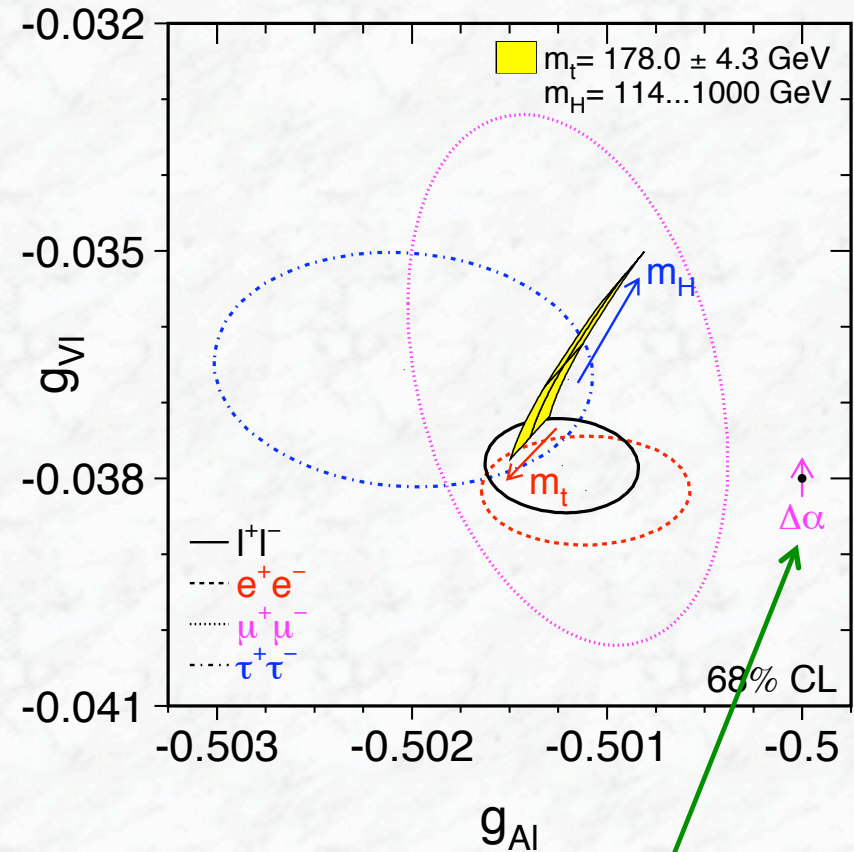
$$A_{\text{FB}} \sim g_A^e g_V^e g_A^f g_V^f$$

since  $g_V^f$  is small  
(in particular for leptons)

- For off-resonance points, the interference term dominates and gives larger contributions

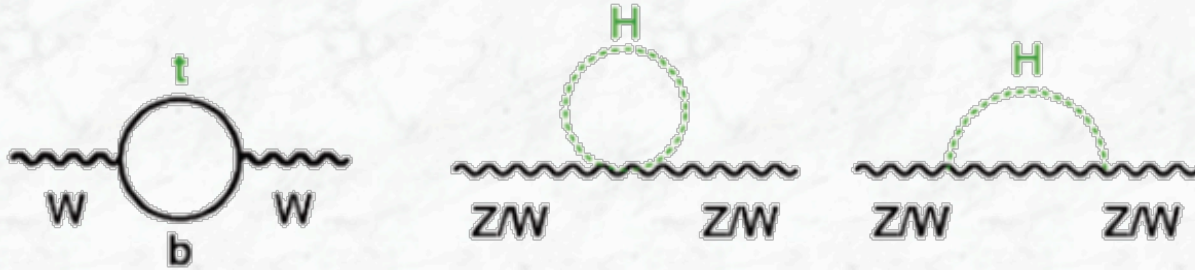
$$A_{\text{FB}} \sim g_A^e g_A^f \cdot \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

- $A_{\text{FB}}$  can be used for the determination of the fermion couplings



LO Standard Model prediction:  
 $g_A = T_3$   
 $g_V = T_3 - 2 Q \sin^2 \theta_W$

# Electroweak radiative corrections



Standard Model relations  
(lowest order)

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F}$$

$$\alpha(0)$$

Relations including  
radiative corrections

$$\vec{\rho} = 1 + \Delta\rho$$

$$\sin^2 \theta_{\text{eff}} = (1 + \Delta\kappa) \sin^2 \theta_W$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} \cdot \frac{1}{(1 - \Delta r)}$$

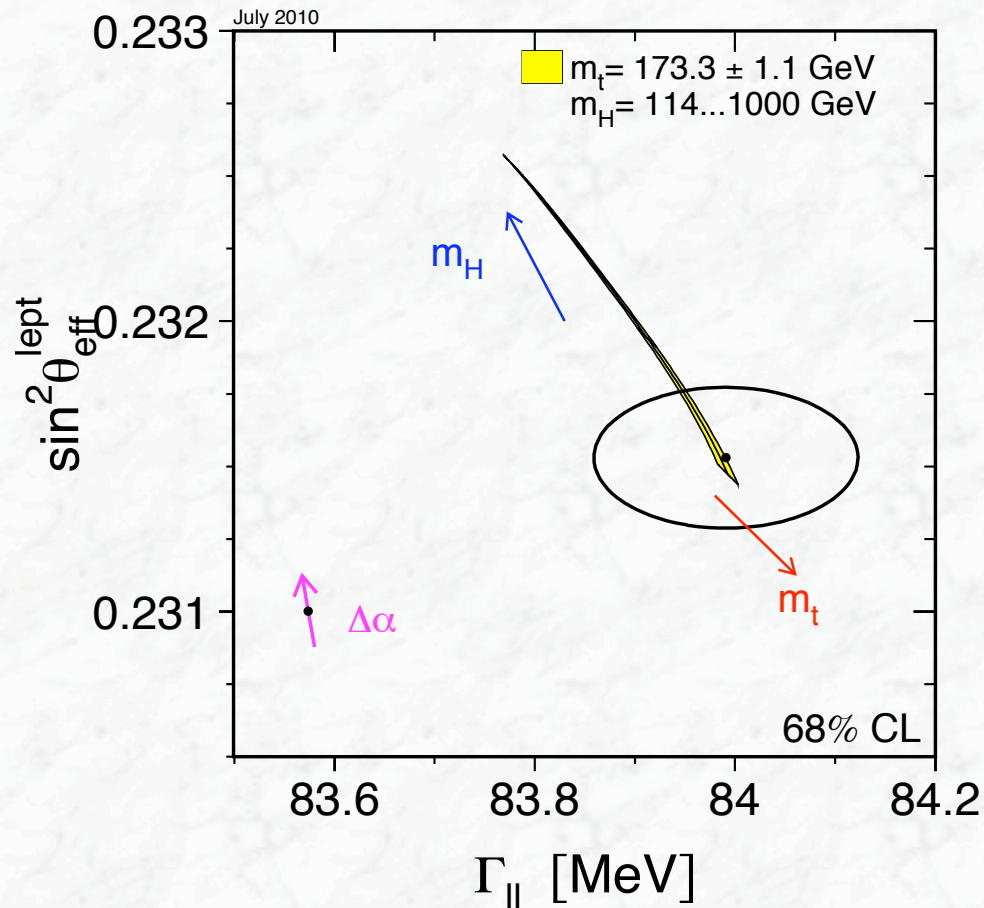
$$\alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta\alpha}$$

$$\Delta\alpha = \Delta\alpha_{\text{lepl}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$$

$$\Delta\rho, \Delta\kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$$

# Results of electroweak precision tests at LEP (cont.)

partial decay width versus  $\sin^2 \theta_W$ :



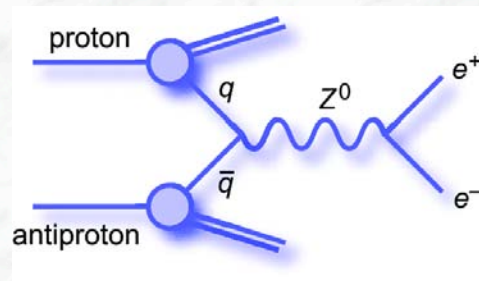
# Results of electroweak precision tests at LEP (cont.)

## Summary of results:

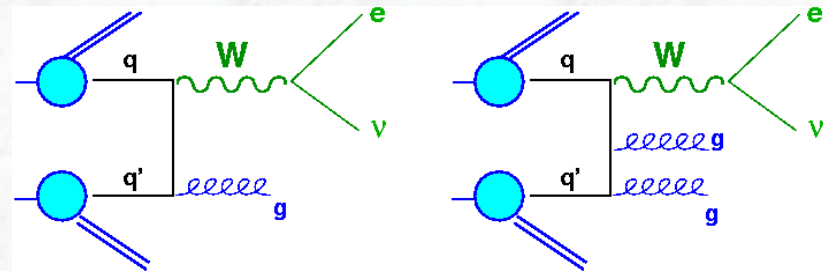
- All measurements in agreement with the Standard Model
- They can be described with a limited set of parameters



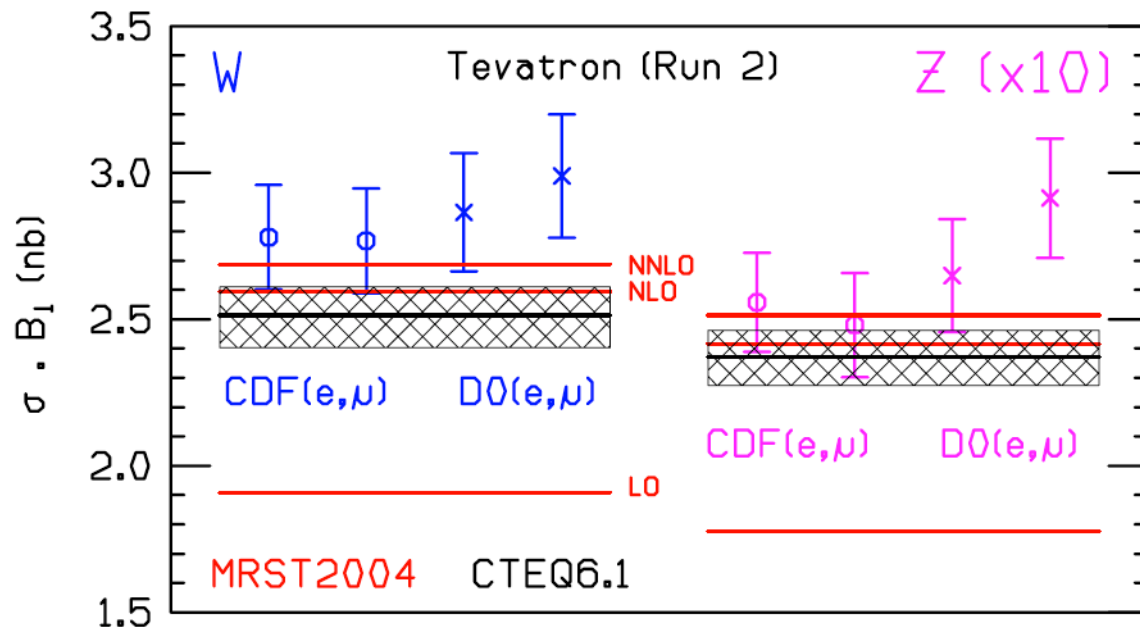
## 6.3 W/Z production at hadron colliders



QCD at work



- Important test of NNLO Drell-Yan QCD prediction for the total cross section
- Test of perturbative QCD in high  $p_T$  region (jet multiplicities,  $p_T$  spectra,....)
- Tuning and „calibration“ of Monte Carlos for background predictions in searches at the LHC



Predictions for the W and Z boson total cross sections at the Tevatron, using the MRST2004 and CTEQ pdfs, compared with measurements from the CDF and D0 collaborations. The predictions are shown at LO, NLO, and NNLO. For the NLO prediction the accompanying pdf uncertainties are shown as band.

## W boson production cross sections at the LHC ( $\sqrt{s} = 7 \text{ TeV}$ )

Program	Non-standard parameters	W charge	Cross Section (pb)
Cross sections for the full kinematic regime			
FEWZ	MSTW2008NNLO	$W^+$	$6160_{-55}^{+49}$ (scale) $\pm 111$ (PDF) $\pm 74$ ( $\alpha_s$ )
		$W^-$	$4301_{-34}^{+34}$ (scale) $\pm 69$ (PDF) $\pm 52$ ( $\alpha_s$ )
		$W^+ + W^-$	$10461_{-94}^{+84}$ (scale) $\pm 167$ (PDF) $\pm 126$ ( $\alpha_s$ )
ZWPRODMS	MSTW2008NNLO	$W^+$	$6189_{-50}^{+33}$ (scale) $\pm 105$ (PDF) $\pm 67$ ( $\alpha_s$ )
		$W^-$	$4316_{-33}^{+25}$ (scale) $\pm 72$ (PDF) $\pm 44$ ( $\alpha_s$ )
		$W^+ + W^-$	$10506_{-83}^{+58}$ (scale) $\pm 173$ (PDF) $\pm 111$ ( $\alpha_s$ )
Cross sections for the kinematic regime of Eq. 1			
FEWZ	MSTW2008NNLO	$W^+$	2907
		$W^-$	1927
		$W^+ + W^-$	4833

Predictions for the  $W \rightarrow l\nu$  cross section at NNLO, calculated for the full kinematic range as well as in the fiducial region (see below).

Major uncertainties: renormalization and factorization scale ( $\sim \pm 1\%$ )  
parton distribution functions ( $\sim \pm 2\%$ )  
uncertainties of  $\alpha_s$  ( $\sim \pm 1\%$ )

Fiducial region:  $PT(l) > 20 \text{ GeV}$ ,  $\eta < 2.47$ , excluding  $1.37 < \eta < 1.52$   
 $E_T^{\text{miss}} > 25 \text{ GeV}$   
 $m_T > 40 \text{ GeV}$



## Z boson production cross sections at the LHC ( $\sqrt{s} = 7$ TeV)

Program	Non-standard parameters	Mass range (GeV)	Cross Section (pb)
Cross section for the full kinematic regime			
FEWZ	MSTW2008NNLO	> 60	$989_{-7}^{+5}$ (scale) $\pm 16$ (PDF) $\pm 10$ ( $\alpha_s$ )
		60 – 120	$978_{-7}^{+5}$ (scale) $\pm 16$ (PDF) $\pm 10$ ( $\alpha_s$ )
		66 – 116	$964_{-7}^{+5}$ (scale) $\pm 15$ (PDF) $\pm 10$ ( $\alpha_s$ )
		70 – 110	$952_{-7}^{+5}$ (scale) $\pm 15$ (PDF) $\pm 10$ ( $\alpha_s$ )
		80 – 100	$904_{-6}^{+5}$ (scale) $\pm 14$ (PDF) $\pm 9$ ( $\alpha_s$ )
		only Z, full range	$970_{-7}^{+5}$ (scale) $\pm 15$ (PDF) $\pm 10$ ( $\alpha_s$ )
ZWPRODMS		only Z, full range	$974_{-6}^{+5}$ (scale) $\pm 16$ (PDF) $\pm 10$ ( $\alpha_s$ )
Cross section for the kinematic regime specified in Eq. 2			
FEWZ	MSTW2008NNLO	66 – 116	420

Predictions for the  $Z / \gamma^* \rightarrow \ell\ell$  cross section at NNLO, calculated for the full kinematic range as well as in the fiducial region (see below).

Major uncertainties: renormalization and factorization scale ( $\sim \pm 1\%$ )  
parton distribution functions ( $\sim \pm 1.5\%$ )  
uncertainties of  $\alpha_s$  ( $\sim \pm 1\%$ )

Fiducial region:  $PT(\ell) > 20$  GeV,  $\eta < 2.47$ , excluding  $1.37 < \eta < 1.52$   
 $66 < m_{\ell\ell} < 116$  GeV

## 6.4 Test of QCD in W/Z production at hadron colliders

As explained, leptons, photons and missing transverse energy are key signatures at hadron colliders

→ Search for leptonic decays:  $W \rightarrow \ell \nu$  (large  $P_T(\ell)$ , large  $P_T^{\text{miss}}$ )  
 $Z \rightarrow \ell \ell$

More difficult:  $W \rightarrow \tau \nu \rightarrow \text{had } \nu \nu$        $Z \rightarrow \tau \tau \rightarrow e(\mu) \nu \nu \text{ had } \nu$

Ingredients for a cross-section measurement:

$$\sigma = \frac{N_{sel} - N_{back}}{L \cdot \varepsilon \cdot \eta}$$

where:  $N_{sel}$  = number of selected events

$N_{back}$  = number of background events in selected events

$L$  = integrated luminosity (measured from machine, reference process)

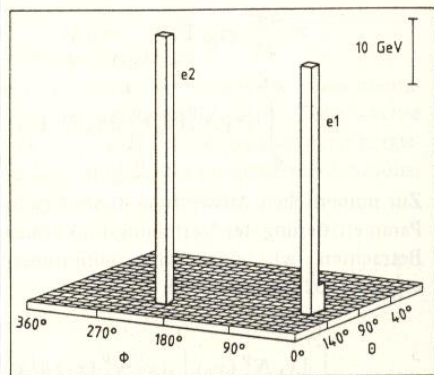
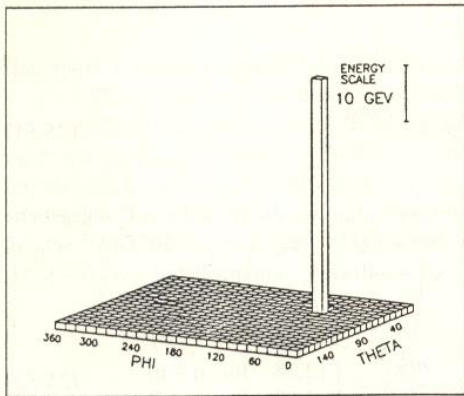
$\varepsilon$  = detection efficiency

$\eta$  = acceptance of fiducial cuts ( $P_T(l)$ ,  $E_T^{\text{miss}}$ ,  $M_T$ ,  $m_{ll}, \dots$ )

# How do W and Z events look like ?

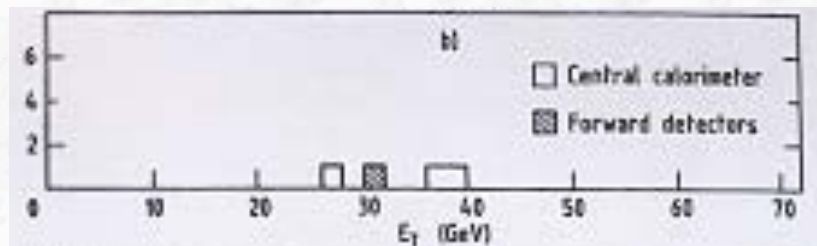
A bit of history: one of the first W and Z events seen (UA2 experiment)

W/Z discovery by the UA1 and UA2 experiments at CERN (1983/84)



Carlo Rubbia (left, UA1) and Luigi Di Lella (right, UA2)

Transverse momentum of the electrons



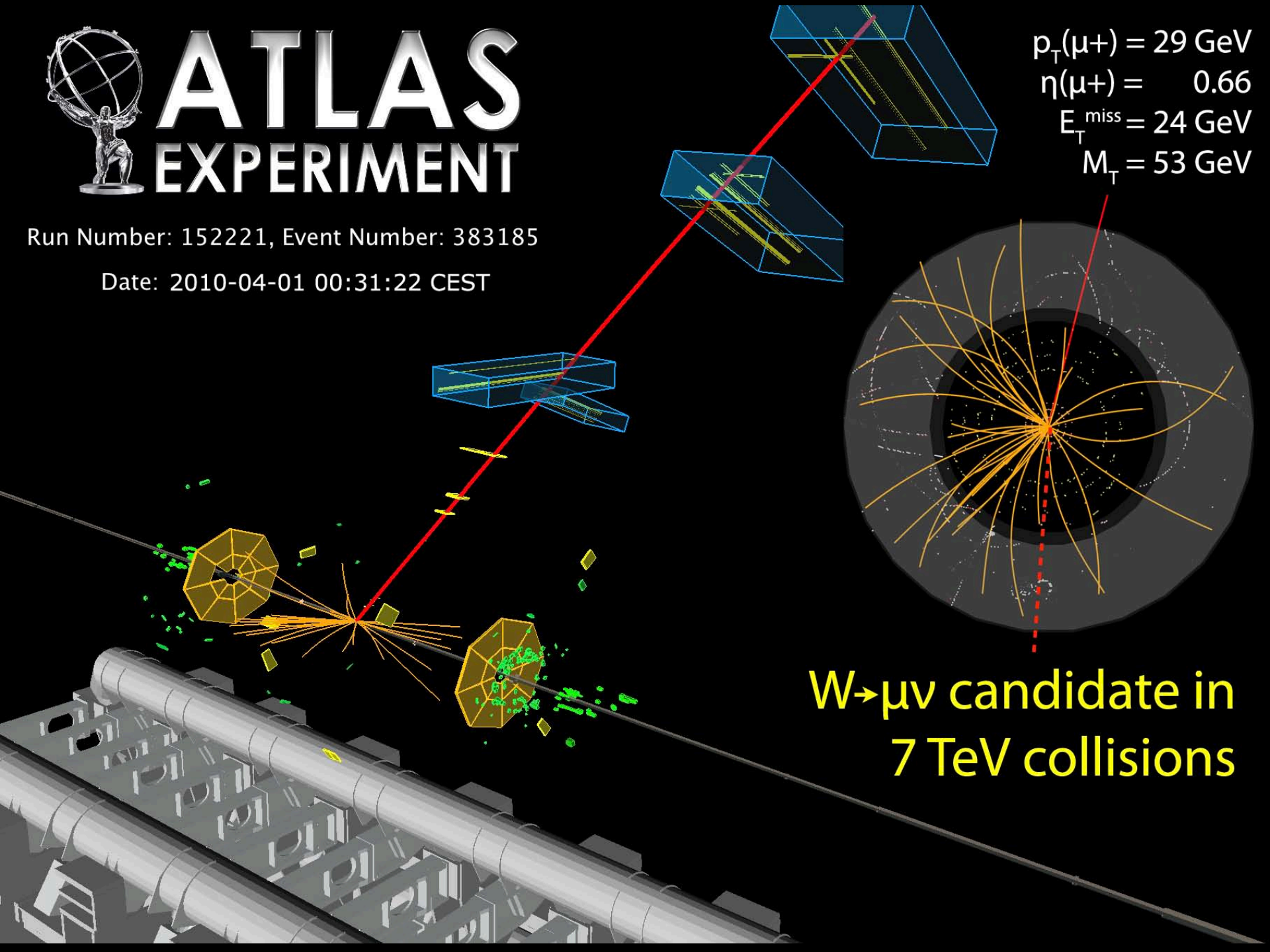


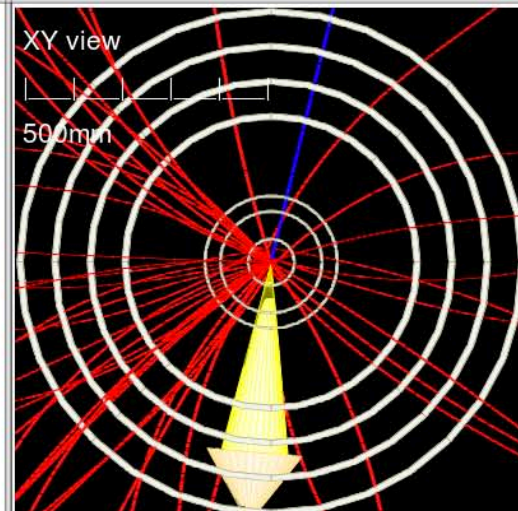
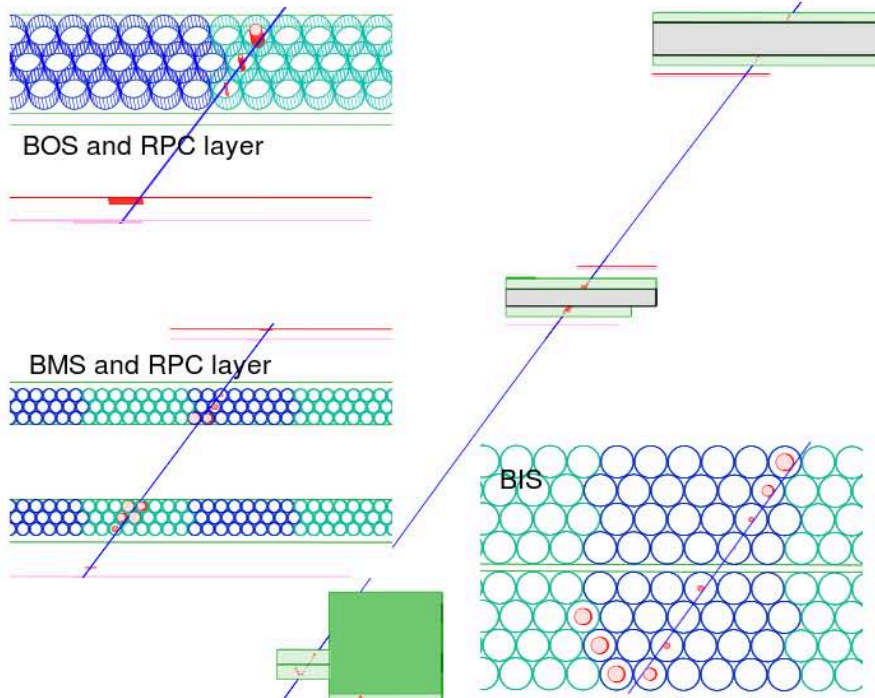
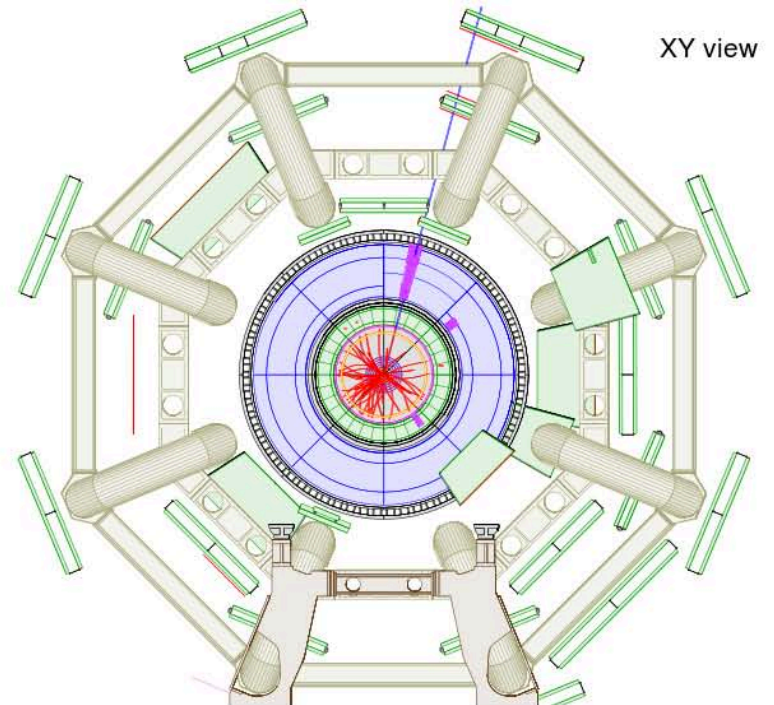
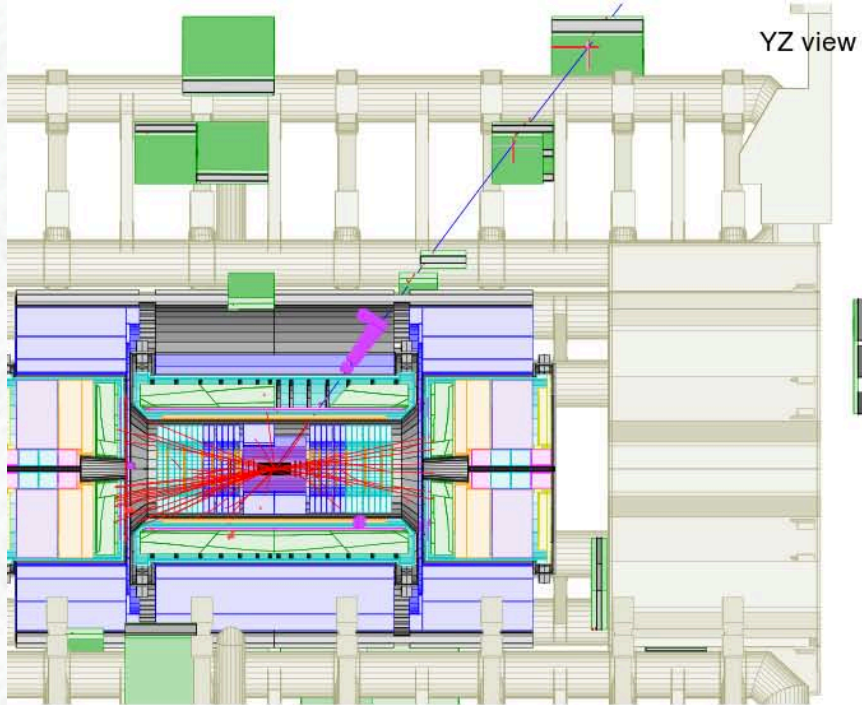
# ATLAS EXPERIMENT

