# Higgs analyses at the LHC

#### Part IIb: VH, H → bb Part III: Higgs boson parameters





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#### Search for VH production with $H \rightarrow bb$ decays



- Exploit three leptonic vector boson decay modes ٠  $\rightarrow$  split analysis in 0, 1, and 2-lepton categories
- Require 2 b-tagged jets (working point for 70% efficiency)
- Major background: W/Z bb, W+jets, tt •
- Signal-to-background ratio improves for "boosted Higgs boson", split analysis in bins of  $p_{T}(V)$

ATLAS: in total 15 categories  $(0,1,2 \text{ jets} \times p_T \text{ bins})$ 

CMS: multivariate analysis







#### Event selection for $H \rightarrow bb$ analyses

#### (i) Basic event selection for the three channels

Object	0-lepton	1-lepton	2-lepton			
Lentons	0 loose leptons	1 tight lepton	1 medium lepton			
Leptons		+ 0 loose leptons	+ 1 loose lepton			
		2 b-tags				
Tets	$p_{\rm T}^{\rm jet_1} > 45 { m ~GeV}$					
Jets		$p_{\rm T}^{\rm jet_2} > 20  { m GeV}$				
	$+ \leq 1$ extra jets					
Missing F_	$E_{\rm T}^{\rm miss} > 120  {\rm GeV}$	$E_{\rm T}^{\rm miss} > 25  {\rm Gev}$	$E_{\rm T}^{\rm miss} < 60 { m GeV}$			
wissing $L_T$	$p_{\rm T}^{\rm miss} > 30  { m GeV}$	-	-			
	$\Delta \phi(E_{\rm T}^{\rm miss}, p_{\rm T}^{\rm miss}) < \pi/2$					
	$\min[\Delta \phi(E_{T}^{\text{miss}}, \text{jet})] > 1.5$					
	$\Delta \phi(E_{\rm T}^{\rm miss}, b\bar{b}) > 2.8$					
Vector Boson	-	$m_{\rm T}^W < 120 { m GeV}$	$83 < m_{\ell\ell} < 99 \text{ GeV}$			

#### (ii) Further topological criteria in intervals of $p_T(V)$

	$p_{\rm T}^V$ [GeV]	0-90	90-120	120-160	160-200	>200
All Channels	$\Delta R(b, \bar{b})$	0.7-3.4	0.7-3.0	0.7-2.3	0.7-1.8	<1.4
1 lonton	$E_{\rm T}^{\rm miss}$ [GeV]	>25				>50
1-lepton	$m_{\rm T}^W$ [GeV]		0			

#### Definition of signal and control regions

. ,	2jets, Itag	3jets, I tag	2jets, 2tag	3jets, 2tag	top e-μ CR
3 p <sub>T</sub> <sup>V</sup> bins x 0-lepton	CR	CR	SR	SR	-
5 p <sub>T</sub> <sup>V</sup> bins x I-lepton	CR	CR	SR	SR	-
5 pT <sup>V</sup> bins x 2-lepton	CR	CR	SR	SR	CR
CR=Control Region (low SR=Signal Region	S/B)			l - m	electron+I muor <sub>«</sub> >40 GeV

- Common nuisance parameters across regions
- Systematic uncertainties on extrapolation between control and signal regions

#### Background normalization, interplay of regions



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#### Reconstructed mass distributions -full data set, 7 and 8 TeV (a selection, high $p_T$ bins)-





#### ATLAS-CONF-2013-079

1 lepton



#### Demonstration of di-boson production with $Z \rightarrow bb$ in ATLAS

combination (all bins, channels) data - background



#### weighted distribution, by S/B ratio



Di-boson signal established (important "calibration" signal; a Standard Model Higgs boson signal is included as background)

Significance  $5.1\sigma$ 

 $\mu_{WZ+WW}$  = 0.90 ± 0.20

ATLAS-CONF-2013-079



#### **Di-boson signal strength**



ATLAS-CONF-2013-079







Weighted (by S/B ratio) backgroundsubtracted m<sub>bb</sub> mass distribution



- Di-boson signal established
- Signal size consistent with expectations from Standard Model

#### CMS PAS HIG-13-012



#### ATLAS results on the search for VH, H $\rightarrow$ bb decays



#### ATLAS-CONF-2013-079



 $\mu_{\rm H}$  = 0.2 ± 0.5 (stat) ± 0.4 (syst)

Probability of obtaining a result more background-like than the observed in the presence of a SM signal ( $\mu$ =1) is 0.11



CMS PAS HIG-13-012





m<sub>H</sub> = 125 GeV:

Observed 95% CL:	1.89 $\sigma_{SM}$
Expected (no Higgs):	0.95 $\sigma_{\text{SM}}$

 $\mu_{\rm H}$  = 1.00 ± 0.49

# Is the new particle the Higgs Boson ?

• Production rates ?