Run Number: 152221, Event Number: 383185

Date: 2010-04-01 00:31:22 CEST

$$\begin{aligned} p_T(\mu^+) &= 29 \text{ GeV} \\ \eta(\mu^+) &= 0.66 \\ E_T^{\text{miss}} &= 24 \text{ GeV} \\ M_T &= 53 \text{ GeV} \end{aligned}$$





W- $\mu$ v candidate in 7 TeV collisions

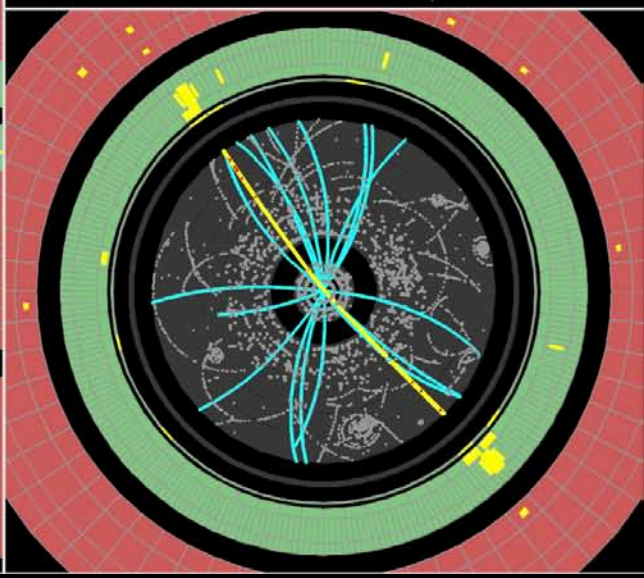
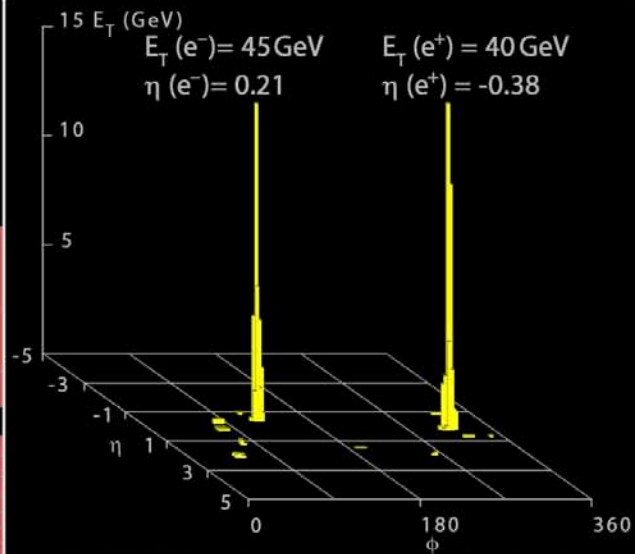
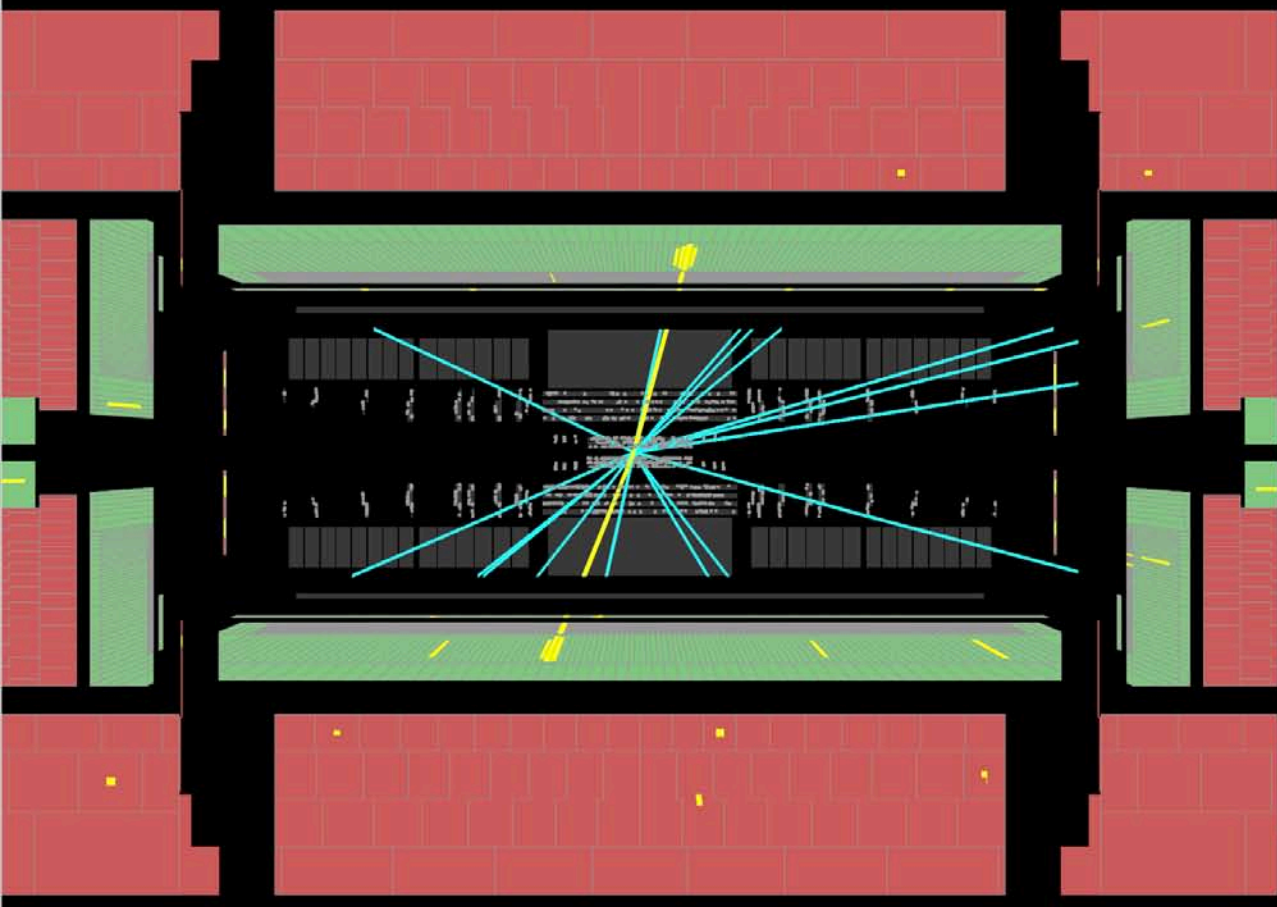
Run Number: 152221, Event Number: 383185  
 Date: 2010-04-01 00:31:22 CEST  
 PT( $\mu$ +) = 29 GeV,  $\eta$  = 0.66  
 ETmis = 24 GeV  
 MT = 53 GeV



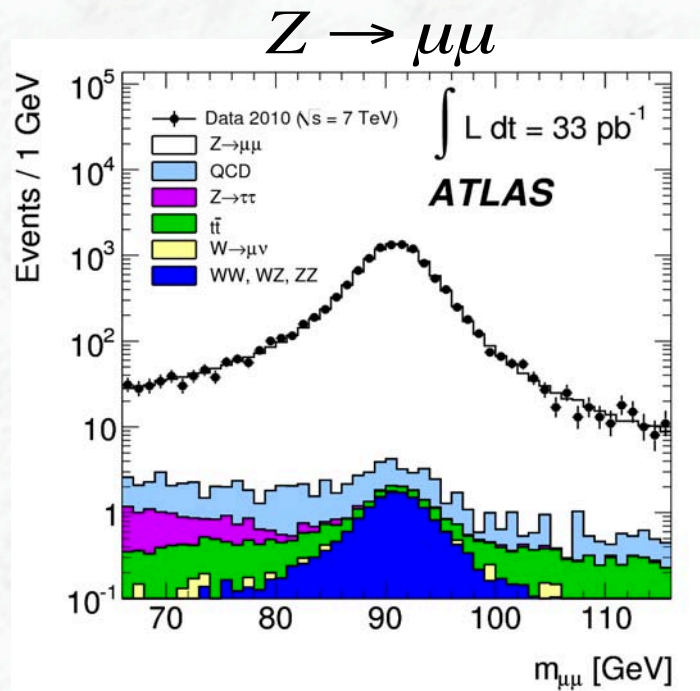
Run Number: 154817, Event Number: 968871  
Date: 2010-05-09 09:41:40 CEST

$M_{ee} = 89 \text{ GeV}$

$Z \rightarrow ee$  candidate in 7 TeV collisions



# Today's $W/Z \rightarrow e\nu(\mu\nu) / ee(\mu\mu)$ signals ATLAS Experiment



## Trigger:

- Single-lepton candidate  $> 10$ - $25 \text{ GeV}/c$  (early 2010 data – Run 2)

## Electrons:

- Isolated el.magn. cluster in the calorimeter
- Shower shape consistent with expectation for electrons
- Matched with tracks
- $P_T > 25 \text{ GeV}/c$

## Muons:

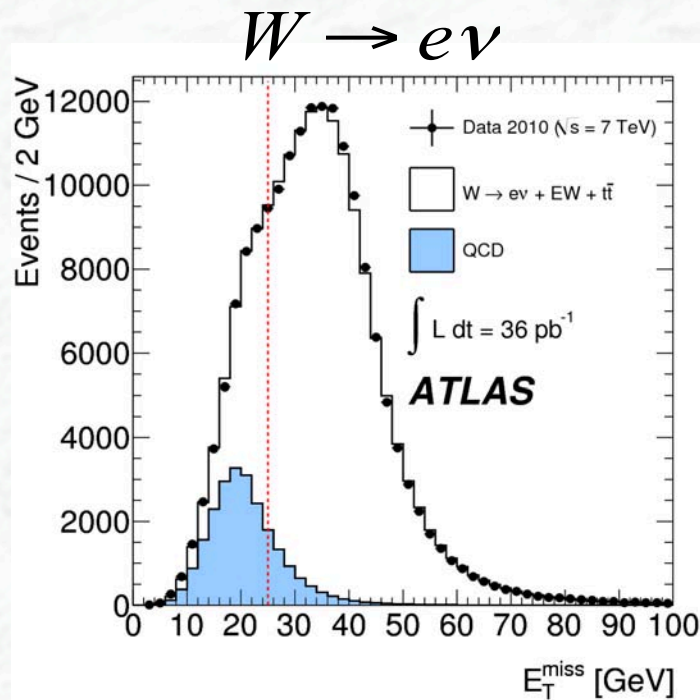
- Inner detector track matched with a muon spectrometer track (segment)

## $Z \rightarrow ee$

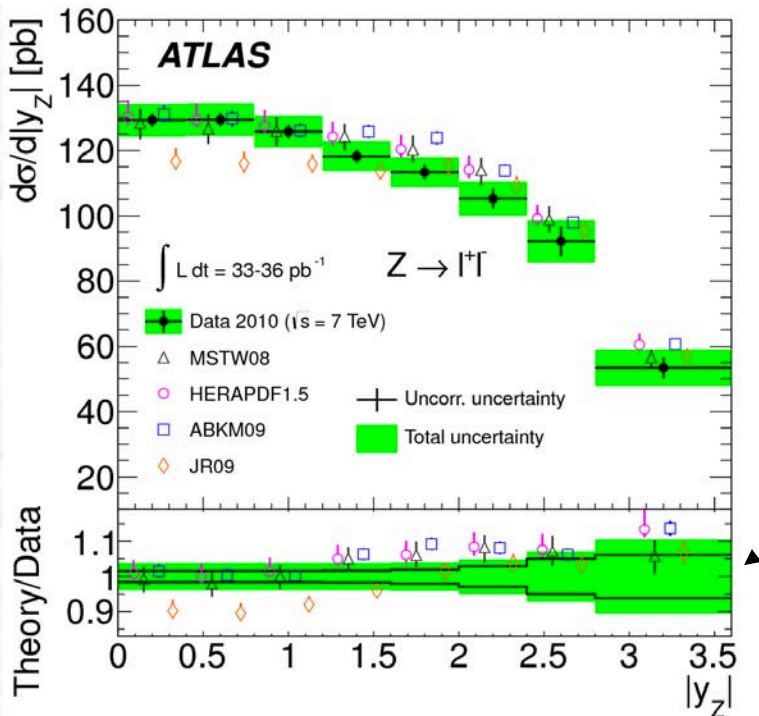
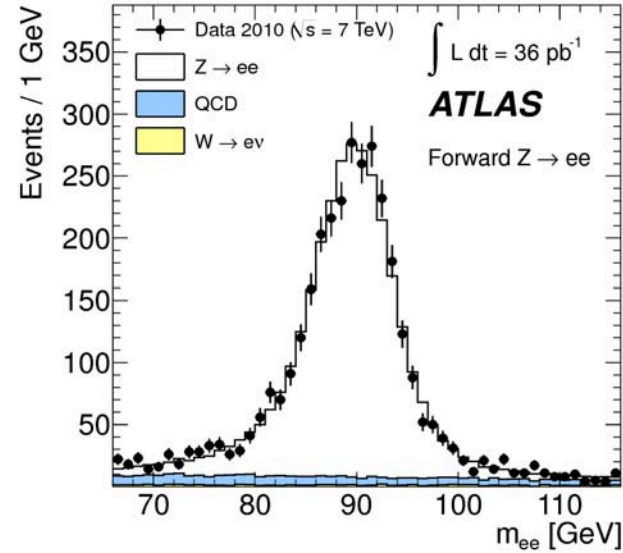
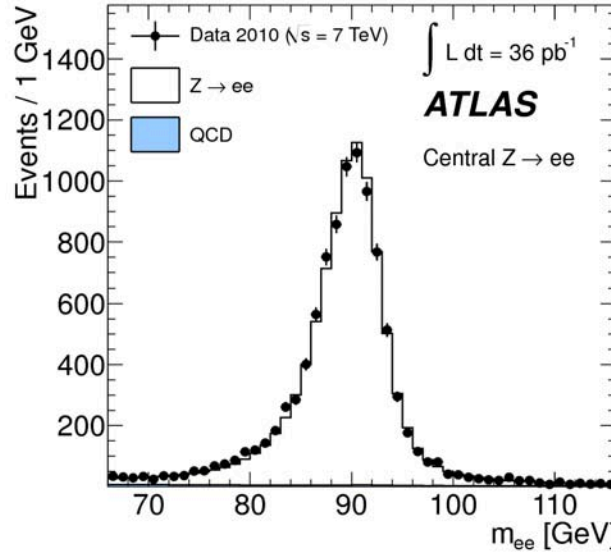
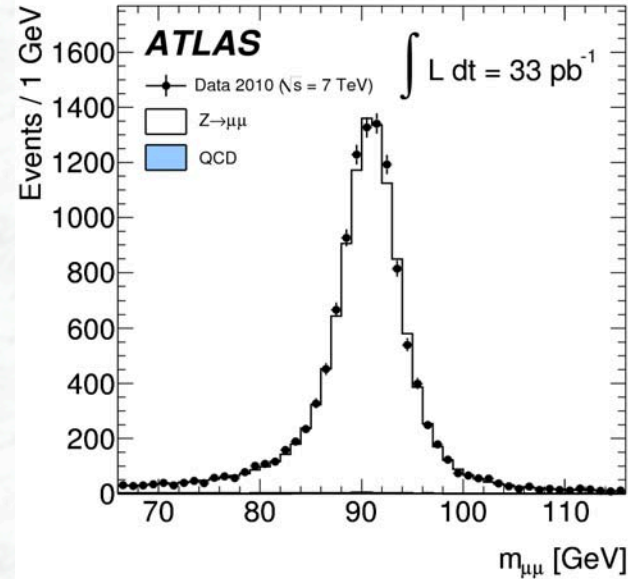
- $70 \text{ GeV}/c^2 < m_{ee} < 110 \text{ GeV}/c^2$

## $W \rightarrow e\nu$

- Missing transverse momentum  $> 25 \text{ GeV}/c$



# Z → ℓℓ cross sections



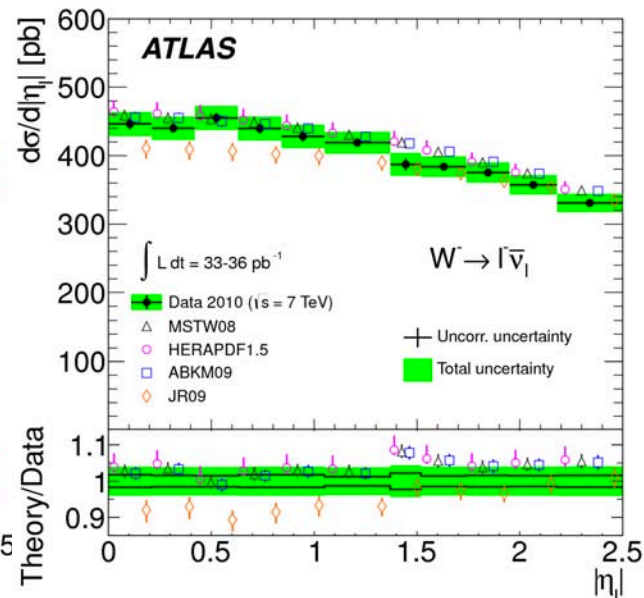
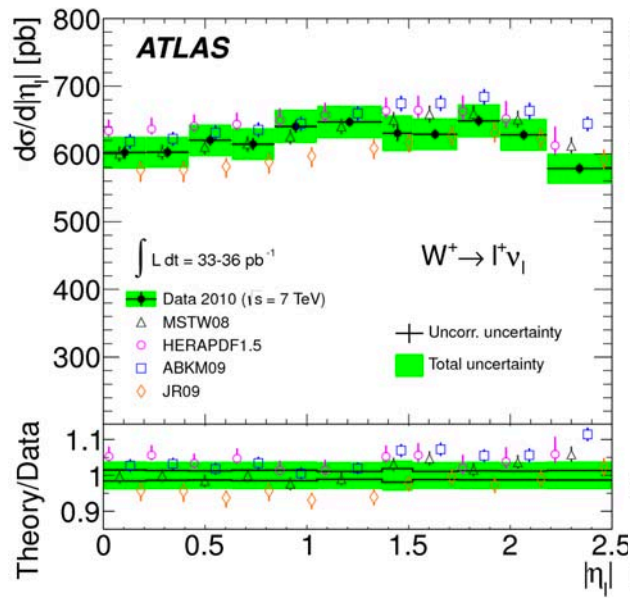
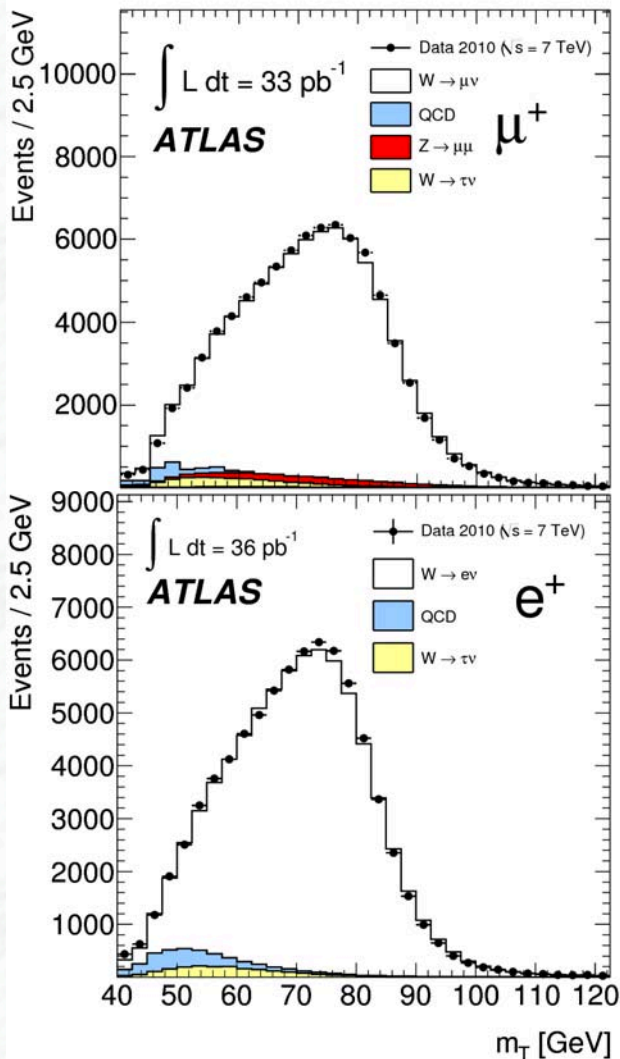
$$x_{1,2} = \frac{M_{W,Z}}{\sqrt{s}} e^{\pm y_{W,Z}}$$

Good agreement with NNLO QCD calculations, QCD corrections are large: factor ~ 1.25

Precision is limited by systematic effects (uncertainties on luminosity, parton densities,...)



# W → ℓν Cross Section



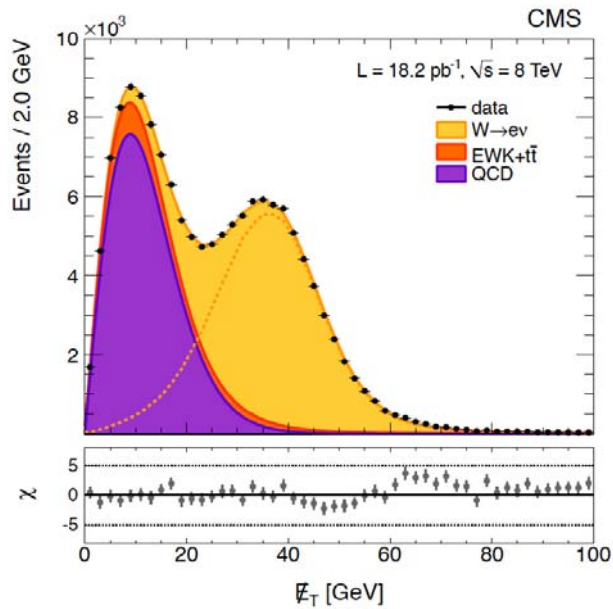
Good agreement with NNLO QCD calculations

$$M_W^T = \sqrt{2 \cdot P_T^l \cdot P_T^{\nu} \cdot (1 - \cos \Delta\phi^{l,\nu})}$$

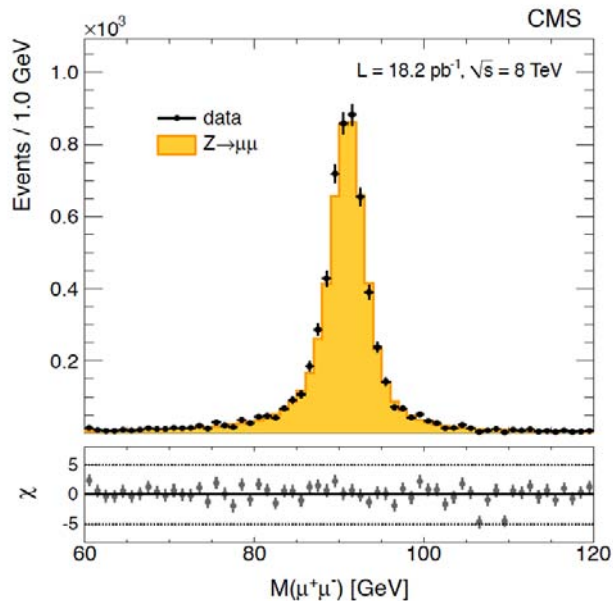
Note: the longitudinal component of the neutrino cannot be measured  
 → only transverse mass can be reconstructed

Precision is limited by systematic effects (uncertainties on luminosity, parton densities,...)