#### Couplings to bosons and fermions





• Spin, J<sup>P</sup> quantum number



# Higgs boson parameters

- After the discovery of the new boson, the most important question is:
  - What are its properties ? (mass, spin, couplings, ...)
  - Is it the Higgs boson of the Standard Model? Or can we finds signs of Physics Beyond the Standard Model by studying its properties?

 Much attention of the LHC (and Tevatron) collaborations and from the theory community has been devoted to these questions during the past year.

# The Higgs boson mass

- The mass of the Higgs boson can be precisely determined by using the measurements in the high resolution H  $\rightarrow \gamma\gamma$  and H  $\rightarrow ZZ^* \rightarrow 4I$  channels
- Relevant here are the lepton/photon energy calibration, its stability and the good lepton / photon energy resolution



#### Higgs boson mass: results from CMS





 $m_{H} = 125.7 \pm 0.3$ (stat)  $\pm 0.3$ (syst) GeV

In the combination the relative signal strength for the two decay modes is constrained by the SM values

#### Higgs boson mass: results from ATLAS

ATLAS

#### AS arXiv:1307.1427



 $m_{H} = 125.5 \pm 0.2 (stat)_{-0.6}^{+0.5} (syst) \text{ GeV}$ 





Likelihood as a function of the mass difference,  $\Delta m_H$ , profiling over the common mass  $m_H$ . The signal strength parameters  $\mu_{\gamma\gamma}$  and  $\mu_{41}$ are allowed to vary independently. Consistency between the fitted mass values from likelihood value for  $\Delta m = 0$  w.r.t. best fit value for  $\Delta m$ .

 $\Delta m = 2.3^{+0.6}_{-0.7}$ (stat) ± 0.6 (syst) GeV

- Probability for disfavoring the  $\Delta m = 0$ hypothesis by more than observed is 1.5% (2.4 $\sigma$ )
- Increases to 8%, by fixing the three principle sources contributing to the  $e/\gamma$  energy scale uncertainty (material, pre-sampler energy scale, calibration procedure) to their  $\pm 1\sigma$  values

#### Consistency with the $H \rightarrow WW^*$ channel (mass vs. signal strength)









#### Signal strength in di-boson decay modes -including full data set-



• Data are consistent with the hypothesis of a Standard Model Higgs boson:

 $\mu = 1.33^{+0.21}_{-0.18}$ 

- Experimental uncertainties are still too large to get excited about "high" γγ signal strength
- Signal strengths in fermionic decay modes have large uncertainties, but are compatible with SM value of 1;

If preliminary  $H \rightarrow \tau\tau$  and  $H \rightarrow$  bb results are included:

 $\mu = 1.23 \pm 0.18$ 

Ratios of production cross sections for the various processes (ggF, VBF,..) fixed to SM values

# Signal strength in all decay and production modes<br/>-including full data set-Combined $\sqrt{s} = 7 \text{ TeV}, L \le 5.1 \text{ fb}^{-1} \sqrt{s} = 8 \text{ TeV}, L \le 19.6 \text{ fb}^{-1}$ Combined $\sqrt{s} = 7 \text{ TeV}, L \le 5.1 \text{ fb}^{-1} \sqrt{s} = 8 \text{ TeV}, L \le 19.6 \text{ fb}^{-1}$ Combined $\sqrt{s} = 7 \text{ TeV}, L \le 5.1 \text{ fb}^{-1} \sqrt{s} = 8 \text{ TeV}, L \le 19.6 \text{ fb}^{-1}$



• Data are consistent with the hypothesis of a Standard Model Higgs boson:

 $\mu = 0.80 \pm 0.14$ 

• Signal strengths in fermionic decay modes have large uncertainties, but are compatible with SM value of 1



- Sensitivity to (ggF + ttH) and (VBF+VH) production fractions, modulo branching ratio factors  $\rm B/B_{SM}$
- Good agreement with the Standard Model, within the large uncertainties
- A combination of the different decay modes is not performed, since it would require introducing hypotheses on the relative branching ratios;



#### Evidence for production via vector boson fusion



- Fit for the ratio of  $\mu_{VBF+VH}$  /  $\mu_{ggF+ttH}$  for the individual channels (model independent)
- Good agreement with SM expectation for individual channels and the combination)

#### Next step: combination of results



#### Evidence for production via vector boson fusion





 $\mu_{VBF} / \mu_{ggF+ttH} = 1.4^{+0.4}_{-0.3} (stat)^{+0.6}_{-0.4} (syst)$ 

 $3.3\sigma$  evidence for VBF production



 $3.2\sigma$  evidence for V-boson mediated production

### Higgs boson couplings

Production and decay involve several couplings



Decays: e.g.  $H \rightarrow \gamma\gamma$  (best example) (Decay widths depends on W and top coupling, destructive interference)





Standard Model couplings are tested by introducing coupling scale factors  $\kappa$ 

 $g_i = \kappa g_i^{SM}$ 

Standard Model tree level amplitudes:

$$\Gamma_{ff} \propto \left(\kappa_f \frac{m_f}{\nu}\right)^2 = \kappa_f^2 \cdot \Gamma_{ff}^{SM}$$

$$\Gamma_{VV} \propto \left(\kappa_f \frac{m_V^2}{v}\right)^2 = \kappa_V^2 \cdot \Gamma_{VV}^{SM}$$

## Higgs boson couplings

• Example:  $H \rightarrow \gamma \gamma$   $\Gamma_{\gamma\gamma} \propto \left| 1.28 \kappa_W - 0.28 \kappa_t \right|^2 \cdot \Gamma_{\gamma\gamma}^{SM}$   $-H - \psi_{w,t} \psi_{w,t} \psi_{w,t} \psi_{w,t}$ 



- Loop scaling factors can be expressed in terms of  $\kappa_{f}$  and  $~\kappa_{V}$
- The analysis is also done in terms of effective loop couplings  $\kappa_{q}$  and  $\kappa_{v}$

#### Higgs boson couplings

- Benchmarks defined by LHC cross section working group (leading-order tree-level framework):
  - Signals observed originate from a single resonance; (mass assumed here is 125.5 GeV)
  - Narrow width approximation:  $\rightarrow$  rates for given channels can be decomposed as:

$$\sigma \cdot B (i \to H \to f) = \frac{\sigma_i \cdot \Gamma_f}{\Gamma_H}$$
 i, f = initial, final state  
 $\Gamma_f, \Gamma_H = partial, total width$ 

- Modifications to coupling strength are considered (coupling scale factors  $\kappa$ ), tensor structure of Lagrangian assumed as in Standard Model

## Scaling of cross sections with $\kappa_{\text{F}} \, \text{and} \, \kappa_{\text{V}} \, \text{factors}$

$$\sigma \cdot \text{BR} (\text{gg} \to \text{H} \to \gamma\gamma) = \sigma_{\text{SM}}(\text{gg} \to \text{H}) \cdot \text{BR}_{\text{SM}}(\text{H} \to \gamma\gamma) \cdot \frac{\kappa_{\text{g}}^2 \cdot \kappa_{\gamma}^2}{\kappa_{\text{H}}^2}$$

$$\begin{split} \sigma(qg \to H) * \mathrm{BR}(H \to \gamma\gamma) &\sim \quad \frac{\kappa_F^2 \cdot \kappa_\gamma^2(\kappa_F, \kappa_V)}{0.75 \cdot \kappa_F^2 + 0.25 \cdot \kappa_V^2} \\ \sigma(qq' \to qq'H) * \mathrm{BR}(H \to \gamma\gamma) &\sim \quad \frac{\kappa_V^2 \cdot \kappa_\gamma^2(\kappa_F, \kappa_V)}{0.75 \cdot \kappa_F^2 + 0.25 \cdot \kappa_V^2} \\ \sigma(gg \to H) * \mathrm{BR}(H \to ZZ^{(*)}, H \to WW^{(*)}) &\sim \quad \frac{\kappa_F^2 \cdot \kappa_V^2}{0.75 \cdot \kappa_F^2 + 0.25 \cdot \kappa_V^2} \\ \sigma(qq' \to qq'H) * \mathrm{BR}(H \to ZZ^{(*)}, H \to WW^{(*)}) &\sim \quad \frac{\kappa_V^2 \cdot \kappa_V^2}{0.75 \cdot \kappa_F^2 + 0.25 \cdot \kappa_V^2} \\ \sigma(qq' \to qq'H, VH) * \mathrm{BR}(H \to \tau\tau, H \to b\bar{b}) &\sim \quad \frac{\kappa_V^2 \cdot \kappa_F^2}{0.75 \cdot \kappa_F^2 + 0.25 \cdot \kappa_V^2} \end{split}$$