# First measurements of W/Z production at the LHC -CMS data from 2010: 36 pb<sup>-1</sup> -

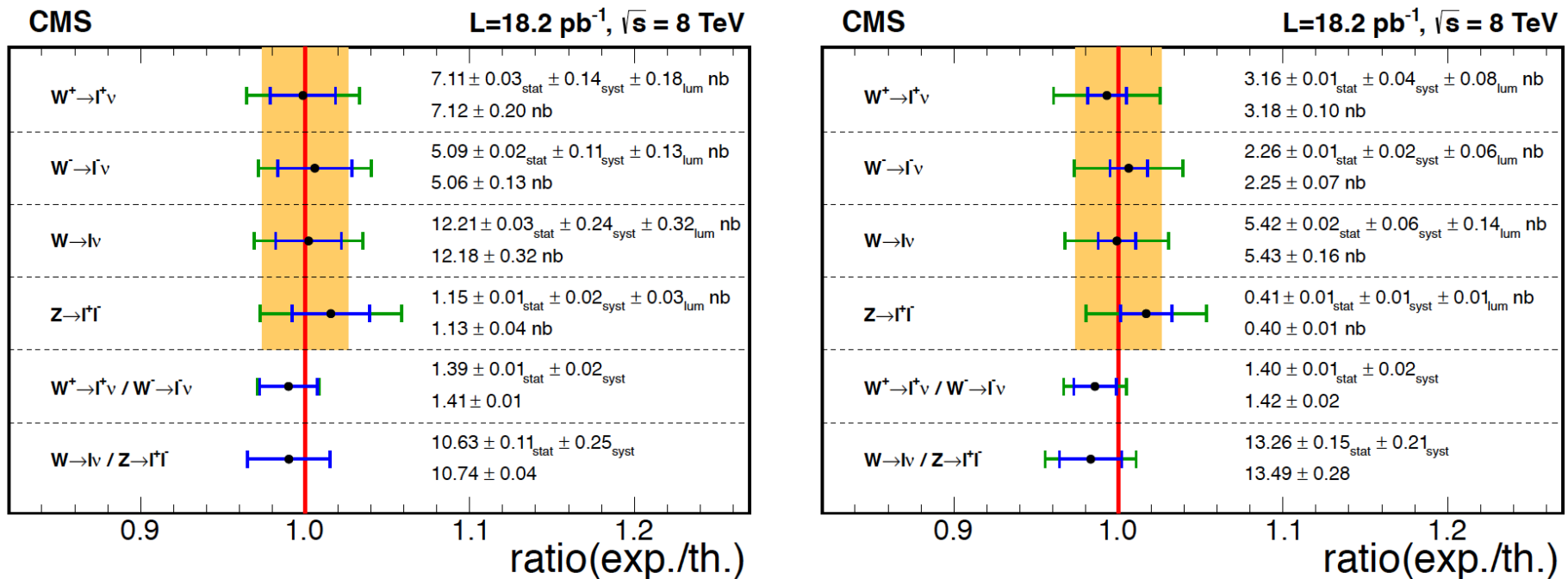


Distributions of the missing transverse energy,  $E_T^{\text{miss}}$ , of electron candidates for data and Monte-Carlo simulation, broken down into the signal and various background components.



Distributions of the invariant di-electron mass,  $m_{ee}$ , for events passing the Z selection. The data are compared to Monte-Carlo simulation, the background is very small.

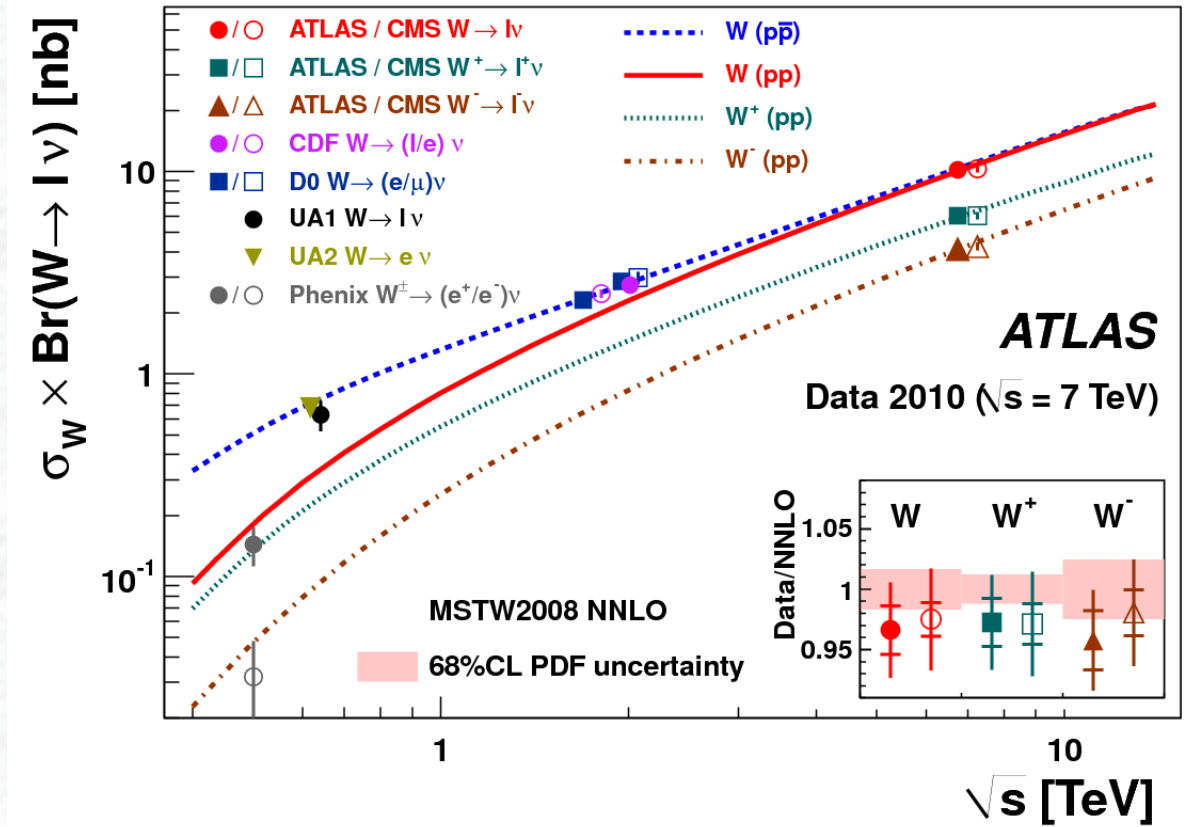
# W and Z production cross sections at LHC



Summary of total (left) and fiducial (right) inclusive  $W^+$ ,  $W^-$ ,  $W$ , and  $Z$  production cross sections times branching fractions,  $W$  to  $Z$  and  $W^+$  to  $W^-$  ratios, and their theoretical predictions. The shaded box indicates the uncertainties in the luminosity measurement. The inner error bars represent the experimental uncertainties, while outer error bars also include the uncertainties in the theoretical predictions. The individual measurements and theoretical predictions are given numerically on the right. The theoretical predictions of cross sections and cross section ratios are computed at NNLO with the program FEWZ and the MSTW2008 set of PDFs.

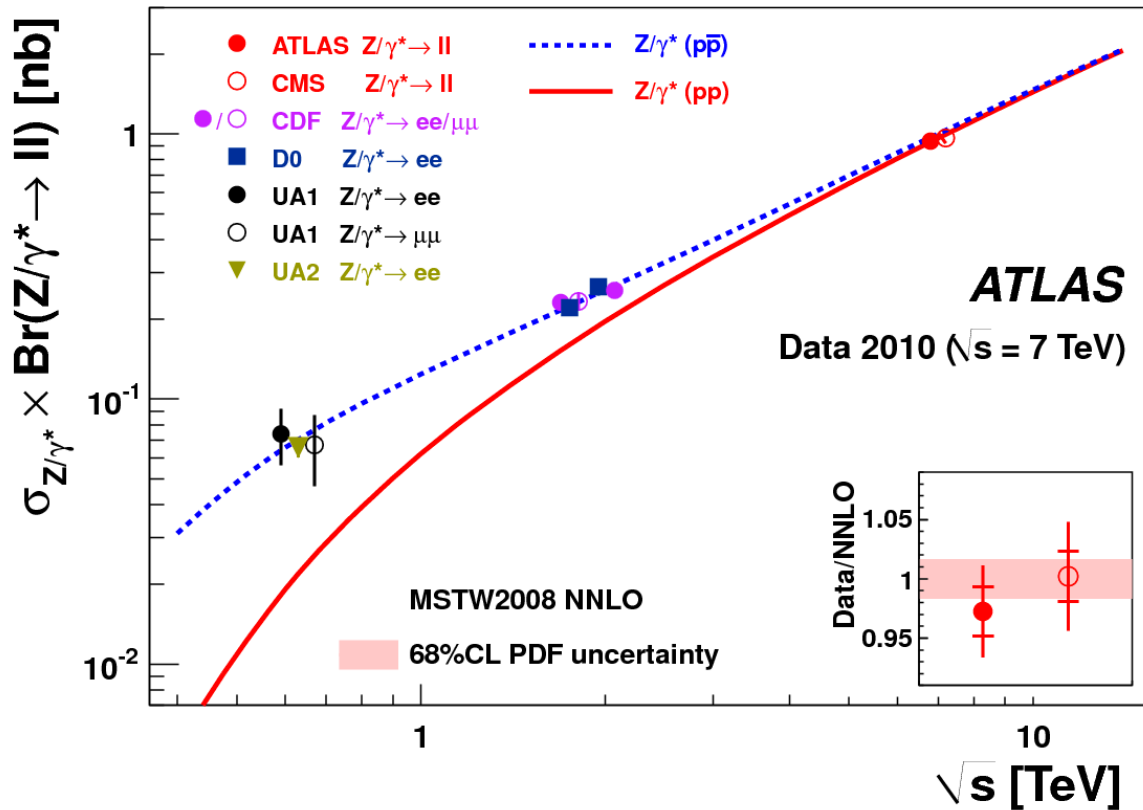
Good agreement between data and NNLO QCD predictions for all measurements

# W production cross sections at hadron colliders



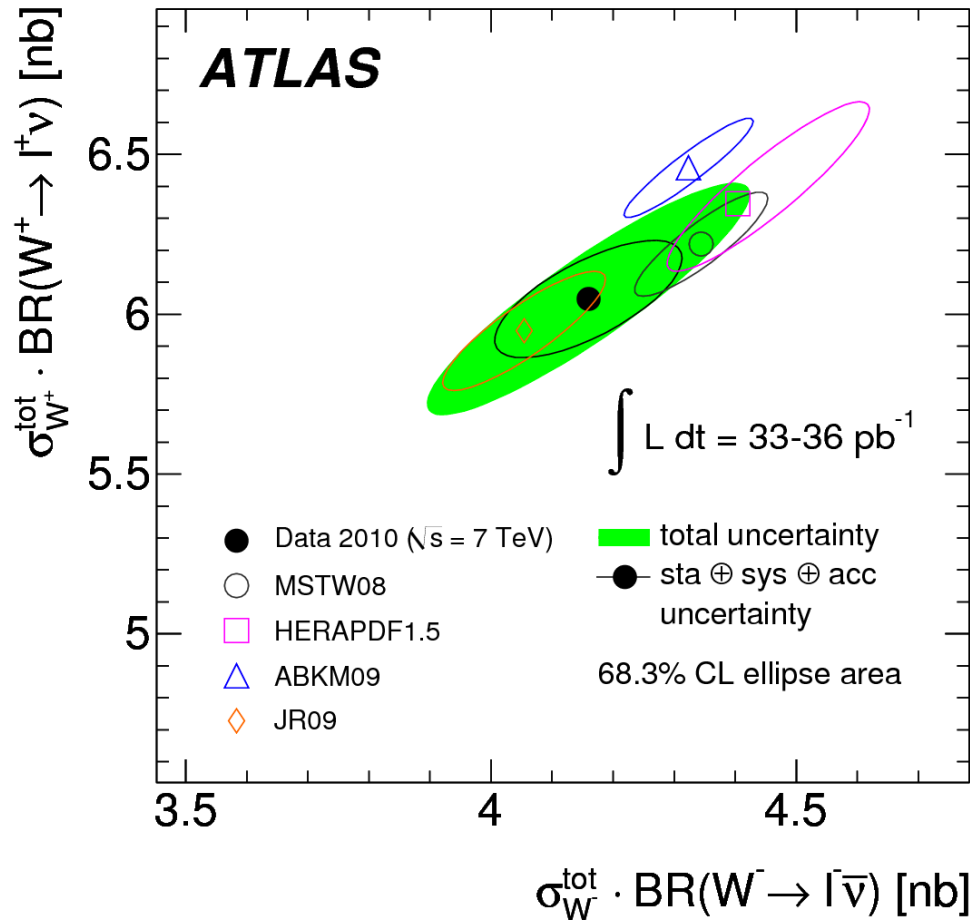
The measured values of  $\sigma(W) \times \text{BR}(W \rightarrow l\nu)$  for  $W^+$ ,  $W^-$  and for their sum compared to the theoretical predictions based on NNLO QCD calculations using the MSTW 2008 PDF set. Results are shown for the combined electron-muon results. The predictions are shown for both proton-proton ( $W^+$ ,  $W^-$  and their sum) and proton-antiproton colliders ( $W$ ) as a function of  $\sqrt{s}$ . In addition, previous measurements at proton-antiproton and proton-proton colliders are shown. The data points at the various energies are staggered to improve visibility. The CDF and D0 measurements are shown for both Tevatron collider energies,  $\sqrt{s} = 1.8$  TeV and  $\sqrt{s} = 1.96$  TeV. All data points in the main plot are displayed with their total uncertainty. The small inset shows the results at  $\sqrt{s} = 7$  TeV for the ATLAS and CMS collaborations as ratio of measurement to NNLO prediction, where the inner error bars denote all but the luminosity uncertainty and the outer error bar the total uncertainty. Only in the insert the theoretical uncertainties are shown, where just the 68 per cent CL PDF errors of MSTW2008 are considered.

# Z production cross sections at hadron colliders



The measured value of  $\sigma(Z/\gamma^*) \times \text{BR}(Z/\gamma^* \rightarrow \ell\ell)$  where the electron and muon channels have been combined, compared to the theoretical predictions based on NNLO QCD calculations using the MSTW 2008 PDF set. The predictions are shown for both proton-proton and proton-antiproton colliders as a function of  $\sqrt{s}$ . In addition, previous measurements at proton-antiproton colliders are shown. The data points at the various energies are staggered to improve readability. The CDF and D0 measurements are shown for both Tevatron collider energies,  $\sqrt{s} = 1.8$  TeV and  $\sqrt{s} = 1.96$  TeV. All data points in the main plot are displayed with their total uncertainty. The small inset shows the results at  $\sqrt{s} = 7$  TeV for the ATLAS and CMS collaborations as ratio of measurement to NNLO prediction, where the inner error bars denote all but the luminosity uncertainty and the outer error bar the total uncertainty. Only in the insert the theoretical uncertainties are shown, where just the 68 per cent CL PDF errors of MSTW2008 are considered. Note also, that various measurements were performed in slightly different invariant mass ranges, while the prediction is for  $66 < m(\ell\ell) < 116$  GeV. The difference are on the order of a few % and not visible on the double-log scale. For the small insert the CMS result is divided by the appropriate prediction for the invariant mass range of the measurement, which was  $60 < m(\ell\ell) < 120$  GeV.

# W cross sections at the LHC, charge separated

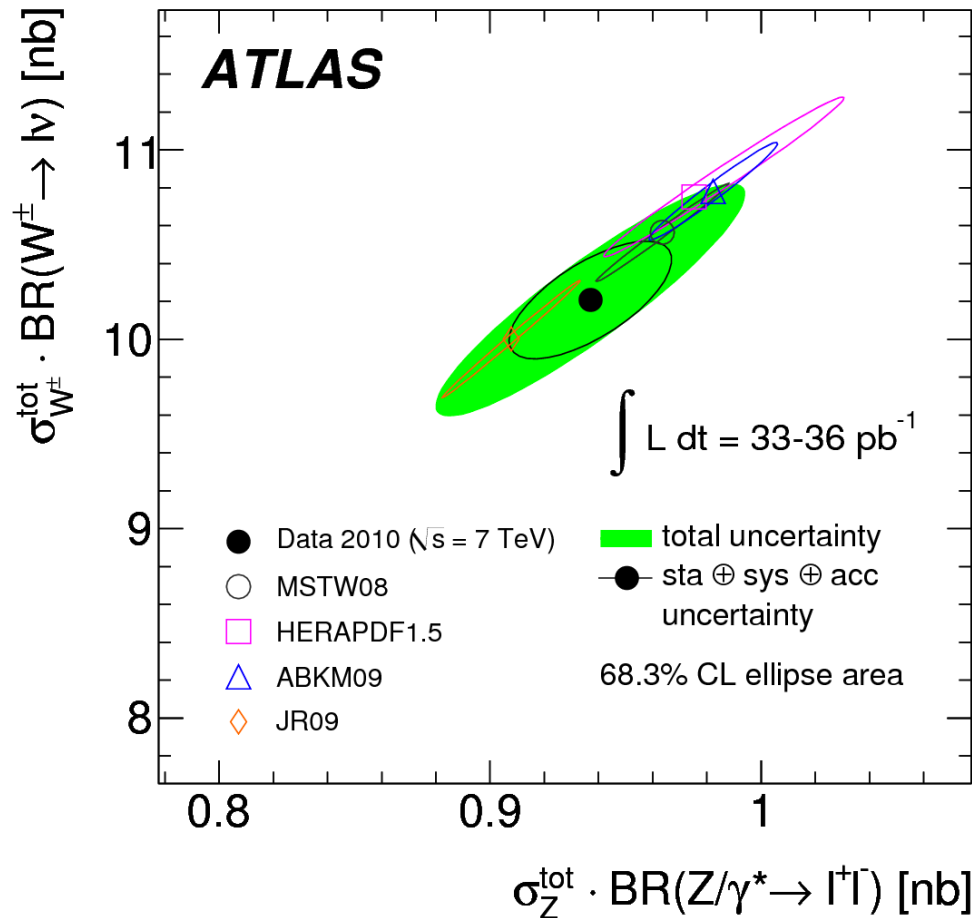


Full ATLAS data set  
from 2010

$L = 36 \text{ pb}^{-1}$

Measured and predicted total cross sections times leptonic branching ratios,  $\sigma(W^+)$  vs  $\sigma(W^-)$ . The ellipses illustrate the 68 per cent CL coverage for total uncertainties (full green) and excluding the luminosity uncertainty (open black). The uncertainties of the theoretical predictions correspond to the PDF uncertainties only.

# W cross sections at the LHC, charge separated

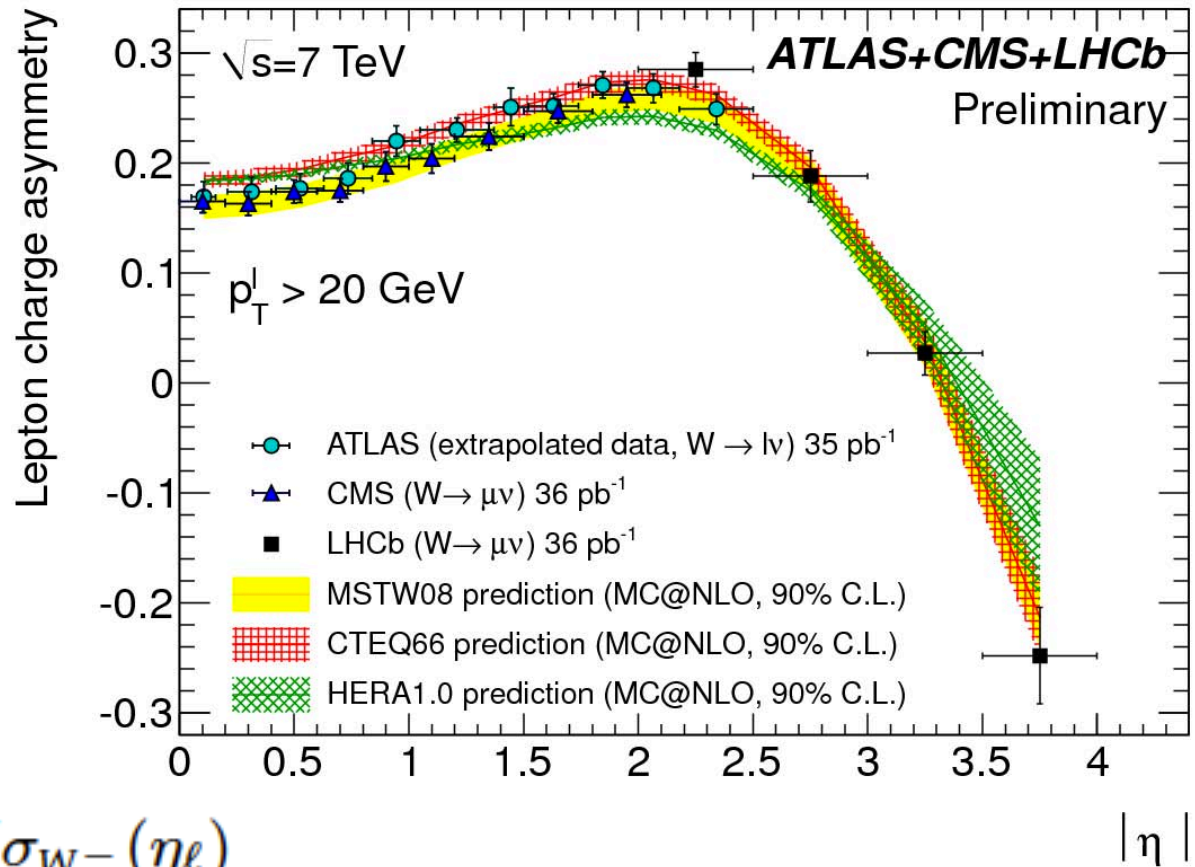


Full ATLAS data set  
from 2010

$L = 36 \text{ pb}^{-1}$

Measured and predicted total cross sections times leptonic branching ratios,  $(\sigma(W^+) + \sigma(W^-))$  vs.  $\sigma(Z/\gamma^*)$ . The ellipses illustrate the 68 per cent CL coverage for total uncertainties (full green) and excluding the luminosity uncertainty (open black). The uncertainties of the theoretical predictions correspond to the PDF uncertainties only.

# W charge asymmetry as a function of pseudorapidity



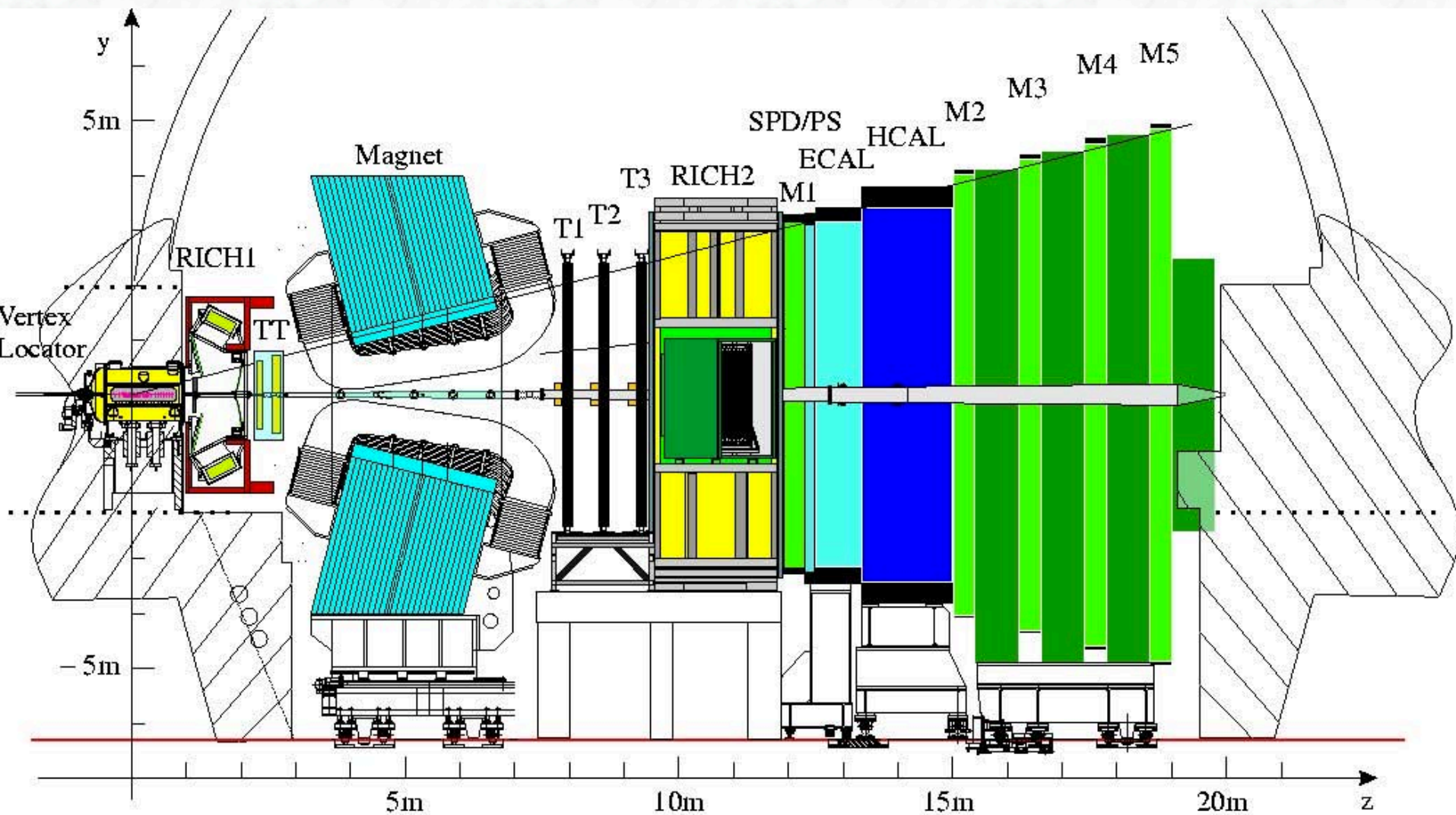
$$A(\eta_e) = \frac{d\sigma_{W^+}(\eta_e) - d\sigma_{W^-}(\eta_e)}{d\sigma_{W^+}(\eta_e) + d\sigma_{W^-}(\eta_e)}$$

- Sensitive to valence quark ( $u\bar{d} \rightarrow W^+ / d\bar{u} \rightarrow W^-$ )
- Usable to constrain  $u_v/d_v$  at low  $x$

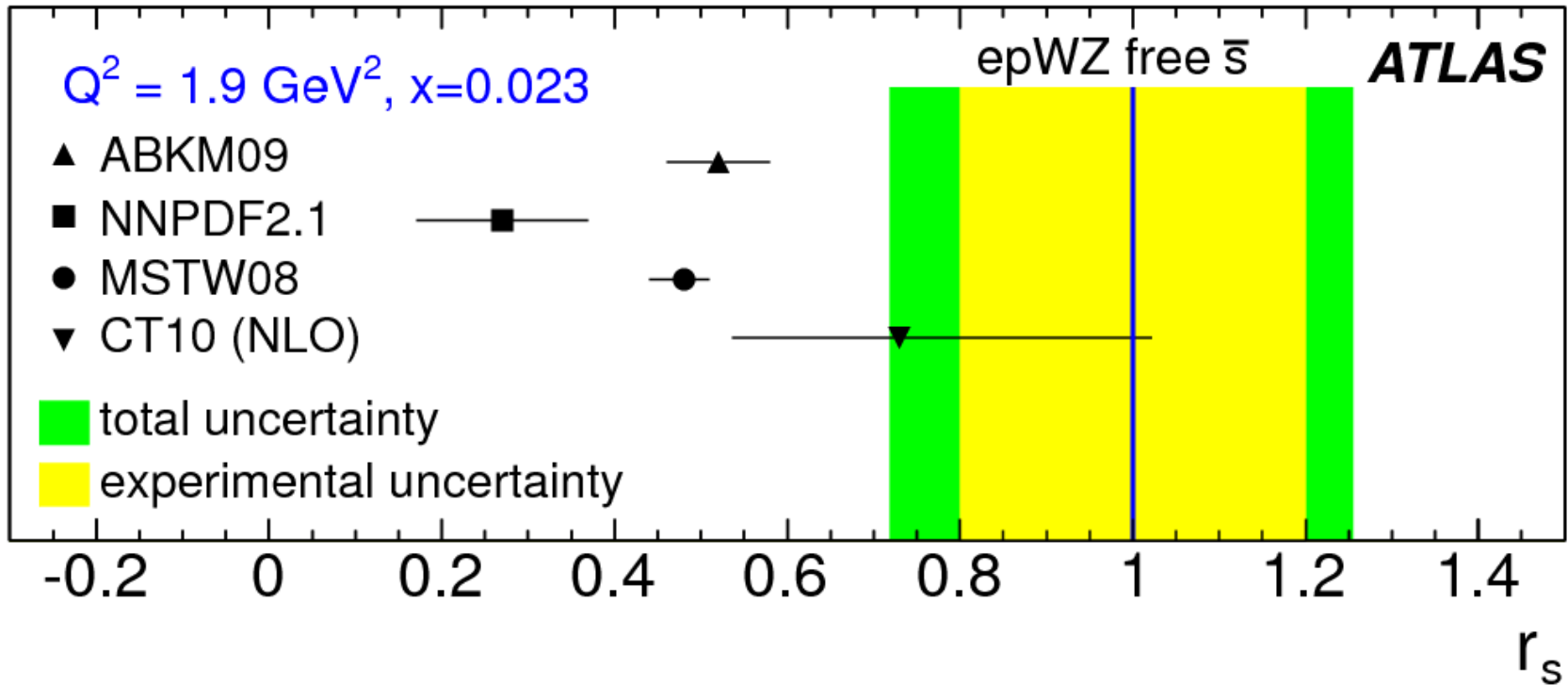
The lepton charge asymmetry from W-boson decays in bins of absolute pseudorapidity for the three different experiments ATLAS, CMS and LHCb.