## (i) Couplings to fermions and bosons

Assume only one scale factor for fermion and vector couplings:

 $\kappa_{V} = \kappa_{W} = \kappa_{Z}$  $\kappa_{F} = \kappa_{t} = \kappa_{b} = \kappa_{\tau}$ 

The size of the current data set is insufficient to quantify all parameters

- Assume that H → γγ and gg → H loops and the total Higgs boson width depend only on κ<sub>V</sub> and κ<sub>F</sub> (no contributions from physics beyond the Standard Model)
- Sensitivity to relative sign between κ<sub>F</sub> and κ<sub>V</sub> only from interference term in H → γγ decays (assume κ<sub>V</sub> > 0)

$$-\underbrace{H}_{w,t} \underbrace{w,t}_{w,t} \underbrace{w,t}_{\gamma}$$

#### (i) Couplings to fermions and bosons (cont).





Results: Data are consistent with the SM expectation;



## (ii) Ratio of couplings to the W and Z bosons



- Relation between  $m_W$  and  $m_Z$  in the Standard Model requires  $\lambda_{WZ} := \kappa_W / \kappa_Z = 1$ ( $\rho$  parameter required to be 1)
- Sensitivity via VBF and VH production and H  $\rightarrow$  WW and H  $\rightarrow$  ZZ rates



68% CL intervals:  $\lambda_{WZ} \in [0.61, 1.04]$ 



## (iii) Constraints on production and decay loops



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- Test on contributions from other particles contributing to loop-induced processes
- Assume nominal couplings for all SM particles  $\kappa_i = 1$ and that the new particles do not contribute to the Higgs boson width
- Fit for effective scale factors  $\kappa_g$  and  $\kappa_\gamma$





## (iv) Constraints on invisible decays (BR<sub>BSM</sub>)



- There might be invisible decays that would increase the total decay width:
- $\Gamma_{\rm H} = \Gamma_{\rm SM} + \Gamma_{\rm BSM}$  (BR<sub>BSM</sub> =  $\Gamma_{\rm BSM} / \Gamma_{\rm H}$ )



Assume  $k_v \le 1$  (motivated by EWSB) Fit for seven parameters



BR<sub>BSM</sub> < 0.64 (95% CL) 0.66 (expected)

BR<sub>BSM</sub> < 0.60 (95% CL) 0.67 (expected)

#### Summary of coupling scale factor measurements



$$\lambda_{FV} = \kappa_F / \kappa_V$$
$$\kappa_{VV} = \kappa_V \kappa_V / \kappa_H$$

If assumption of no contributions from new particles to the Higgs boson width is relaxed, only the ratio of  $k_F/k_V$  can be measured

Extended fit, decouple H  $\rightarrow \gamma \gamma$ event rate from the measurement of  $\lambda_{WZ}$ 

- $\kappa_V$  constrained at ±10% level
- Couplings to fermions indirectly observed (5σ)
- $\kappa_W/\kappa_Z$  found to be consistent with one
- No evidence for significant anomalous contributions to the gg → H and H → γγ loops

(for fixed nominal couplings of SM particles and no BSM contributions to Higgs width)

# **Spin and Parity**



Wolfgang Pauli and Niels Bohr studying the motion of a gyro (1952, at the opening of the institute for theoretical physics in Lund /Sweden) Standard Model Higgs boson:  $J^P = 0^+$ 

→ strategy is to falsify other hypotheses (0<sup>-</sup>, 1<sup>-</sup>, 1<sup>+</sup>, 2<sup>-</sup>, 2<sup>+</sup>)

and demonstrate consistency with the 0<sup>+</sup> hypothesis

Spin 1: strongly dis-favoured by observed  $H \rightarrow \gamma\gamma$  decays, Landau-Yang theorem



#### J<sup>P</sup> = 0<sup>-</sup> versus J<sup>P</sup>=0<sup>+</sup>

#### $(H \rightarrow ZZ^{(*)} \rightarrow 4\ell \text{ events})$

- Sensitive variables:
  - Masses of the two Z bosons
  - Production angle  $\,\theta^*$
  - Four decay angles  $\Phi_1$ ,  $\Phi$ ,  $\theta_1$  and  $\theta_2$
- Perform multivariate analysis (Boosted decision tree, similar sensitivity using matrix-element method)







#### $J^{P} = 0^{-}$ versus $J^{P} = 0^{+}$

#### $(H \rightarrow ZZ^{(*)} \rightarrow 4\ell \text{ events})$

- Sensitive variables:
  - Masses of the two Z bosons
  - Production angle  $\theta^*$
  - Four decay angles  $\Phi_1, \Phi, \theta_1$  and  $\theta_2$
- Matrix-Element based discriminant D<sub>JP</sub>









# Further CMS results based on $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ decays

#### CMS PAS HIG-13-005

$J^p$	production	comment	expect (µ=1)	obs. 0 <sup>+</sup>	obs. $J^p$	CLs
0-	$gg \rightarrow X$	pseudoscalar	<b>2.6</b> σ (2.8σ)	$0.5\sigma$	$3.3\sigma$	0.16%
$0_h^+$	$gg \rightarrow X$	higher dim operators	$1.7\sigma (1.8\sigma)$	$0.0\sigma$	$1.7\sigma$	8.1%
$2^+_{mgg}$	$gg \rightarrow X$	minimal couplings	$1.8\sigma$ ( $1.9\sigma$ )	$0.8\sigma$	$2.7\sigma$	1.5%
$2^+_{mq\bar{q}}$	$q\bar{q} \rightarrow X$	minimal couplings	$1.7\sigma (1.9\sigma)$	$1.8\sigma$	$4.0\sigma$	<0.1%
1- "	$q\bar{q} \rightarrow X$	exotic vector	$2.8\sigma$ (3.1 $\sigma$ )	$1.4\sigma$	$>4.0\sigma$	<0.1%
1+	$q\bar{q} \rightarrow X$	exotic pseudovector	<b>2.3</b> σ <b>(2.6</b> σ)	$1.7\sigma$	$>4.0\sigma$	<0.1%





#### $J^{P} = 1^{+/-} \text{ versus } J^{P} = 0^{+}$ (H $\rightarrow$ ZZ<sup>\*</sup> and H $\rightarrow$ WW<sup>\*</sup> events)

- $H \rightarrow ZZ^*$ , as before: BDT separation based on masses and angles
- $H \rightarrow WW^*$ :  $m_{\ell\ell}, \Delta \phi_{\ell\ell}, p_T(\ell\ell), m_T$  carry information on spin, combine these variables using a BDT analysis



	p <sub>0</sub> (0+)	CL (1 <sup>+</sup> ) Exclusion	p <sub>0</sub> (0+)	CL (1 <sup>-</sup> ) Exclusion
H → ZZ*	0.55	99.8%	0.1	94%
$H \rightarrow WW^*$	0.70	92%	0.66	98%
Combination	0.62	99.97%	0.33	99.7%

$$q = \log \frac{\mathcal{L}(J^P = 0^+, \hat{\hat{\mu}}_{0^+}, \hat{\hat{\theta}}_{0^+})}{\mathcal{L}(J^P_{\text{alt}}, \hat{\hat{\mu}}_{J^P_{\text{alt}}}, \hat{\hat{\theta}}_{J^P_{\text{alt}}})}$$

q = test statistics to discriminate between two spin hypotheses



## $J^{P} = 2^{+} \text{ versus } J^{P} = 0^{+}$ (H $\rightarrow \gamma\gamma$ , H $\rightarrow ZZ^{*}$ , and H $\rightarrow WW^{*}$ events)

- Spin 2: consider graviton-like tensor, equivalent to a Kaluza-Klein graviton (Y. Gao et al, Phys. Rev. D81 (2010) 075022)
- Production via gluon fusion and qq annihilation possible; Studies are performed as a function of the qq annihilation fraction (f<sub>qq</sub>)
- Specific model 2<sup>+</sup><sub>m</sub>:
   minimal couplings to SM particles
   (f<sub>aa</sub> = 4% at LO, however, large uncertainties)



 Observed exclusion (combination of γγ, ZZ\* and WW\*) of J<sup>P</sup> = 2<sup>+</sup> (versus the SM J<sup>P</sup> =0<sup>+</sup>) exceeds 99.9%, independent of f<sub>qq</sub>; Complementary behaviour of the different channels



#### Example: $H \rightarrow \gamma\gamma$ contribution

Use decay angle w.r.t. collision axis





Exclude  $J^P=2^+$  (produced via gluon fusion,  $f_{qq}=0$ ) (vs. 0<sup>+</sup>) via  $H \rightarrow \gamma\gamma$  decays with 99.3% CL

 $\cos \theta^*$  distribution in signal region, after background subtraction

#### CMS summary on coupling results





CMS PAS HIG-13-005

"The consistency of the couplings of the observed boson with those predicted for the Standard Model Higgs boson is tested in various ways, and no significant deviations are found."