# Reminder: The LHCb Detector



# Extraction of strange-quark density

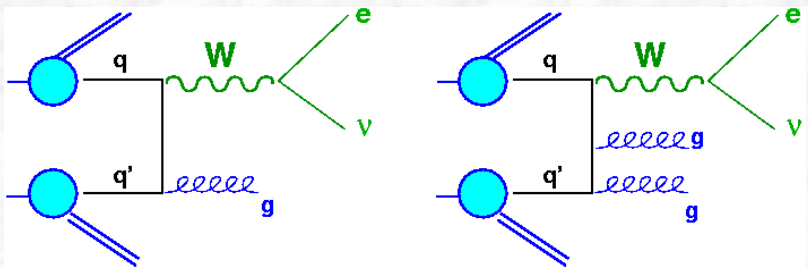


## Little is known about strange quark density

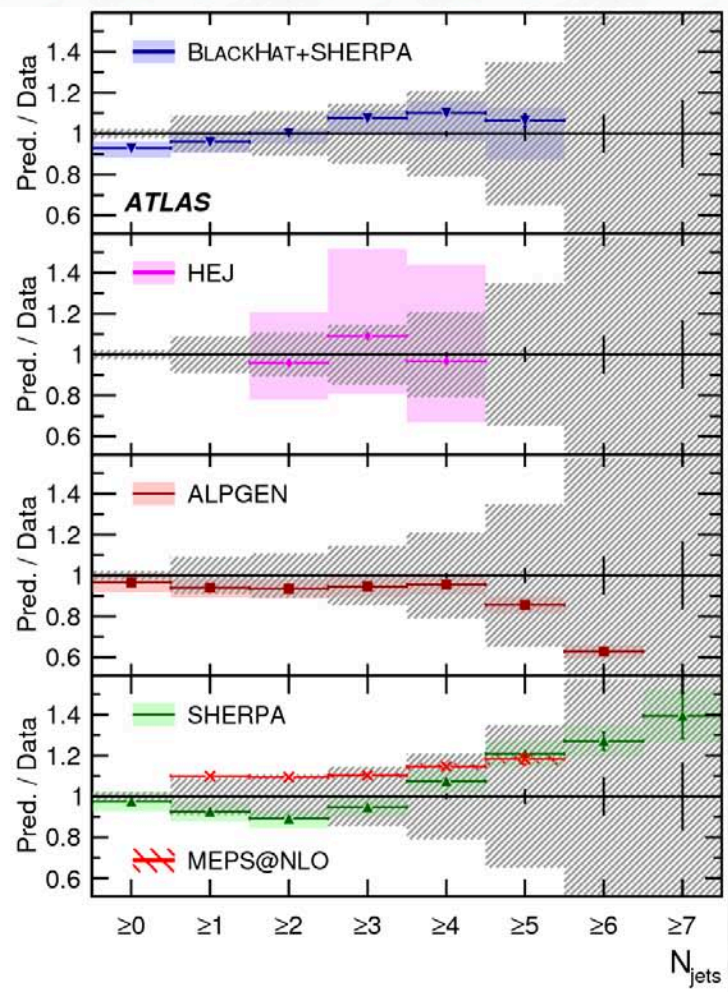
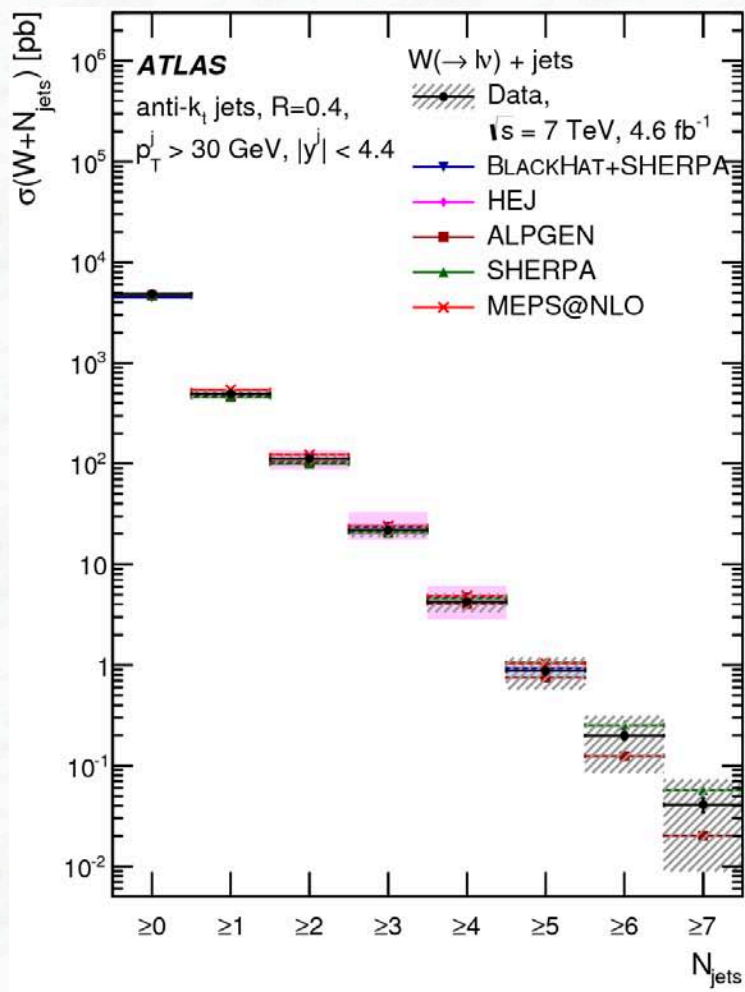
- Flavor SU(3) suggests equal ( $u, d, s$ ) in sea
- $m_s > m_{u,d}$ ; flavor SU(3) not exact symmetry
- $s$  suppression assumed in many PDFs because of  $s$  mass ( $r_s \sim 0.5$ )

$$r_s = 0.5(s + \bar{s})/\bar{d}$$

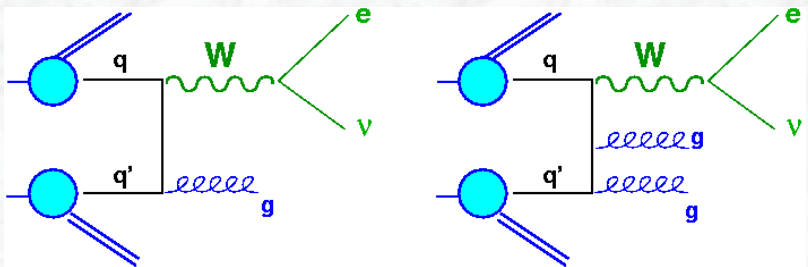
# Test of QCD in W/Z + jet production



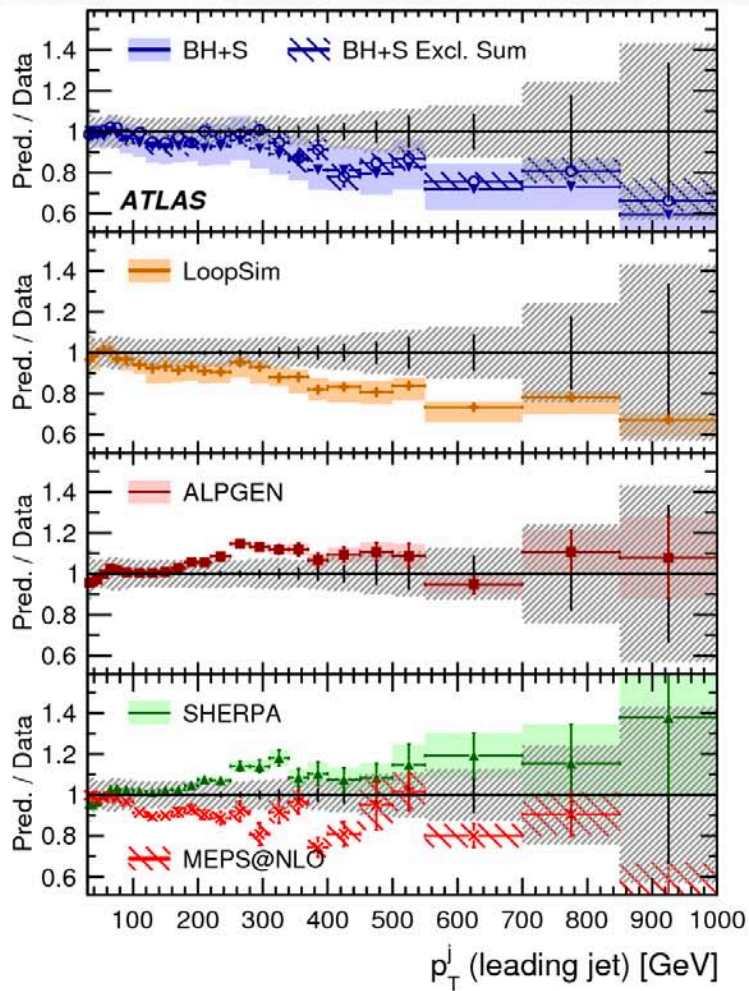
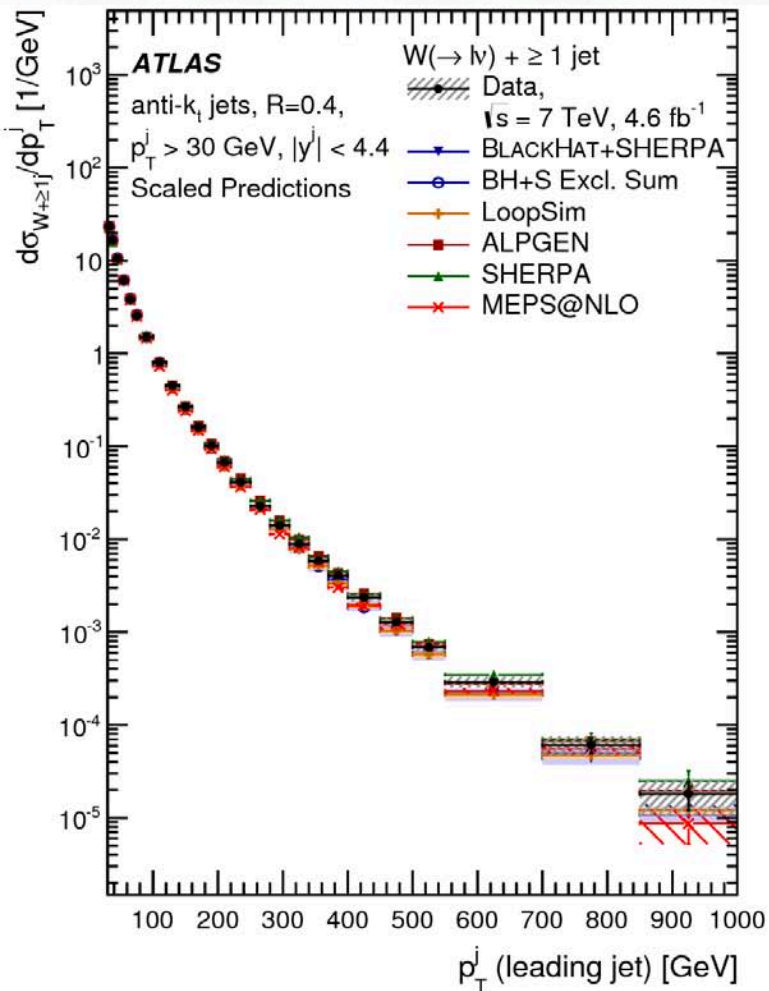
- Very large statistical power
- Agreement with several NLO and multi-lep predictions



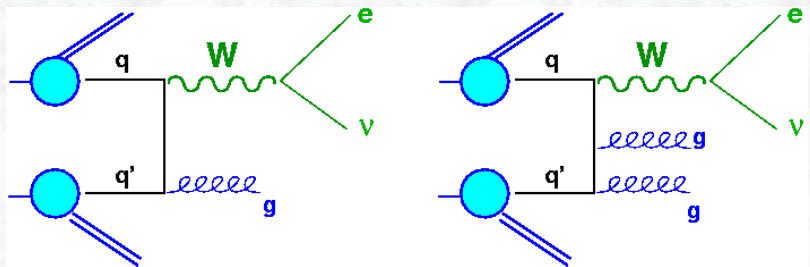
# Test of QCD in W/Z + jet production



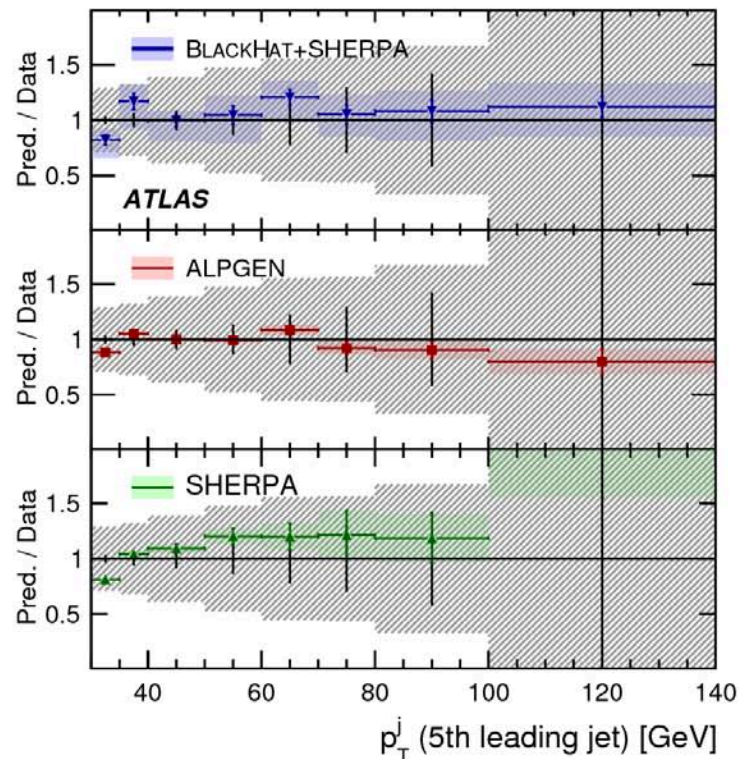
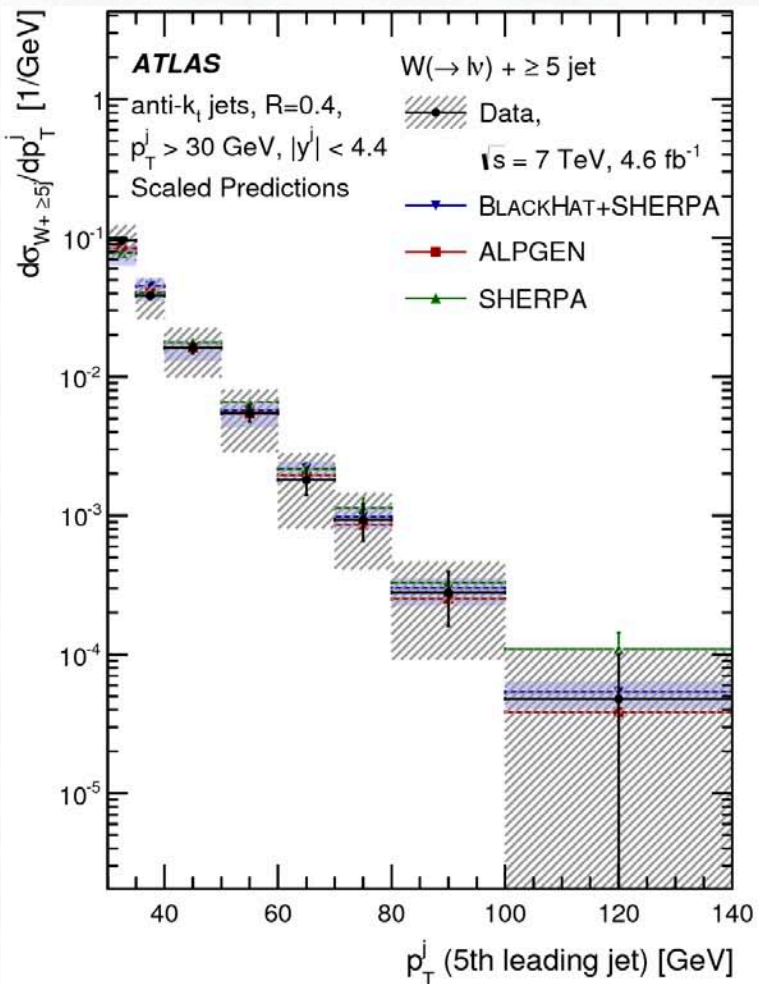
- Very large statistical power
- Agreement over many orders of magnitude



# Test of QCD in W/Z + jet production



- Very large statistical power
- Agreement over many orders of magnitude



Both jet rates and  $p_T$  spectra are well described by perturbative QCD calculations

# Summary of W/Z (+jets) measurements

## Vector Boson + X Cross Section Measurements

Status: March 2015

$\sigma^{\text{fid}}(\gamma+X)$  [ $|\eta^\gamma| < 1.37$ ]  
 - [ $1.52 < |\eta^\gamma| < 2.37$ ]

$\sigma^{\text{fid}}(Z \rightarrow ee, \mu\mu)$

- [ $n_{\text{jet}} \geq 1$ ]

- [ $n_{\text{jet}} \geq 2$ ]

- [ $n_{\text{jet}} \geq 3$ ]

- [ $n_{\text{jet}} \geq 4$ ]

- [ $n_{b\text{-jet}} \geq 1$ ]

- [ $n_{b\text{-jet}} \geq 2$ ]

-  $\sigma^{\text{fid}}(Z_{\text{jj}}^{\text{EWK}})$

$\sigma^{\text{fid}}(Z \rightarrow \tau\tau)$

$\sigma^{\text{fid}}(Z \rightarrow b\bar{b})$

$\sigma^{\text{fid}}(W \rightarrow e\nu, \mu\nu)$

- [ $n_{\text{jet}} \geq 1$ ]

- [ $n_{\text{jet}} \geq 2$ ]

- [ $n_{\text{jet}} \geq 3$ ]

- [ $n_{\text{jet}} \geq 4$ ]

- [ $n_{\text{jet}} \geq 1, n_{b\text{-jet}} = 1$ ]

- [ $n_{\text{jet}} \geq 2, n_{b\text{-jet}} = 1$ ]

$\sigma^{\text{fid}}(W \rightarrow e\nu, \mu\nu) / \sigma^{\text{fid}}(Z \rightarrow ee, \mu\mu)$

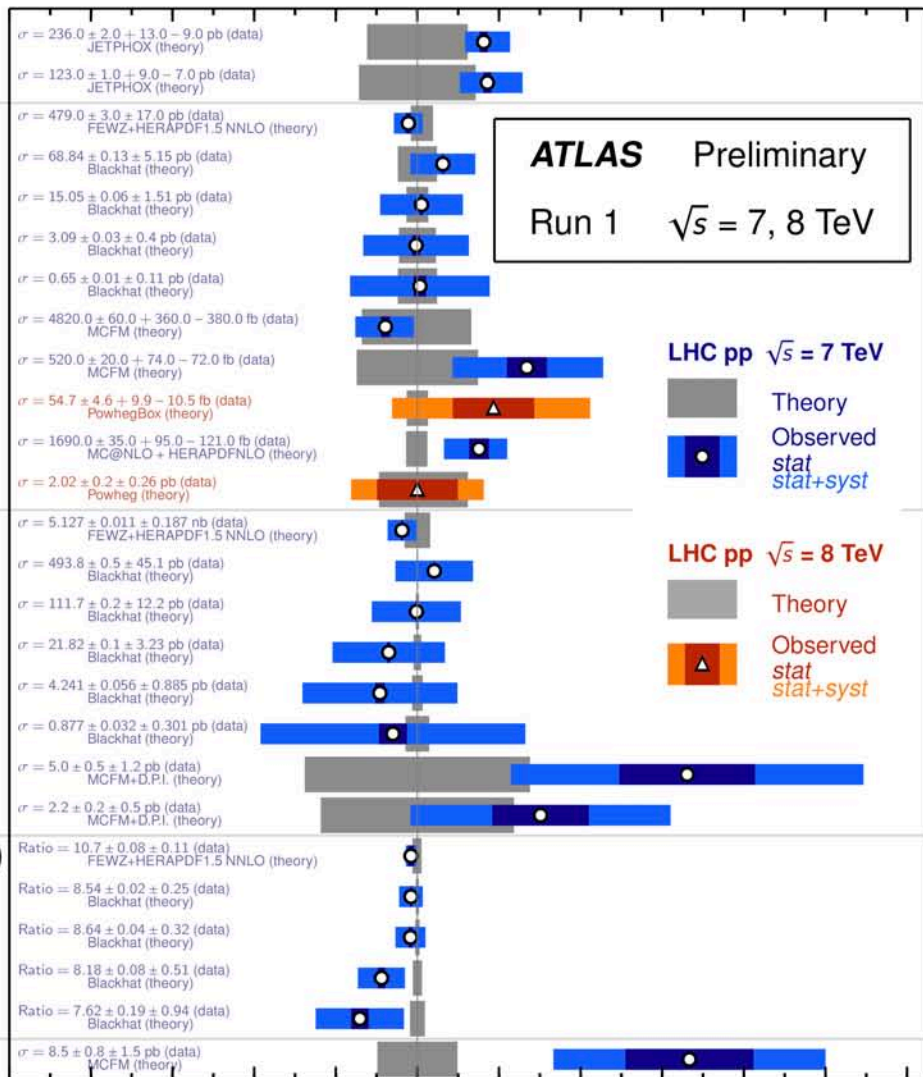
- [ $n_{\text{jet}} \geq 1$ ]

- [ $n_{\text{jet}} \geq 2$ ]

- [ $n_{\text{jet}} \geq 3$ ]

- [ $n_{\text{jet}} \geq 4$ ]

$\sigma^{\text{fid}}(W+Z \rightarrow q\bar{q})$



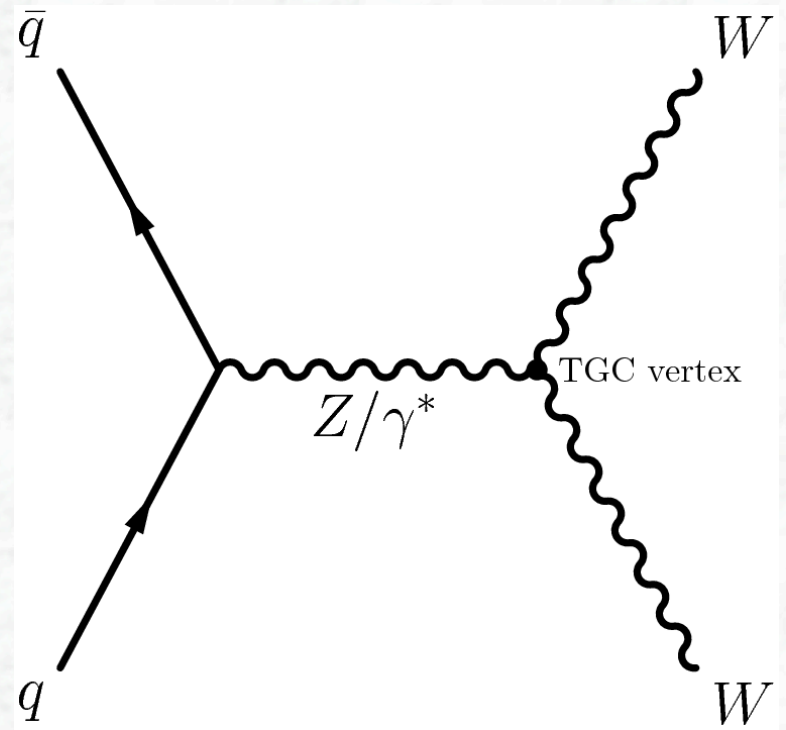
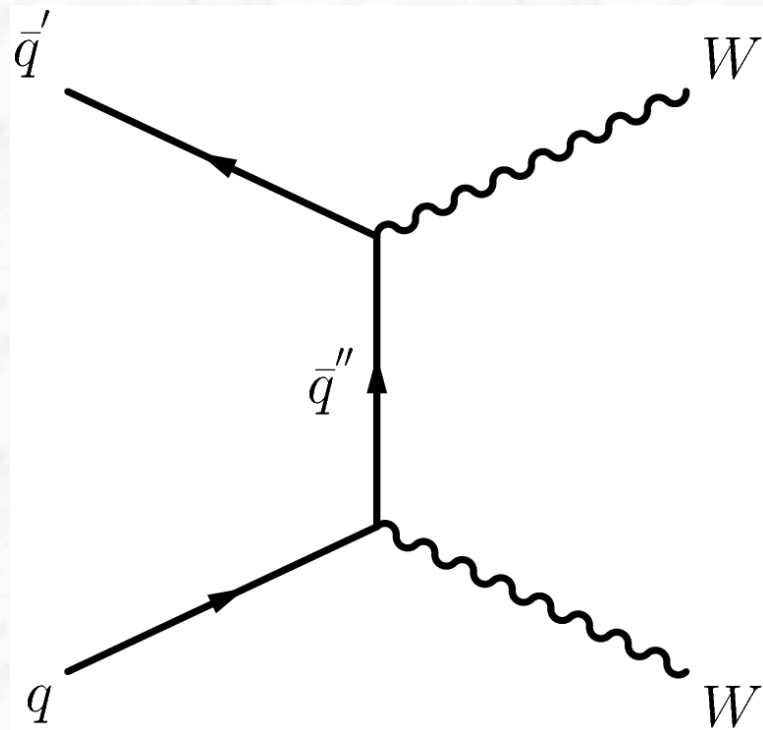
$\int \mathcal{L} dt$   
 [ $\text{fb}^{-1}$ ]

Reference

4.6	PRD 89, 052004 (2014)
4.6	PRD 89, 052004 (2014)
0.035	PRD 85, 072004 (2012)
4.6	JHEP 07, 032 (2013)
4.6	JHEP 07, 032 (2013)
4.6	JHEP 07, 032 (2013)
4.6	JHEP 10, 141, (2014)
4.6	JHEP 10, 141, (2014)
20.3	JHEP 04, 031 (2014)
4.6	arXiv:1407.0573 [hep-ex]
19.5	PLB 738, 25-43, (2014)
0.035	PRD 85, 072004 (2012)
4.6	arXiv:1409.8639 [hep-ex]
4.6	arXiv:1409.8639 [hep-ex]
4.6	arXiv:1409.8639 [hep-ex]
4.6	arXiv:1409.8639 [hep-ex]
4.6	JHEP 06, 084 (2013)
4.6	JHEP 06, 084 (2013)
0.035	PRD 85, 072004 (2012)
4.6	Eur. Phys. J. C 74: 3168 (2014)
4.6	Eur. Phys. J. C 74: 3168 (2014)
4.6	Eur. Phys. J. C 74: 3168 (2014)
4.6	Eur. Phys. J. C 74: 3168 (2014)
4.6	New J. Phys. 16, 113013 (2014)

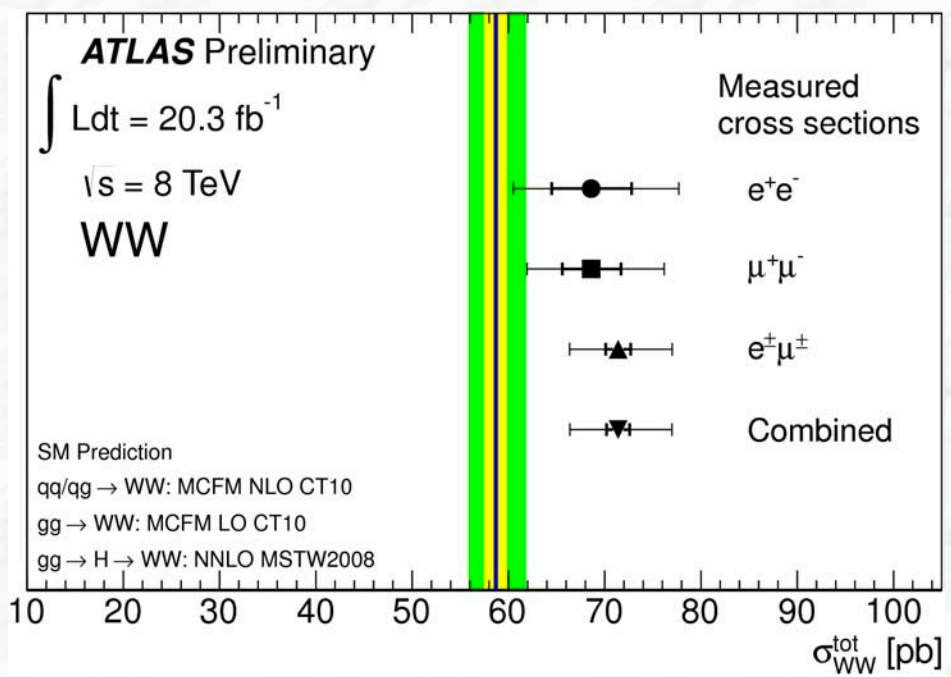
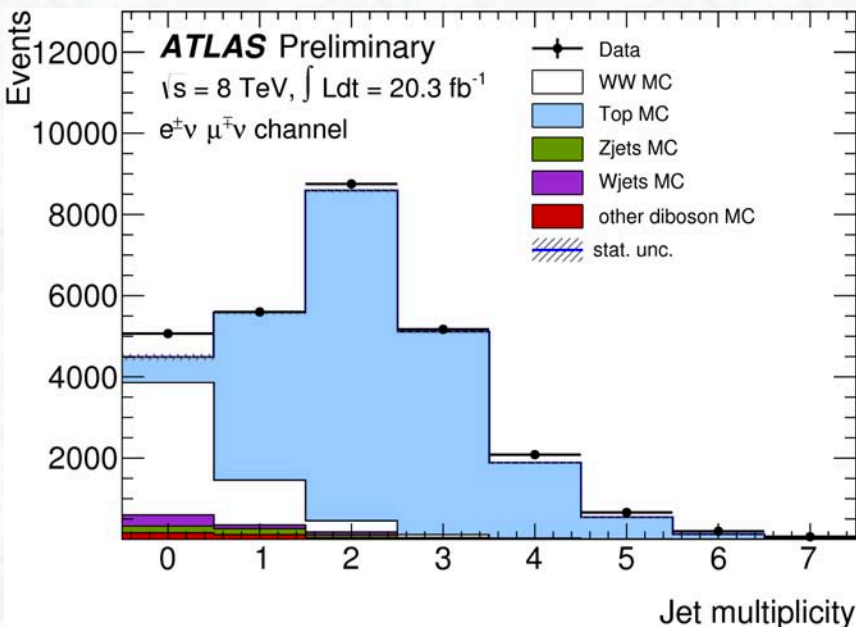
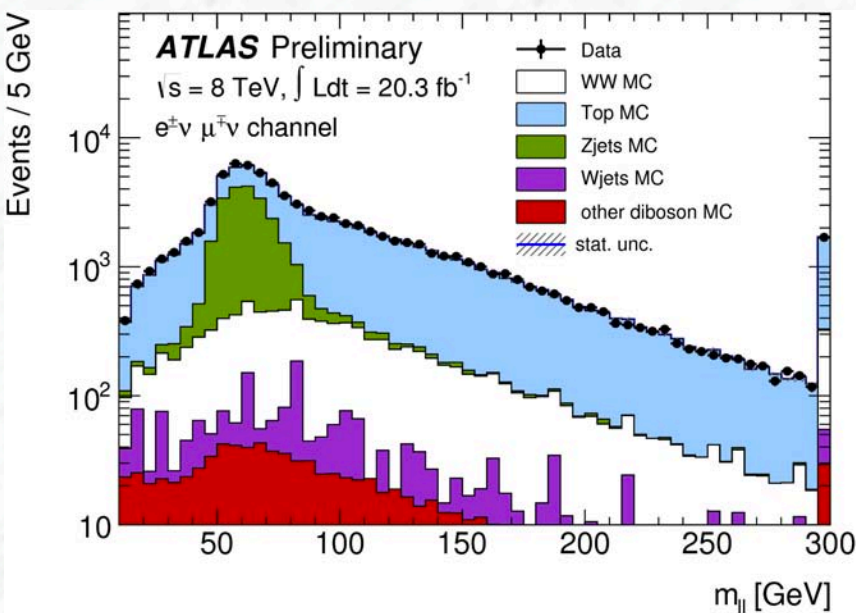
0.0 0.2 0.4 0.6 0.8 1.0 1.2 1.4 1.6 1.8 2.0 2.2  
 observed/theory

## 6.5 Di-boson measurements



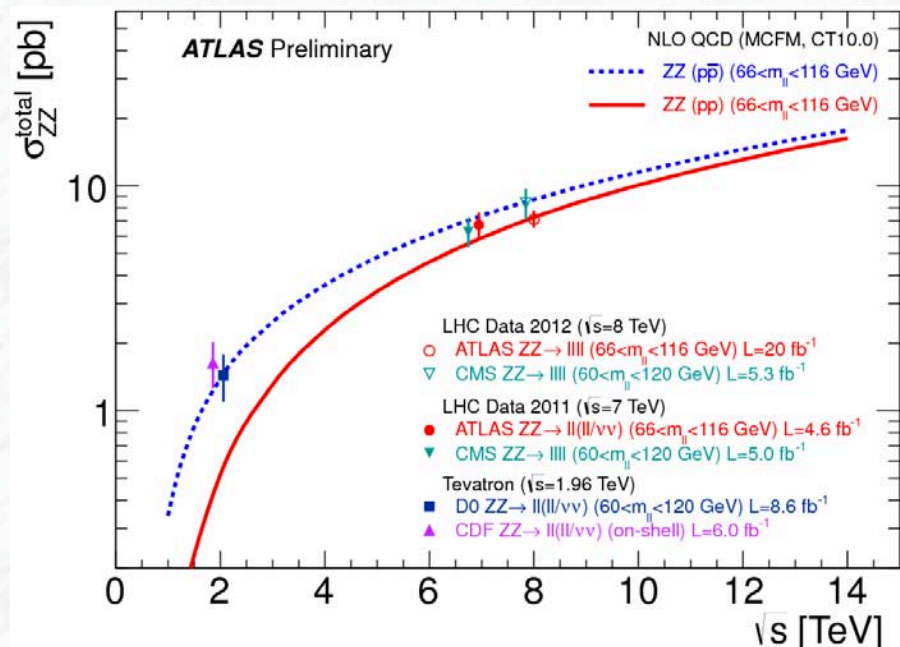
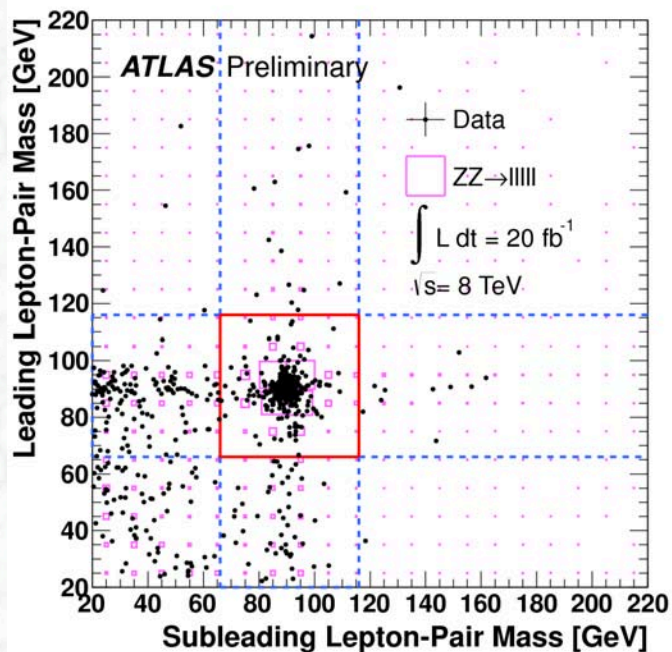
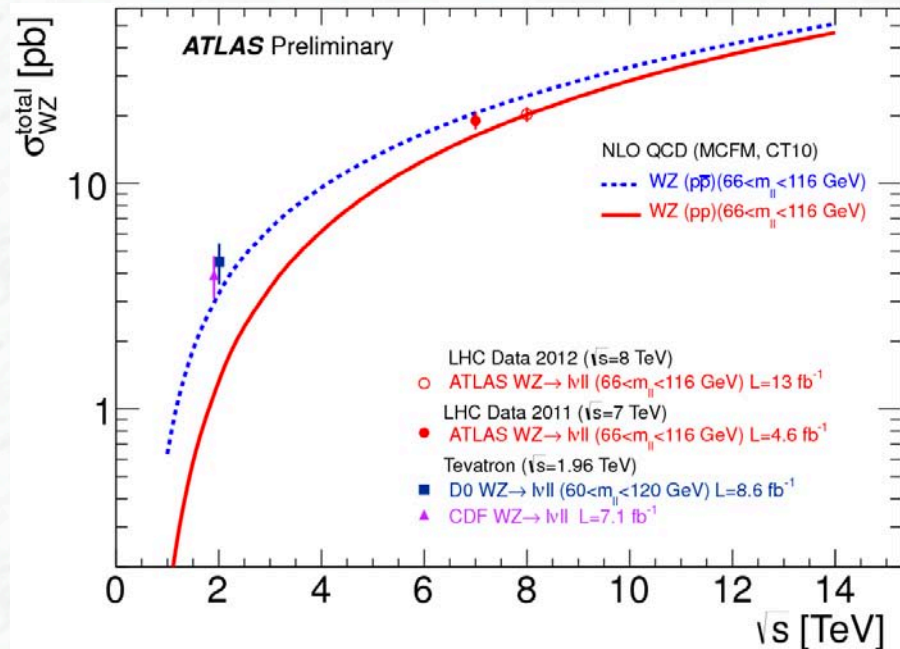
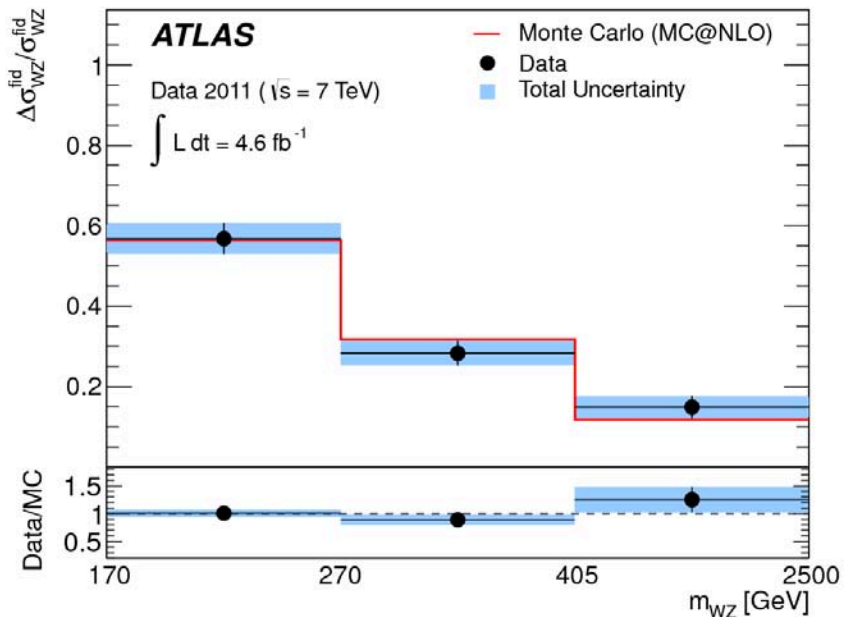
# WW Production

- Cross section measurement only in 0-jet bin
- Result is factor 1.21 higher than NLO prediction; but not significant
- Newer calculations seem to (partially) reconcile this
- Measure also 1-jet bin in the future

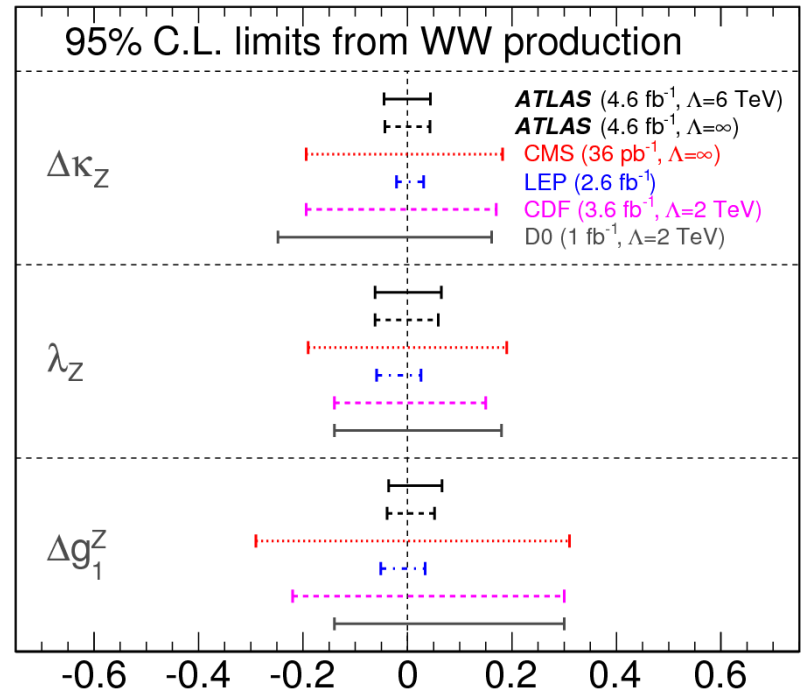
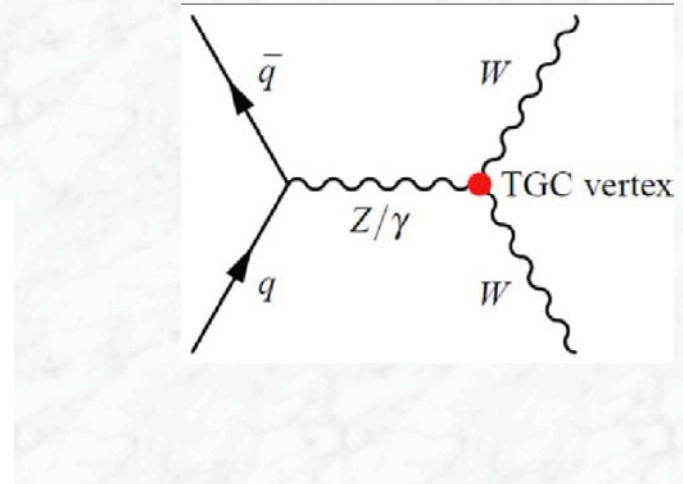
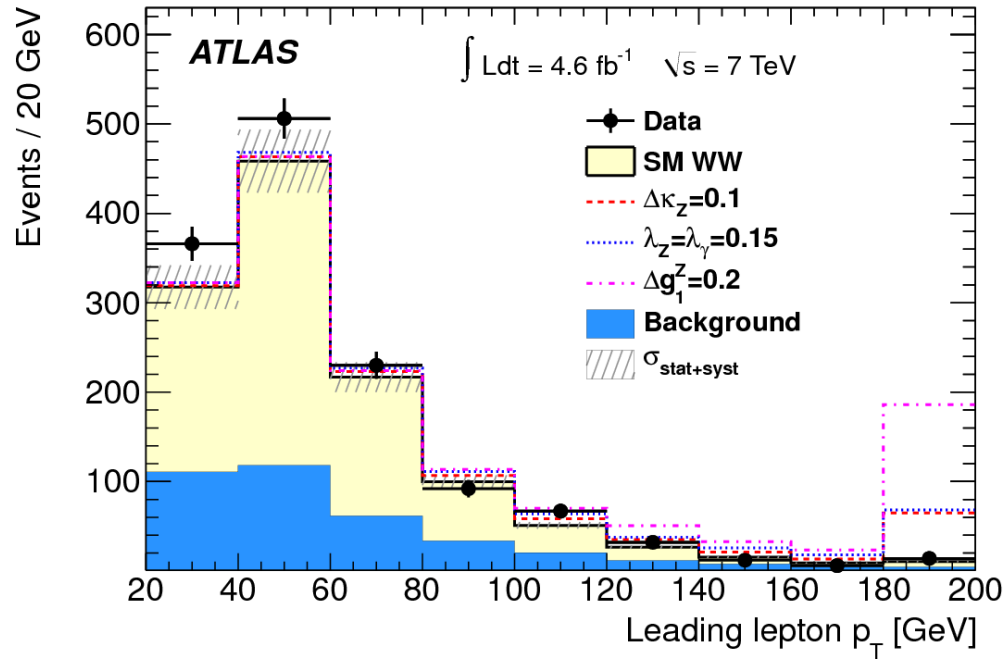




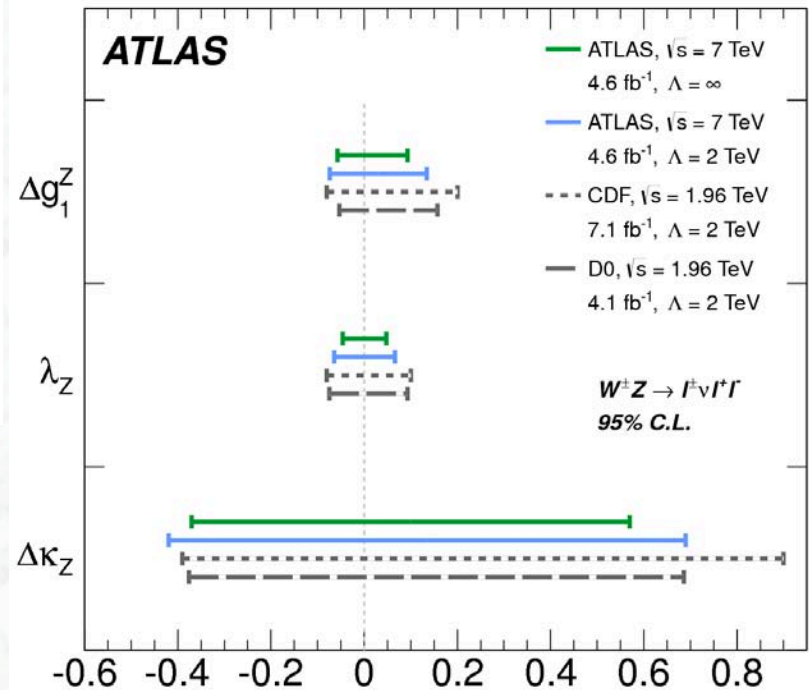
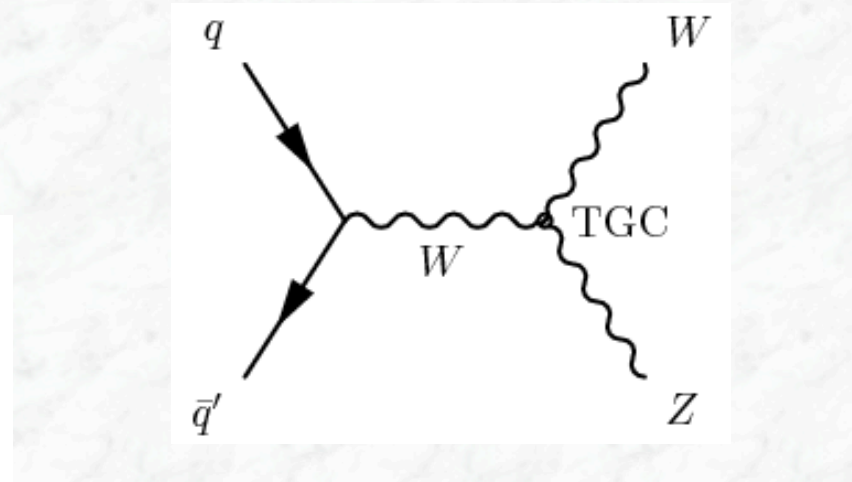
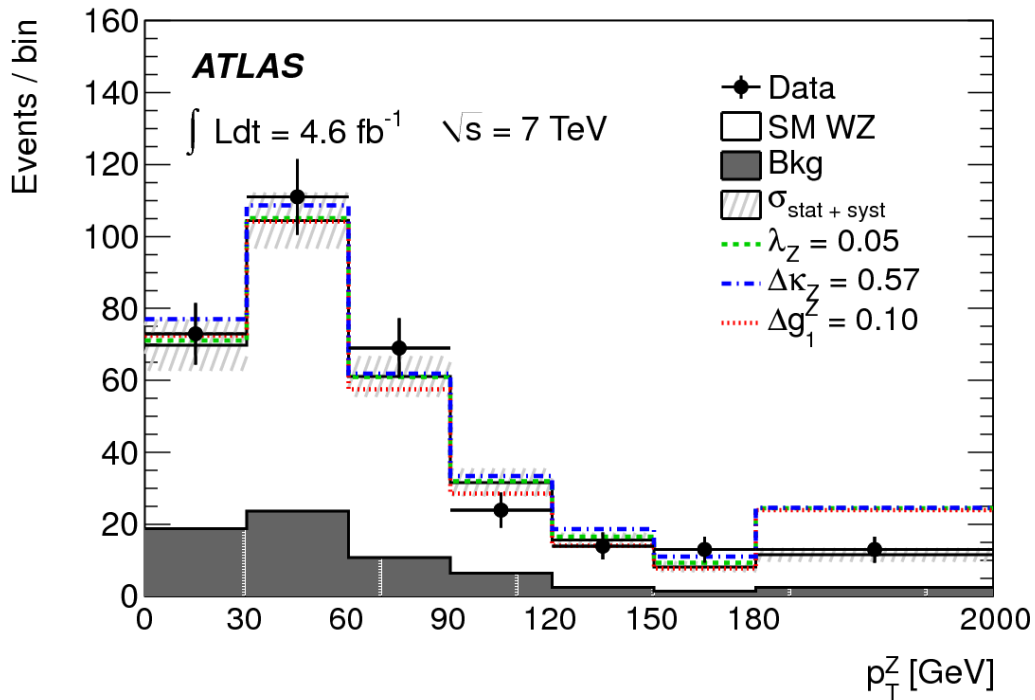
# WZ and ZZ cross-sections



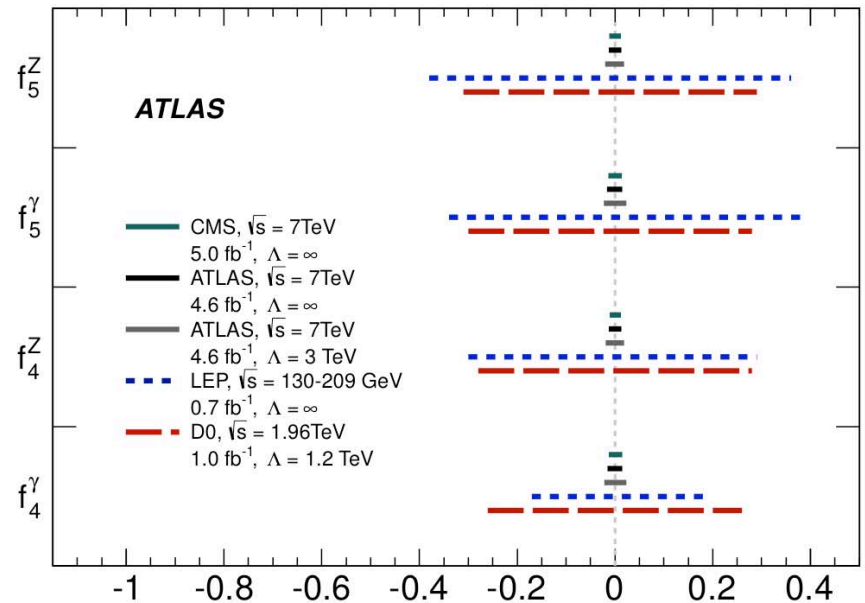
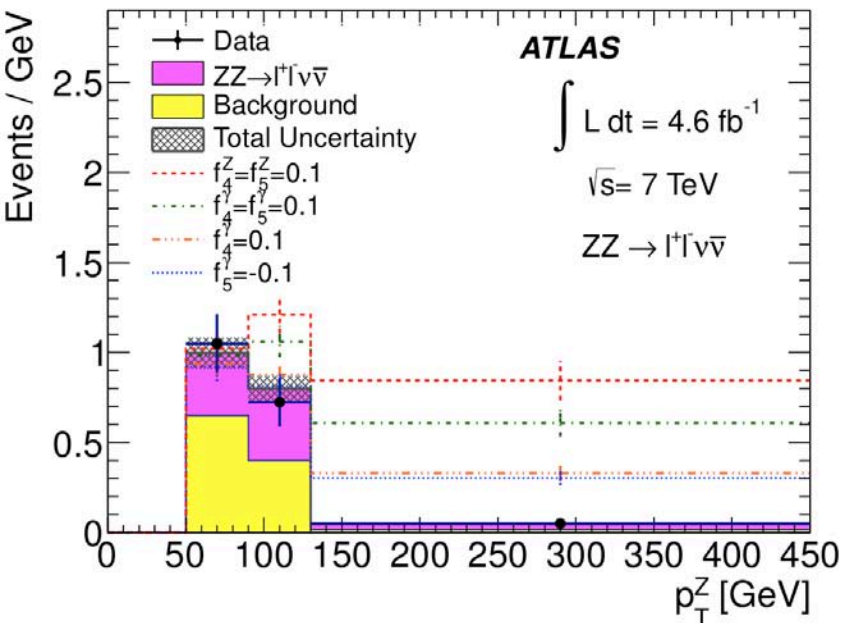
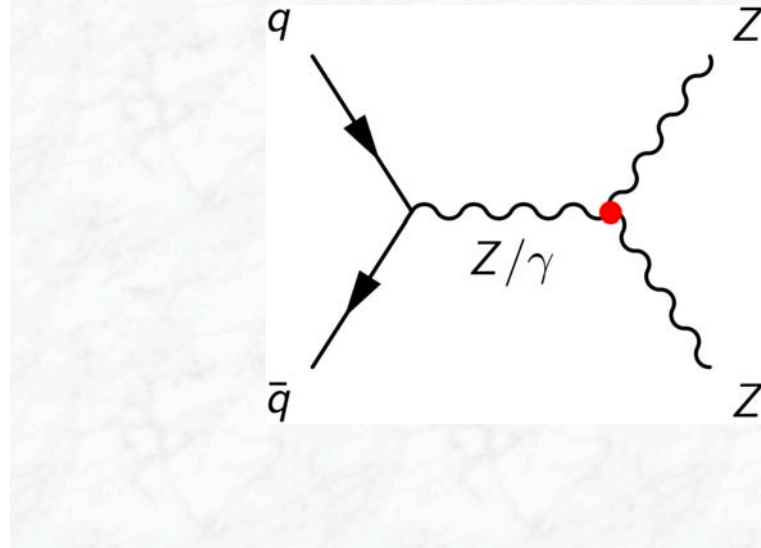
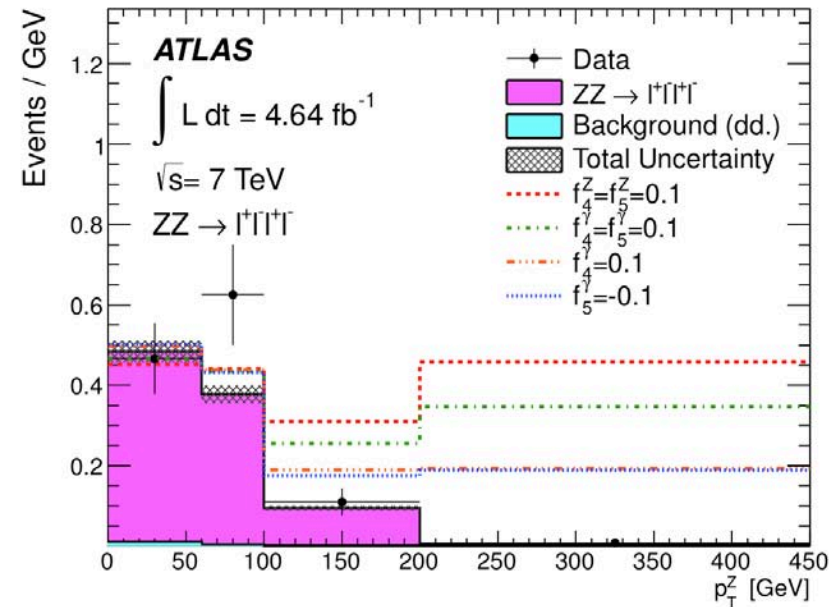
# Limits on anomalous couplings: **WW**



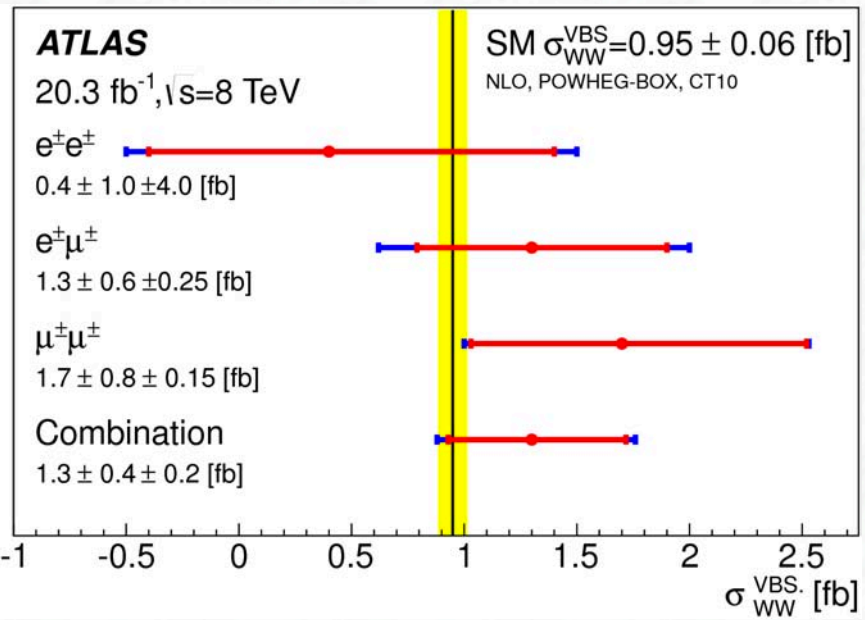
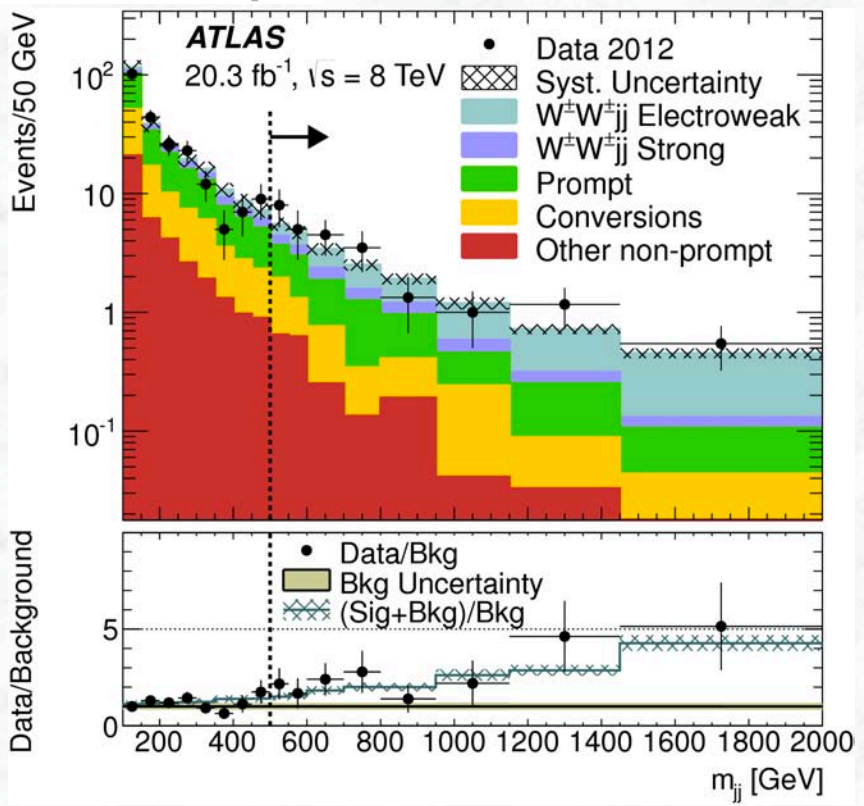
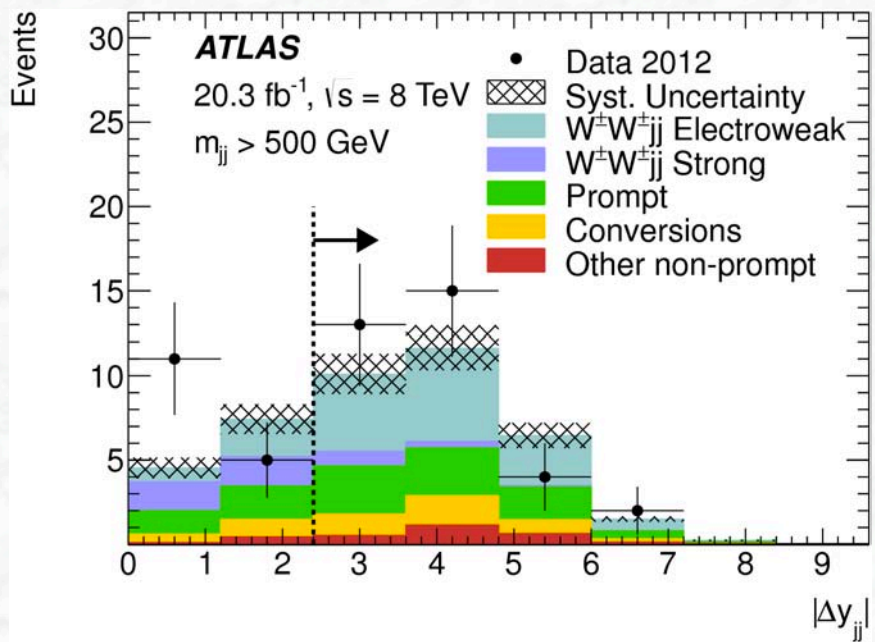
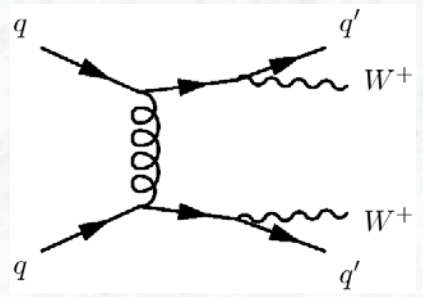
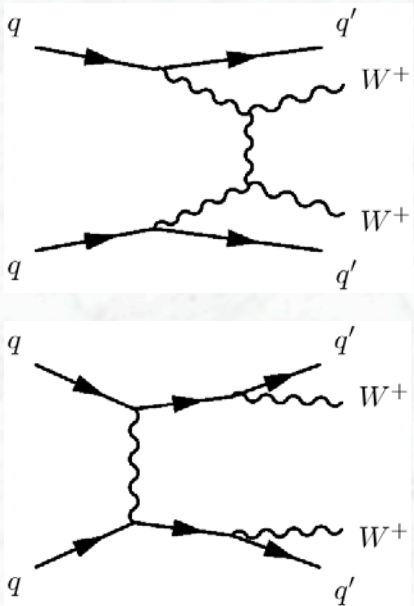
# Limits on anomalous couplings: **WZ**



# Limits on anomalous couplings: ZZ



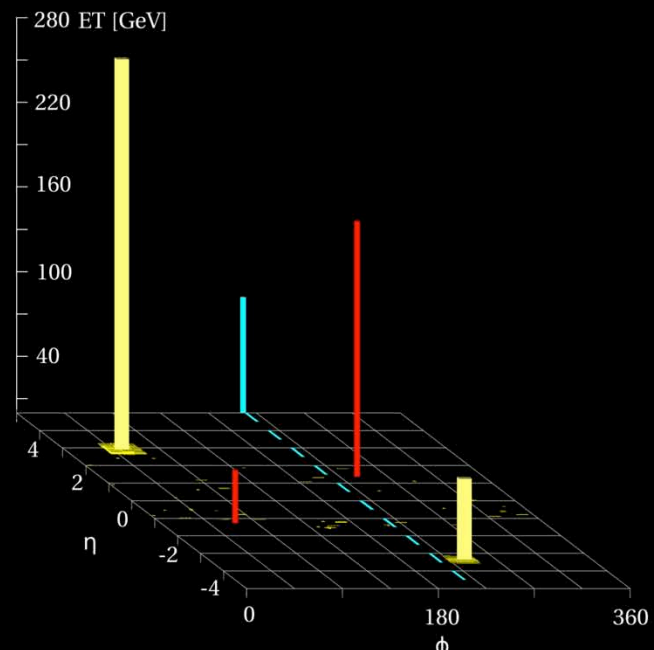
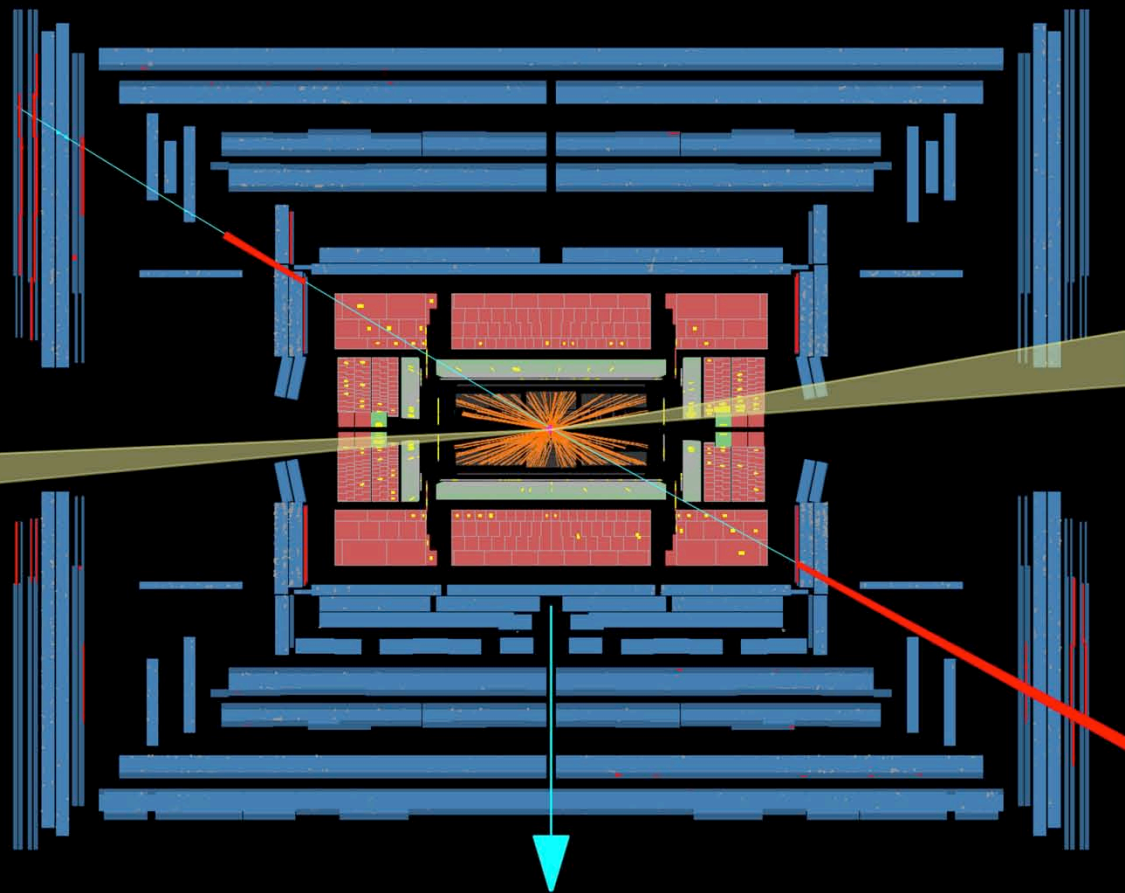
# Electroweak production of $W^+W^+jj$



# $\mu^+\mu^+jj$ Candidate Event

$m_{jj} = 2800$  GeV

$|\Delta y_{jj}| = 6.3$

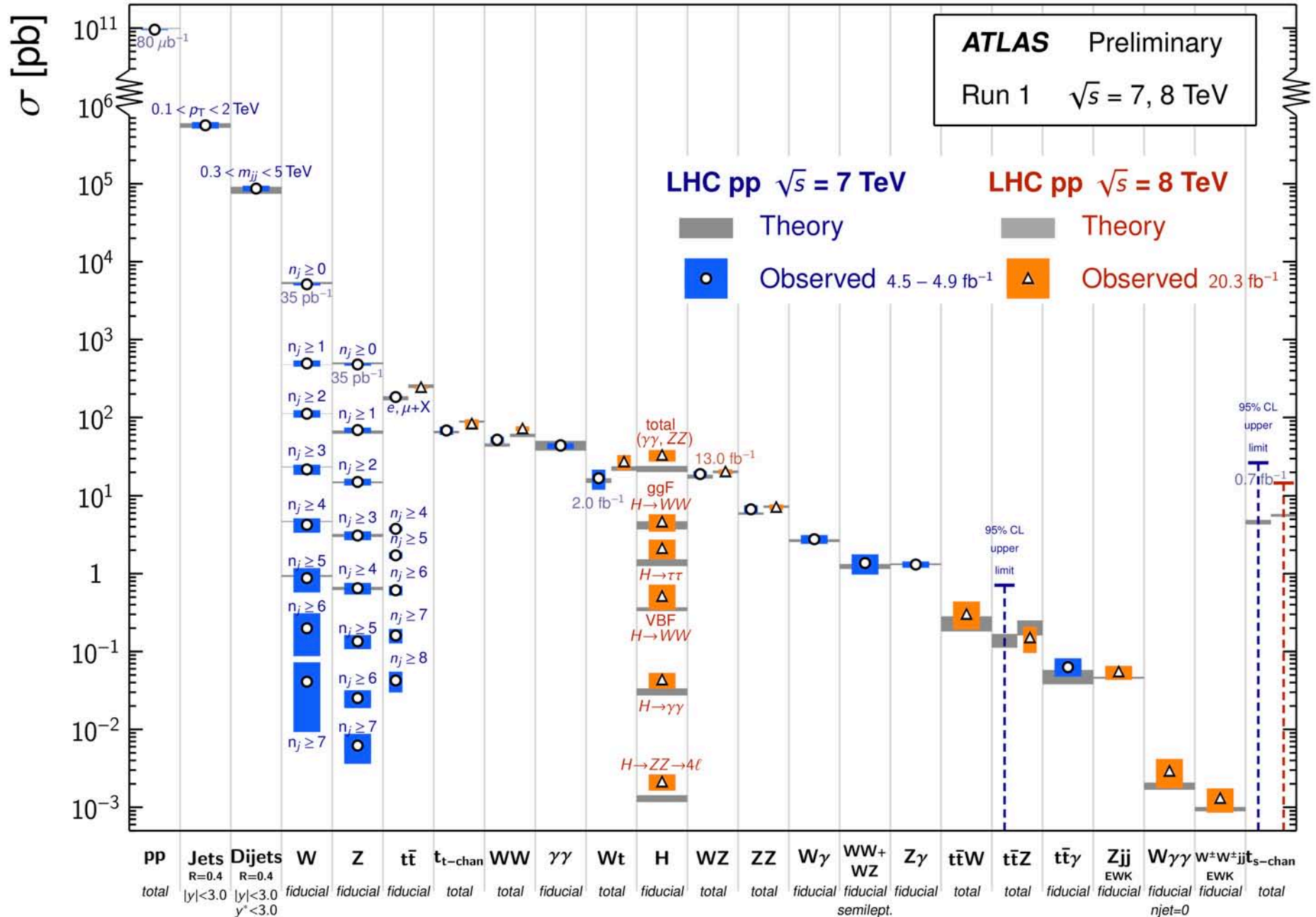


Run Number: 207490, Event Number: 33152138

Date: 2012-07-26 04:16:35 UTC

# Standard Model Production Cross Section Measurements

Status: March 2015



## 6.6 W mass measurement

Major contributions: LEP-II, direct mass reconstruction

Hadron collider: Tevatron and LHC (in the future)



# Precision measurements of $m_W$ and $m_{\text{top}}$

## Motivation:

W mass and top quark mass are **fundamental parameters** of the Standard Model;  
 The standard theory provides well defined **relations between  $m_W$ ,  $m_{\text{top}}$  and  $m_H$**

Electromagnetic constant  
 measured in atomic transitions,  
 $e^+e^-$  machines, etc.

$$m_W = \left( \frac{\pi \alpha_{EM}}{\sqrt{2} G_F} \right)^{1/2} \frac{1}{\sin \theta_W \sqrt{1 - \Delta r}}$$

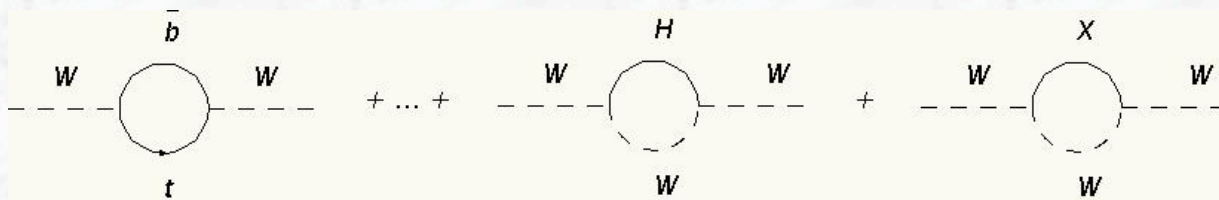
Fermi constant  
 measured in muon  
 decay

weak mixing angle  
 measured at  
 LEP/SLC

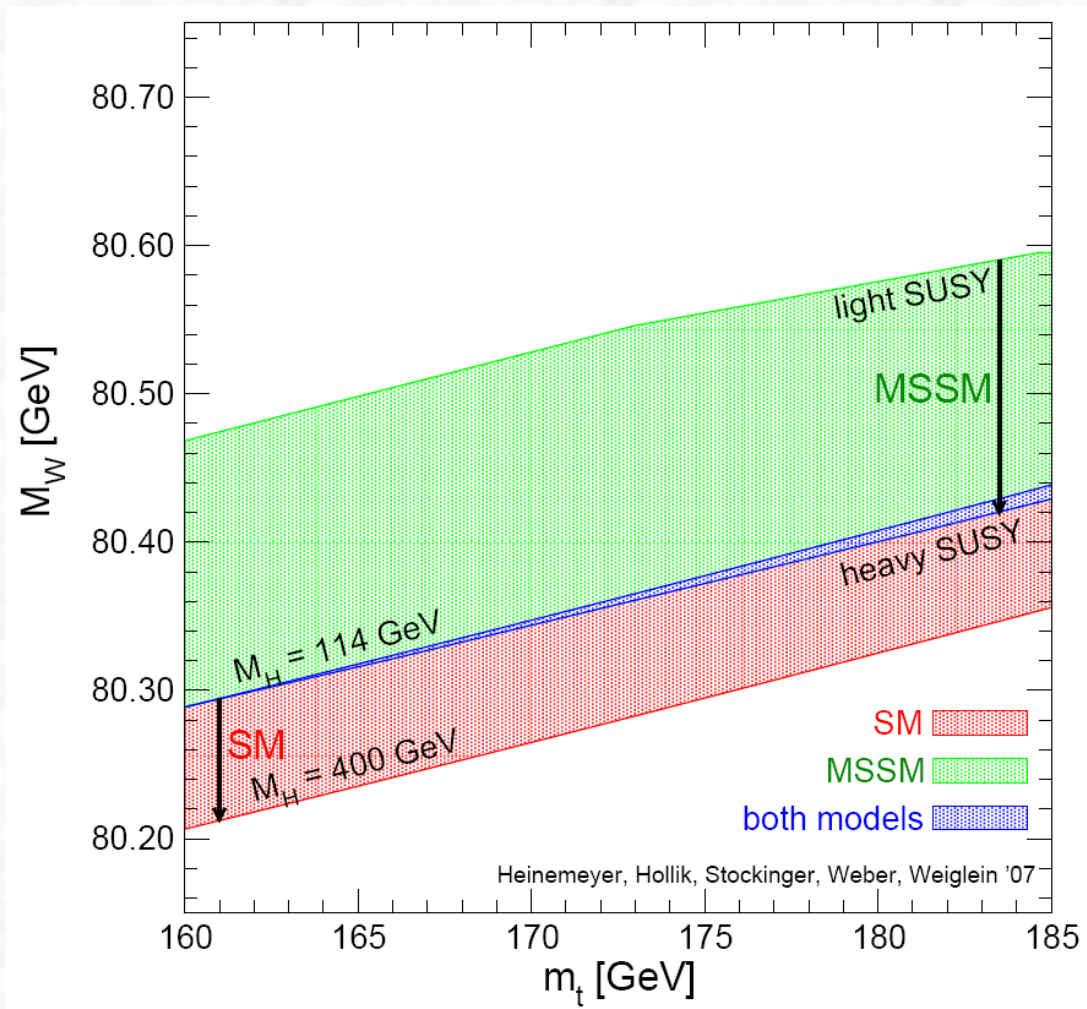
radiative corrections  
 $\Delta r \sim f(m_{\text{top}}^2, \log m_H)$   
 $\Delta r \approx 3\%$

$G_F, \alpha_{EM}, \sin \theta_W$   
 are known with high precision

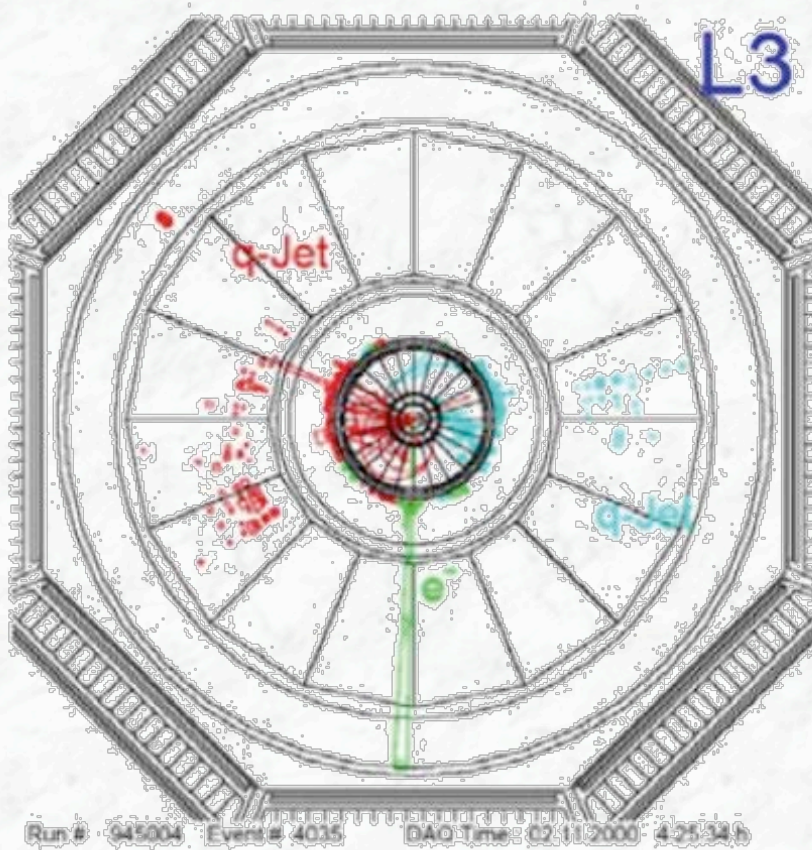
Precise measurements of the  
 W mass and the top-quark  
 mass constrain the Higgs-  
 boson mass  
 (and/or the theory,  
 radiative corrections)



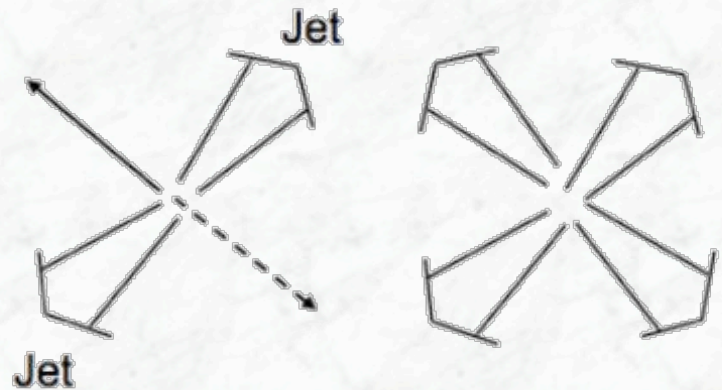
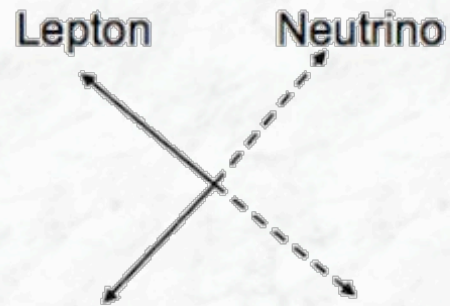
# Relation between $m_W$ , $m_t$ , and $m_H$



# W bosons at LEP – II

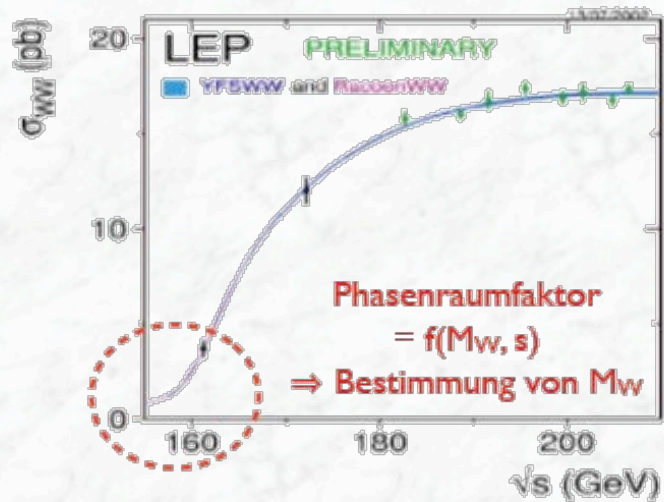


$WW \rightarrow$   $\left\{ \begin{array}{l} qq\ell\nu \text{ 44\%} \\ qqqq \text{ 45\%} \\ \ell\nu\ell\nu \text{ 11\%} \end{array} \right.$

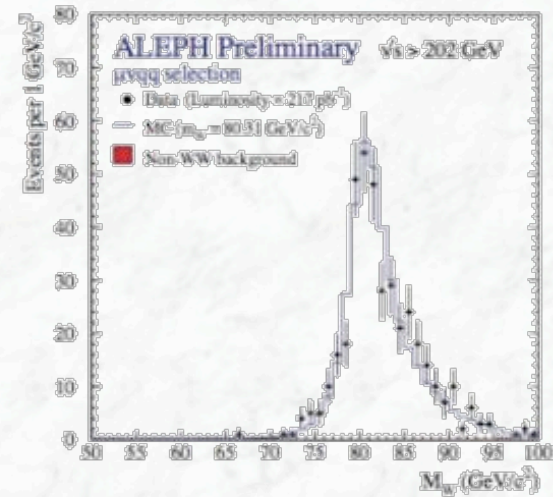


# W mass measurement

(1) Messung des WQs an der WW-Produktionsschwelle

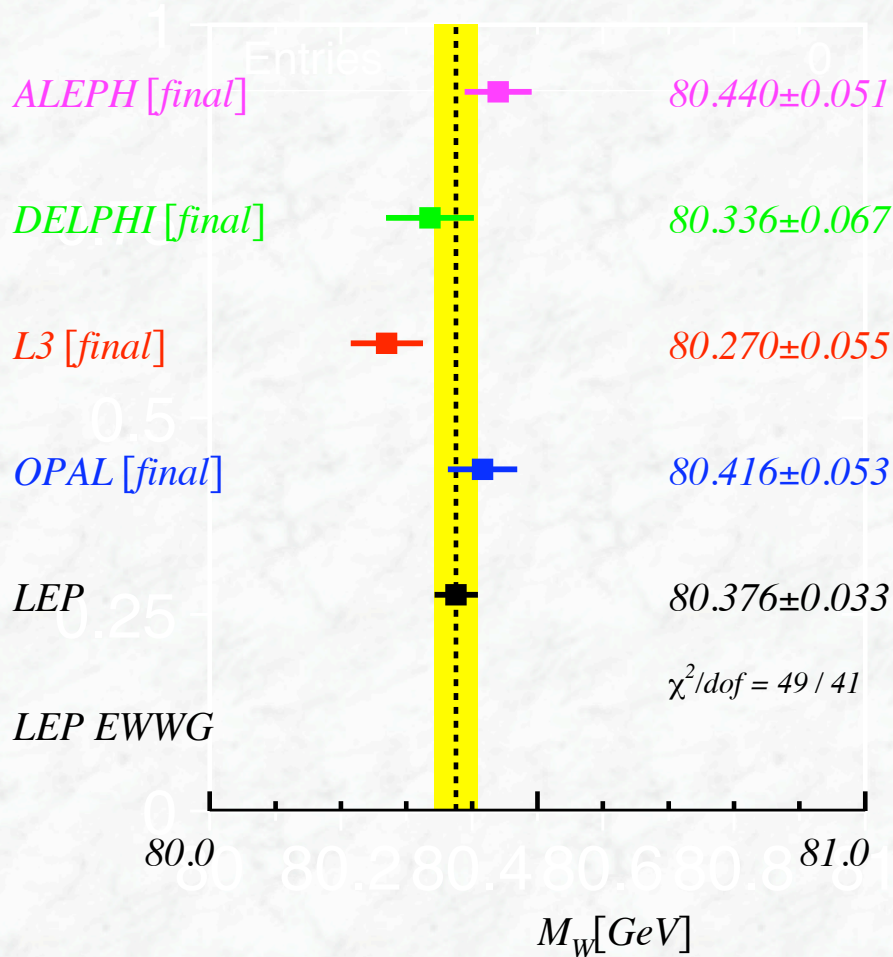


(2) LEP-II: Direkte Rekonstruktion der invarianten Masse des W-Bosons:



# Results from W mass measurements at LEP-II

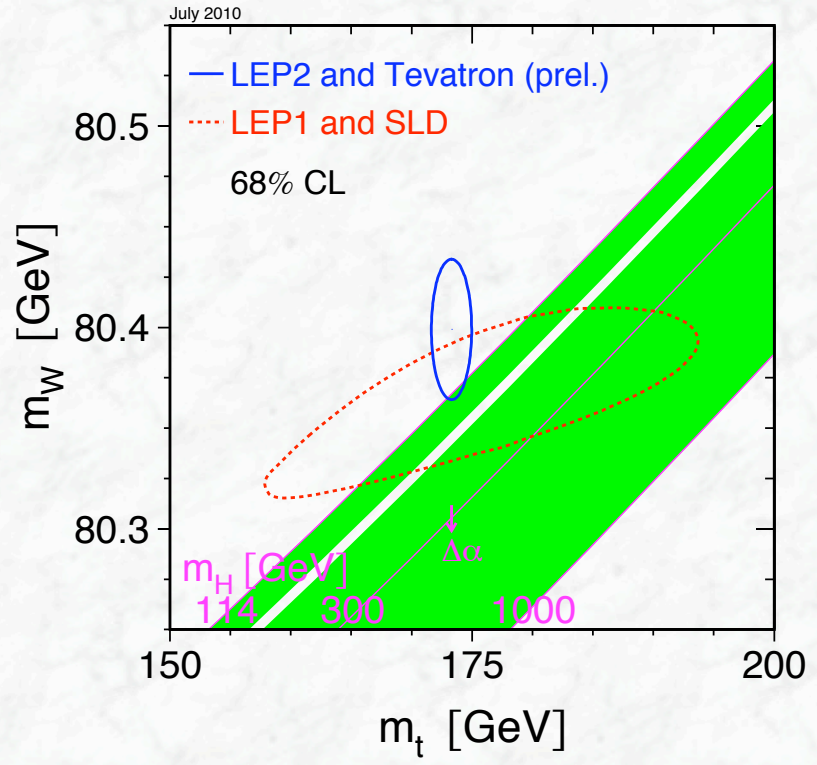
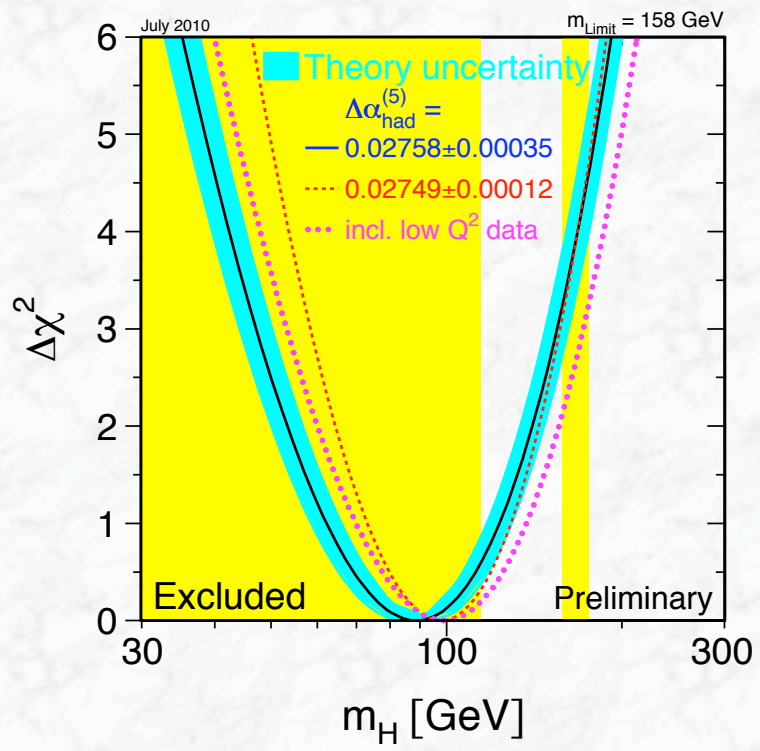
Summer 2006 - LEP Preliminary



- Results from all four LEP experiments are consistent
- Statistical error is dominant
- Total precision from LEP-II

$$\Delta m_W = \pm 33 \text{ MeV}$$

# Results of electroweak precision tests at LEP (cont.)



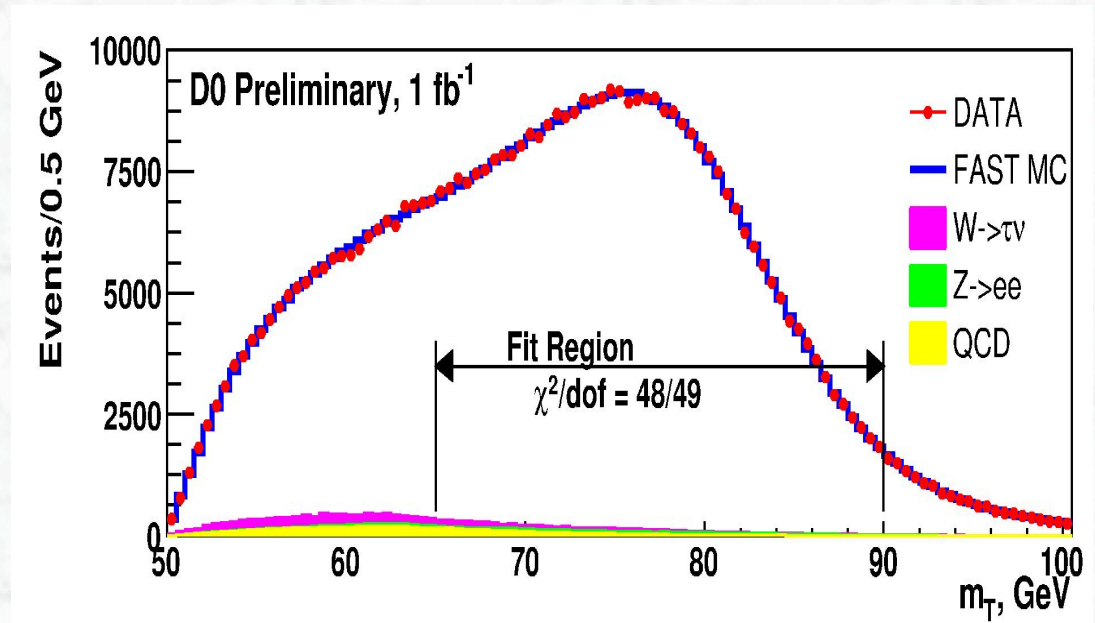
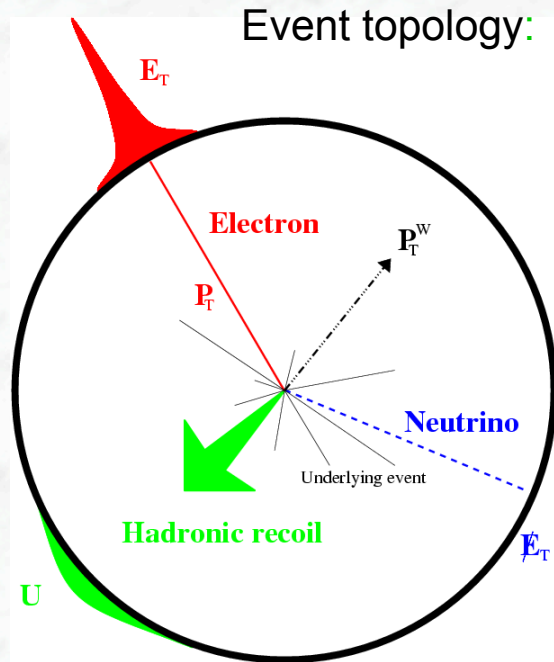
- Radiative corrections (loop, quantum corrections) can be used to constrain yet unobserved particles (however, sensitivity to  $m_H$  only through log terms)
- Main reason for continued precision improvements in  $m_t$ ,  $m_W$

What can hadron collider contribute ?

How can  $W$  mass be measured at a hadron collider ?



# Technique used for W mass measurement at hadron colliders:



Observables:  $P_T(e)$ ,  $P_T(\text{had})$

$$\Rightarrow P_T(\nu) = - ( P_T(e) + P_T(\text{had}) )$$

long. component cannot be

$$\Rightarrow M_W^T = \sqrt{2 \cdot P_T^e \cdot P_T^{\nu} \cdot (1 - \cos \Delta\phi^{e,\nu})}$$

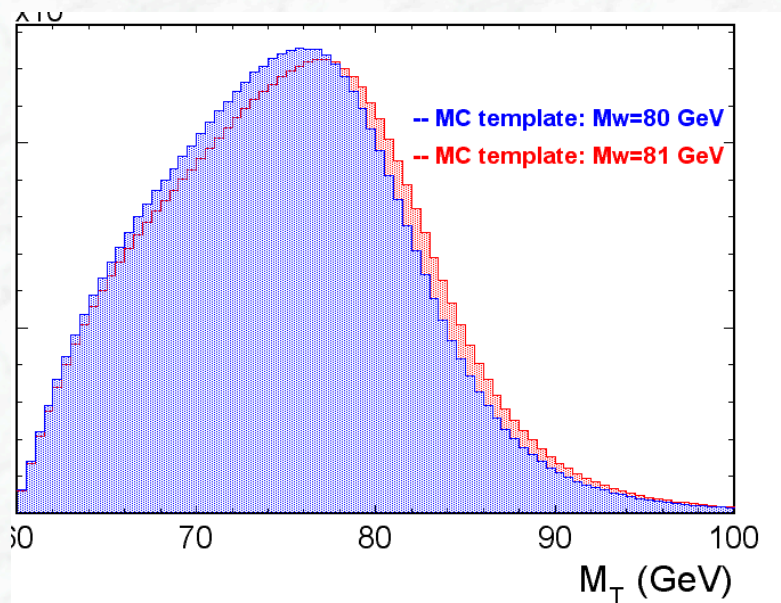
measured

In general the **transverse mass**  $M_T$  is used for the determination of the W mass

(smallest systematic uncertainty).



Shape of the transverse mass distribution is sensitive to  $m_W$ , the measured distribution is fitted with Monte Carlo predictions, where  $m_W$  is a parameter



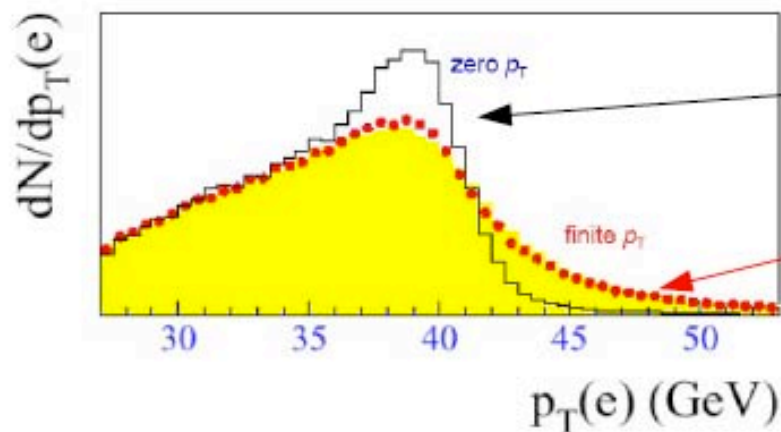
### Main uncertainties:

Ability of the Monte Carlo to reproduce real life:

- Detector performance  
(energy resolution, energy scale, ....)
- Physics: production model  
 $p_T(W)$ ,  $\Gamma_W$ , .....
- Backgrounds

In principle any distribution that is sensitive to  $m_W$  can be used for the measurement;

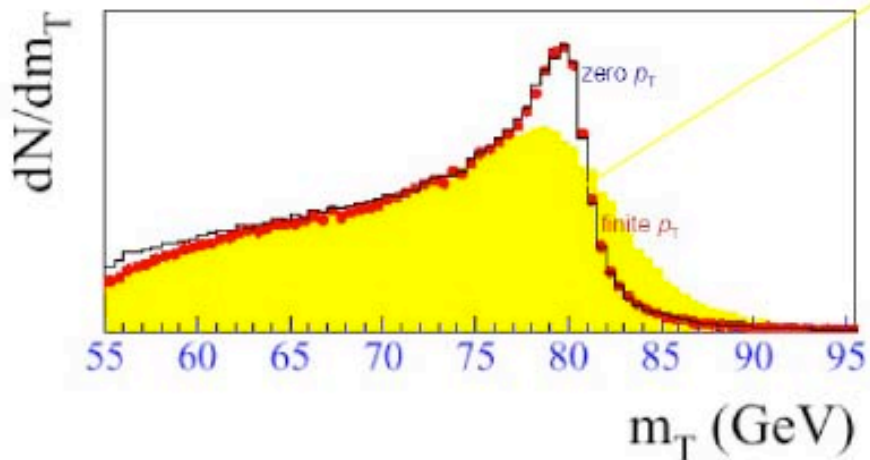
Systematic uncertainties are different for the various observables.



**True distribution**

**Including  $p_T(W)$  effects**

**Including detector effects**

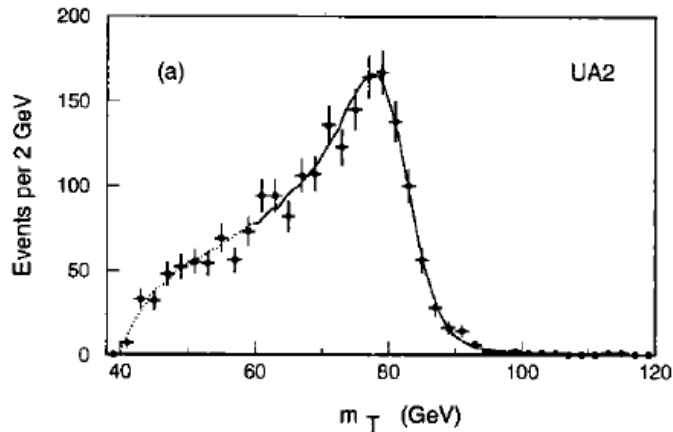


**$p_T(e)$  not sensitive to detector effects, requires  $p_T(W)$  knowledge**

**Transverse mass less sensitive to  $p_T(W)$ , requires good modeling of missing  $E_T$**

# W mass measurements

## The beginning

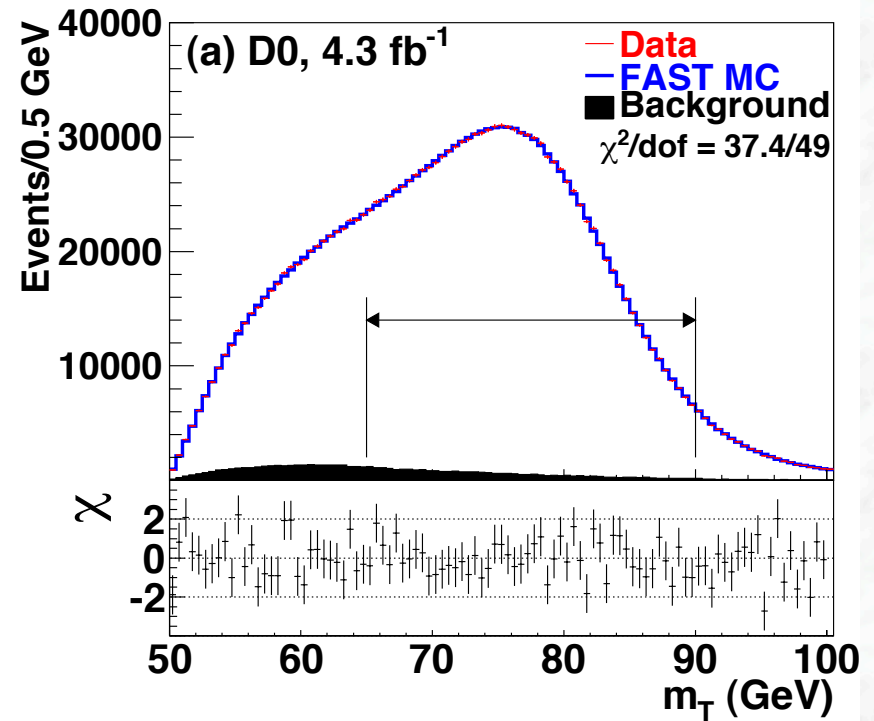


$$m_W = 80.35 \pm 0.33 \pm 0.17 \text{ GeV}$$

## State of the art, today



1.68 M events, electrons  $|\eta| < 1.05$



$$m_W = 80.371 \pm 0.013 \text{ (stat.) GeV}$$

## Systematic uncertainties:

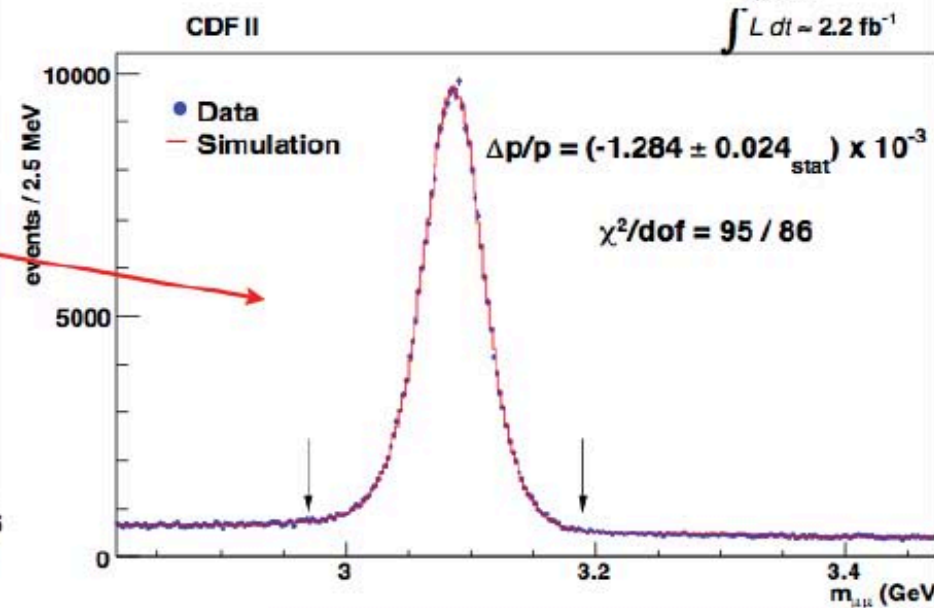
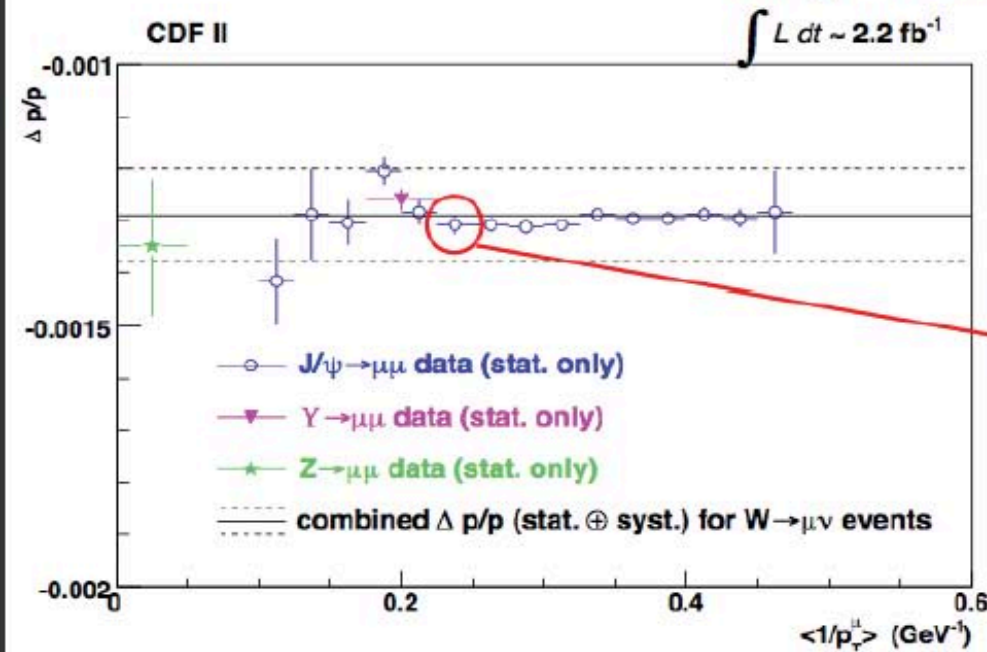
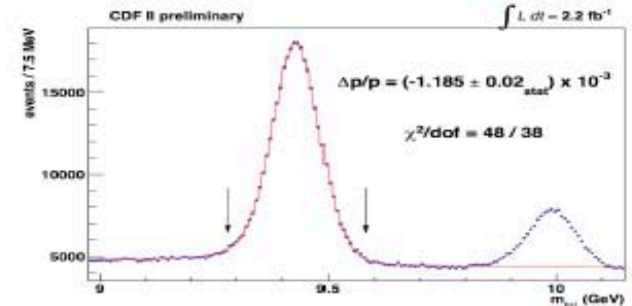
### New CDF Result ( $2.2 \text{ fb}^{-1}$ ) Transverse Mass Fit Uncertainties (MeV)

	<i>electrons</i>	<i>muons</i>	<i>common</i>
W statistics	19	16	0
Lepton energy scale	10	7	5
Lepton resolution	4	1	0
Recoil energy scale	5	5	5
Recoil energy resolution	7	7	7
Selection bias	0	0	0
Lepton removal	3	2	2
Backgrounds	4	3	0
pT(W) model	3	3	3
Parton dist. Functions	10	10	10
QED rad. Corrections	4	4	4
Total systematic	18	16	15
Total	26	23	

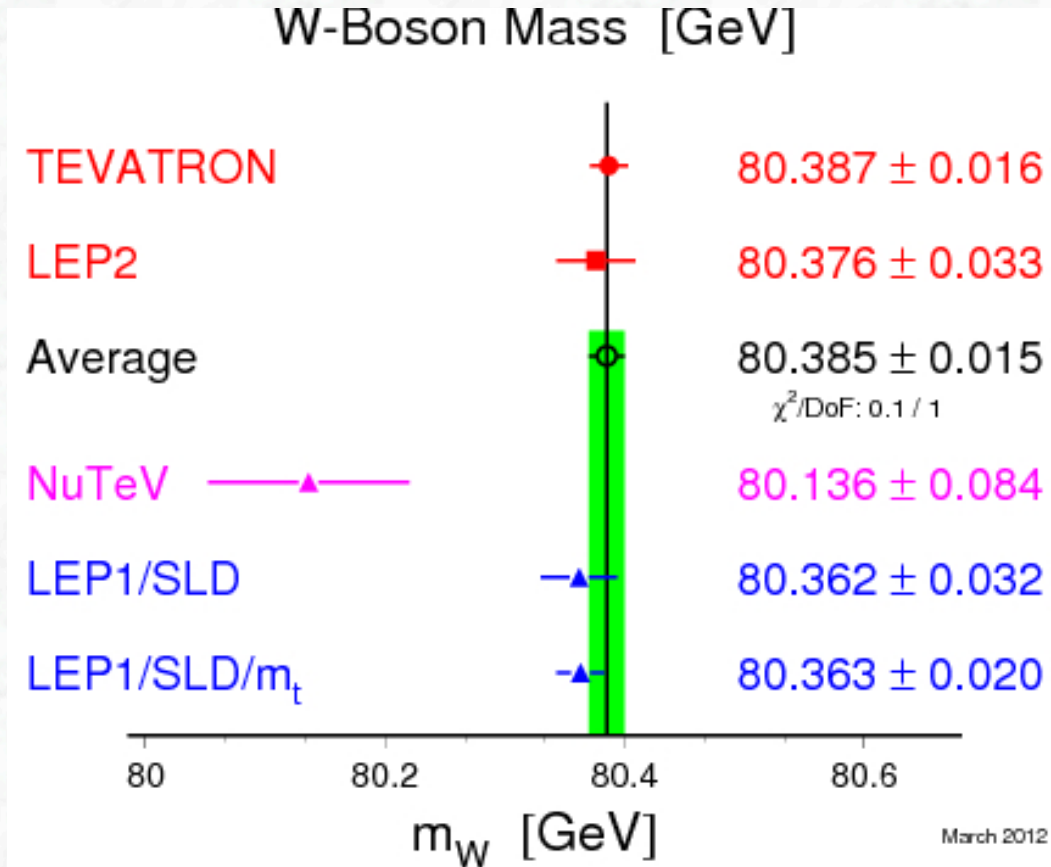
# Momentum Scale Calibration



- “Back bone” of CDF analysis is track  $p_T$  measurement in drift chamber (COT)
- Perform alignment using cosmic ray data:  $\sim 50\mu\text{m} \rightarrow \sim 5\mu\text{m}$  residual
- Calibrate momentum scale using samples of dimuon resonances ( $J/\psi$ ,  $\Upsilon$ ,  $Z$ )
  - Span a large range of  $p_T$
  - Flatness is a test of  $dE/dx$  modeling
- Final scale error of  $9 \times 10^{-5}$ :  $\Delta m_W = 7 \text{ MeV}$



# Summary of W-mass measurements

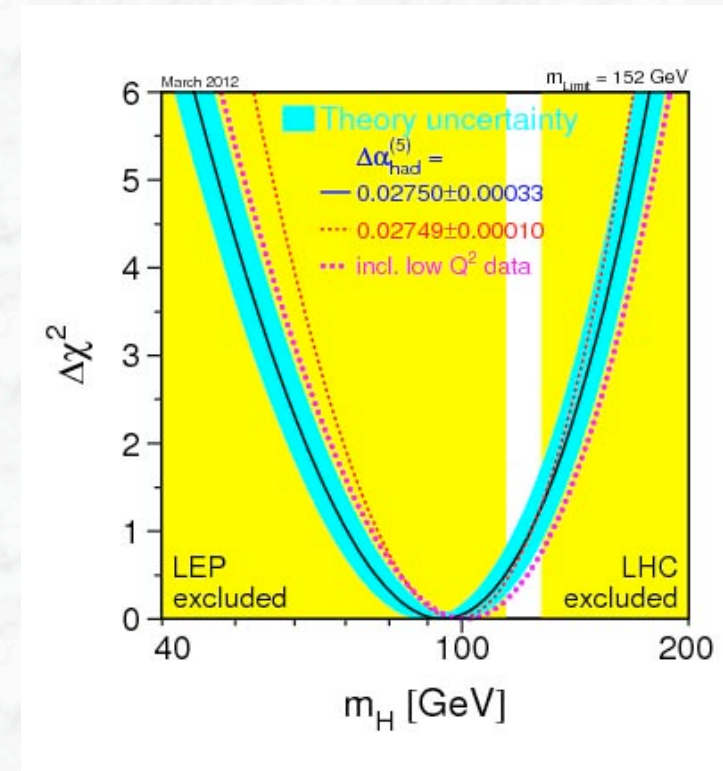
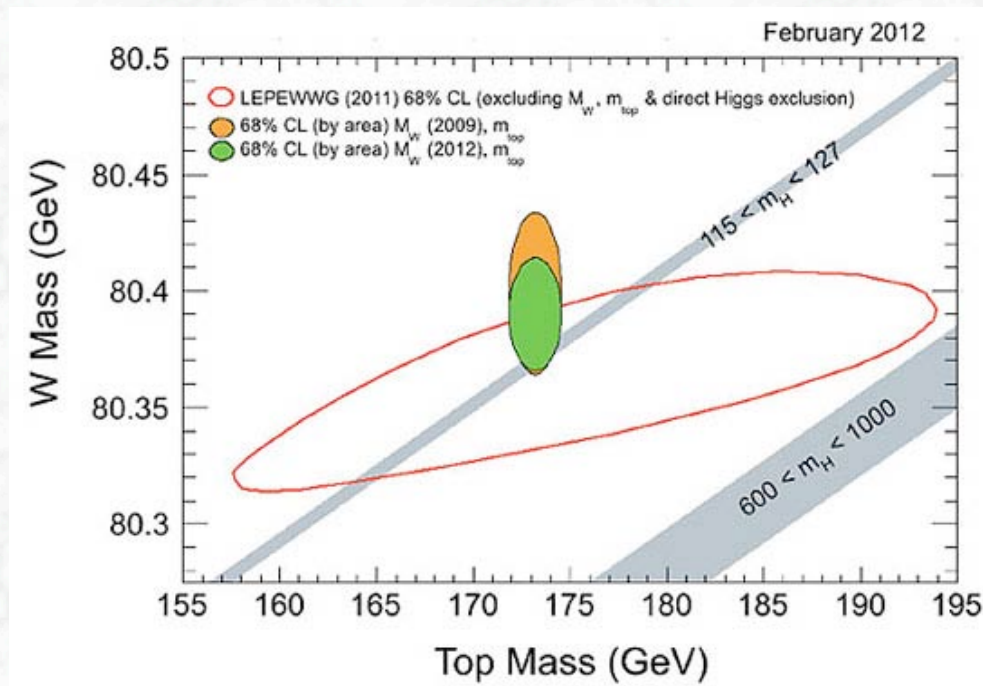


Precision obtained at the Tevatron is superior to the LEP-II precision

□  $m_W$  (from LEP2 + Tevatron) =  $80.385 \pm 0.015$  GeV

$2 \cdot 10^{-4}$

# Indirect limits from electroweak precision measurements

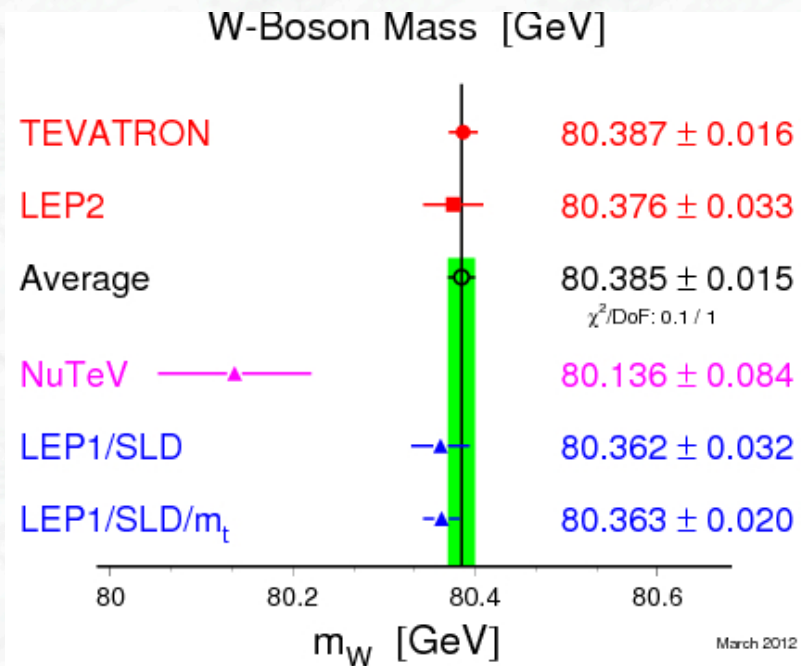


Impressive precision in  $W$  mass from the Tevatron (February 2012)

$$m_H = 94^{+29}_{-24} \text{ GeV}/c^2$$

$$m_H < 152 \text{ GeV}/c^2 \quad (95 \% \text{ C.L.})$$

The main story of 2011: eliminate 470 GeV of Higgs boson mass range



## Systematic uncertainties:

### New CDF Result ( $2.2 \text{ fb}^{-1}$ ) Transverse Mass Fit Uncertainties (MeV)

	<i>electrons</i>	<i>muons</i>	<i>common</i>
W statistics	19	16	0
Lepton energy scale	10	7	5
Lepton resolution	4	1	0
Recoil energy scale	5	5	5
Recoil energy resolution	7	7	7
Selection bias	0	0	0
Lepton removal	3	2	2
Backgrounds	4	3	0
<i>pT(W) model</i>	3	3	3
Parton dist. Functions	10	10	10
QED rad. Corrections	4	4	4
Total systematic	18	16	15
Total	26	23	

Can the LHC improve on this?

In principle yes, but probably not soon .and. not with 30 pileup events

- Very challenging (e-scale, hadronic recoil,  $p_T(W)$ ,... )
- However there is potential for reduction of uncertainties
  - statistics
  - statistically limited systematic uncertainties (marked in green above)
  - pdfs, energy scale, ....., recoil(?)



# What precision can be reached in Run II and at the LHC ?

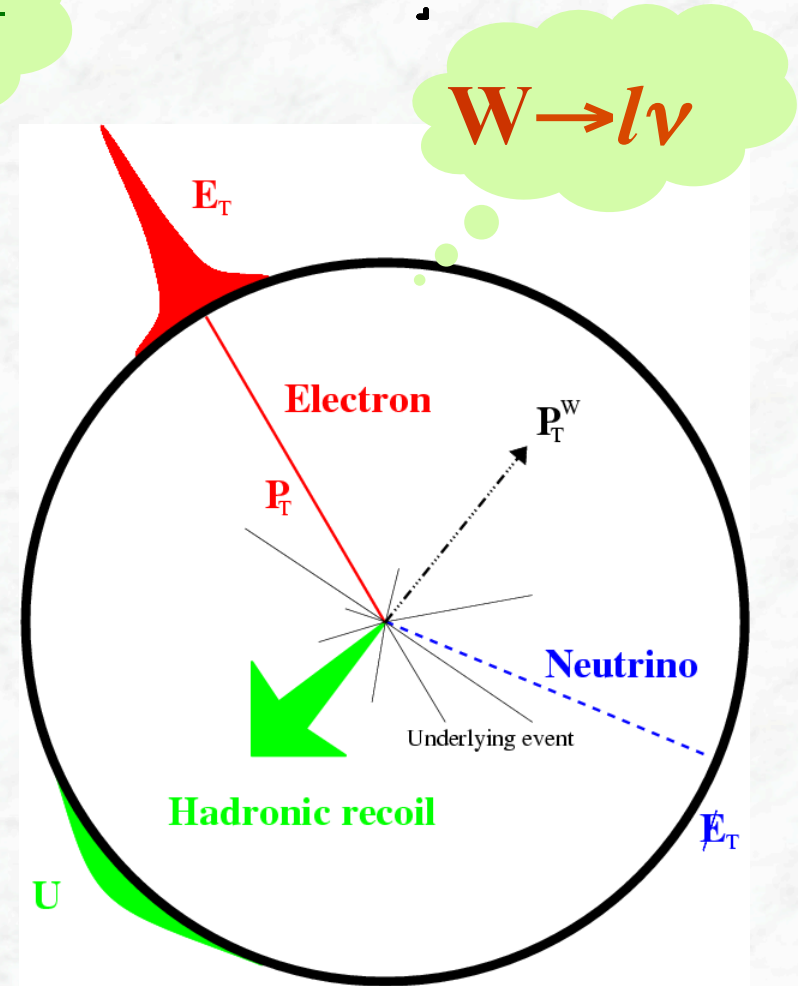
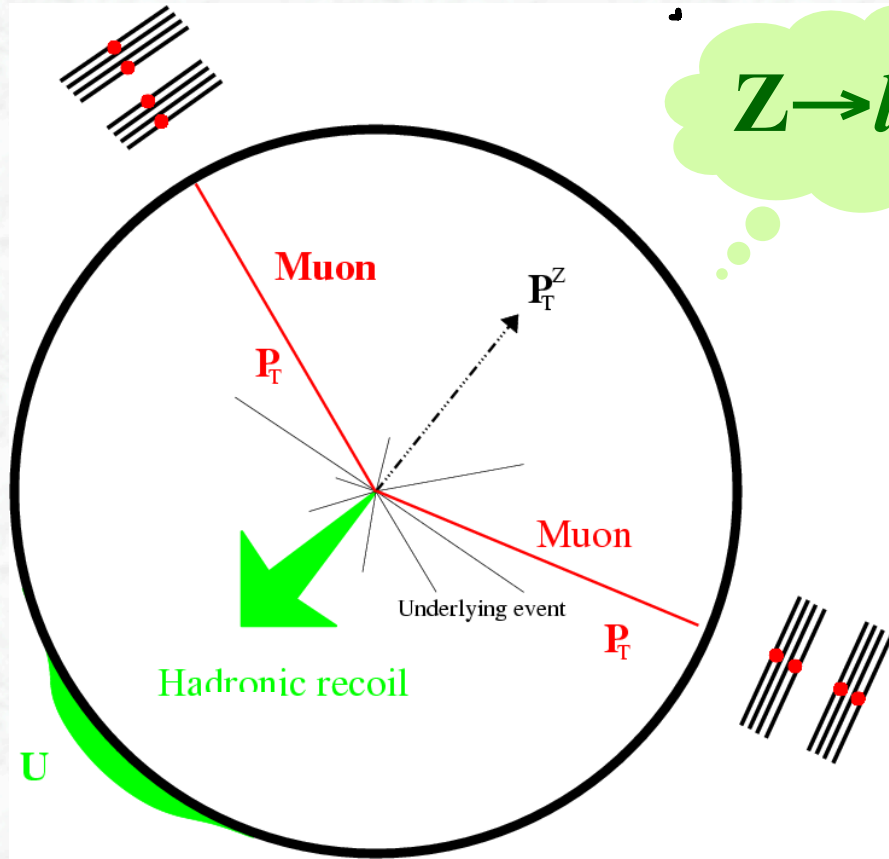
	Int. Luminosity	CDF	DØ	LHC
Numbers for a single decay channel		0.2 fb <sup>-1</sup>	1 fb <sup>-1</sup>	10 fb <sup>-1</sup>
	<b>Stat. error</b>	<b>48 MeV</b>	<b>23 MeV</b>	<b>2 MeV</b>
	Energy scale, lepton res.	30 MeV	34 MeV	4 MeV
W → eν	Monte Carlo model (P <sub>T</sub> <sup>W</sup> , structure functions, photon-radiation....)	16 MeV	12 MeV	7 MeV
	Background	8 MeV	2 MeV	2 MeV
	<b>Tot. Syst. error</b>	<b>39 MeV</b>	<b>37 MeV</b>	<b>8 MeV</b>
	<b>Total error</b>	<b>62 MeV</b>	<b>44 MeV</b>	<b>~10 MeV</b>

- Tevatron numbers are based on real data analyses
- LHC numbers should be considered as „ambitious goal“
  - Many systematic uncertainties can be controlled in situ, using the large Z → ℓℓ sample (p<sub>T</sub>(W), recoil model, resolution)
  - Lepton energy scale of ± 0.02% has to be achieved to reach the quoted numbers

Combining both experiments (ATLAS + CMS, 10 fb<sup>-1</sup>), both lepton species and assuming a scale uncertainty of ± 0.02% a total error in the order of

⇒ Δ m<sub>W</sub> ~ ± 10 MeV might be reached.

# Signature of Z and W decays



# What precision can be reached in Run II and at the LHC ?

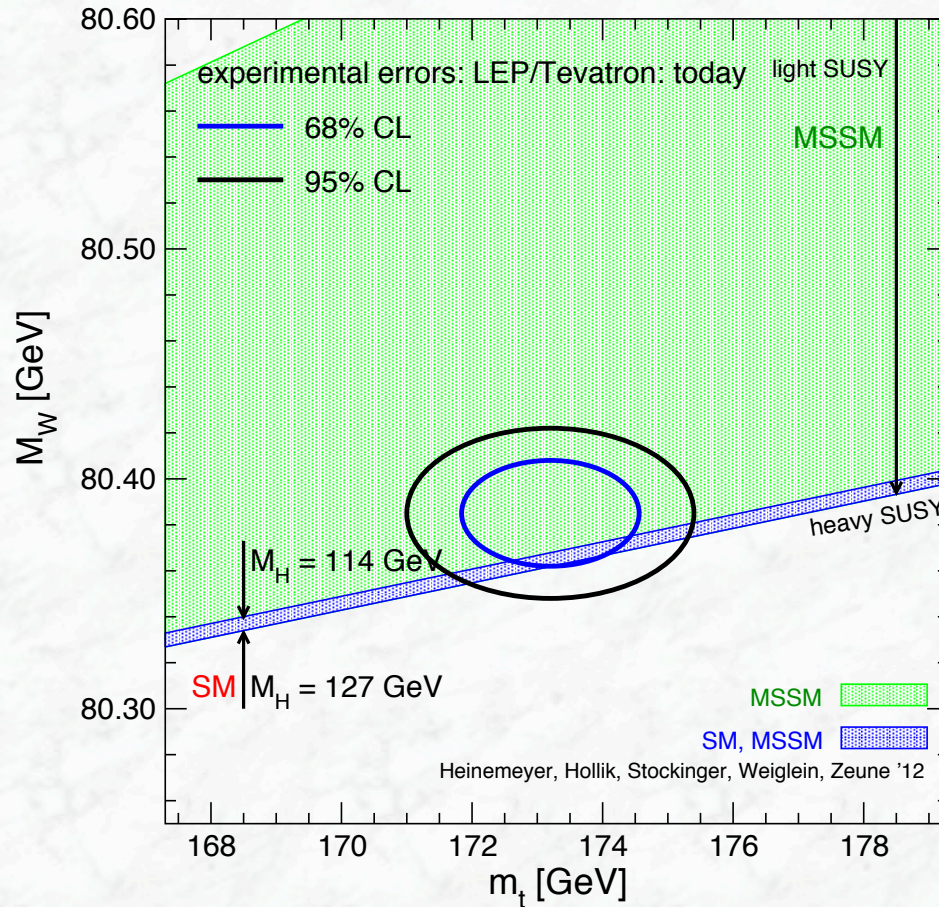
	Int. Luminosity	CDF	DØ	LHC
Numbers for a single decay channel		0.2 fb <sup>-1</sup>	1 fb <sup>-1</sup>	10 fb <sup>-1</sup>
	<b>Stat. error</b>	<b>48 MeV</b>	<b>23 MeV</b>	<b>2 MeV</b>
	Energy scale, lepton res.	30 MeV	34 MeV	4 MeV
W → eν	Monte Carlo model (P <sub>T</sub> <sup>W</sup> , structure functions, photon-radiation....)	16 MeV	12 MeV	7 MeV
	Background	8 MeV	2 MeV	2 MeV
	<b>Tot. Syst. error</b>	<b>39 MeV</b>	<b>37 MeV</b>	<b>8 MeV</b>
	<b>Total error</b>	<b>62 MeV</b>	<b>44 MeV</b>	<b>~10 MeV</b>

- Tevatron numbers are based on real data analyses
- LHC numbers should be considered as „ambitious goal“
  - Many systematic uncertainties can be controlled in situ, using the large Z → ℓℓ sample (p<sub>T</sub>(W), recoil model, resolution)
  - Lepton energy scale of ± 0.02% has to be achieved to reach the quoted numbers

Combining both experiments (ATLAS + CMS, 10 fb<sup>-1</sup>), both lepton species and assuming a scale uncertainty of ± 0.02% a total error in the order of

⇒ **Δ m<sub>W</sub> ~ ± 10 MeV** might be reached.

2012



Ultimate test of the Standard Model:

Compare direct prediction of the Higgs boson mass with direct observation