#### ATLAS summary on spin results

ATLAS arXiv:1307:1432





"These studies provide evidence for the spin-0 nature of the Higgs boson, with positive parity being strongly preferred."

# Summary

- A milestone discovery announced in July 2012
- Signals have been impressively confirmed with additional data; The discovery phase has turned into the measurement phase
- ATLAS and CMS data are consistent with the expectations for the Standard Model Higgs boson (within present uncertainties)
  - Production rates and coupling strengths
  - Evidence for VBF production
  - Evidence for spin-0 (0<sup>-</sup> disfavoured)
- Exciting times ahead of us to study the Higgs boson with higher precision (> 2015) and to look for surprises (deviations? more Higgs bosons? ...)



# End of lectures



# Backup Slides





# Scaling of cross sections with $\kappa_{\text{F}}$ and $\kappa_{\text{V}}$ factors with no assumptions on the total width

In this case only ratios of coupling scale factors can be measured. Hence there are the following free parameters:

$$\lambda_{FV} = \kappa_F / \kappa_V \tag{16}$$

$$\kappa_{VV} = \kappa_V \cdot \kappa_V / \kappa_{\rm H} \quad . \tag{17}$$

 $\lambda_{FV}$  is the ratio of the fermion and vector coupling scale factors, and  $\kappa_{VV}$  an overall scale that includes the total width and applies to all rates. For the most relevant Higgs boson production and decay modes the following proportionality is found:

$$\sigma(qg \to H) * \mathrm{BR}(H \to \gamma\gamma) \sim \lambda_{FV}^2 \cdot \kappa_{VV}^2 \cdot \kappa_{\gamma}^2(\lambda_{FV}, 1)$$

$$\sigma(qq' \to qq'H) * \mathrm{BR}(H \to \gamma\gamma) \sim \kappa_{VV}^2 \cdot \kappa_{\gamma}^2(\lambda_{FV}, 1)$$

$$\sigma(qg \to H) * \mathrm{BR}(H \to ZZ^{(*)}, H \to WW^{(*)}) \sim \lambda_{FV}^2 \cdot \kappa_{VV}^2 \qquad (18)$$

$$\sigma(qq' \to qq'H) * \mathrm{BR}(H \to ZZ^{(*)}, H \to WW^{(*)}) \sim \kappa_{VV}^2$$

$$\sigma(qq' \to qq'H, VH) * \mathrm{BR}(H \to \tau\tau, H \to b\bar{b}) \sim \kappa_{VV}^2 \cdot \lambda_{FV}^2 \quad ,$$

#### General amplitude for a decay of a spin-2 resonance

#### (Y. Gao et al, Phys. Rev. D81 (2010) 075022)

For a decay of a spin-two resonance to two vector bosons, including ZZ, WW, and  $\gamma\gamma$ , the scattering amplitude has the following general form

$$\begin{aligned} A(X \to V_{1}V_{2}) &= \Lambda^{-1} \left[ 2g_{1}^{(2)}t_{\mu\nu}f^{*(1)\mu\alpha}f^{*(2)\nu\alpha} + 2g_{2}^{(2)}t_{\mu\nu}\frac{q_{\alpha}q_{\beta}}{\Lambda^{2}}f^{*(1)\mu\alpha}f^{*(2)\nu\beta} + g_{3}^{(2)}\frac{\tilde{q}^{\beta}\tilde{q}^{\alpha}}{\Lambda^{2}}t_{\beta\nu} \left( f^{*(1)\mu\nu}f^{*(2)}_{\mu\alpha} + f^{*(2)\mu\nu}f^{*(1)}_{\mu\alpha} \right) \\ &+ g_{4}^{(2)}\frac{\tilde{q}^{\nu}\tilde{q}^{\mu}}{\Lambda^{2}}t_{\mu\nu}f^{*(1)\alpha\beta}f^{*(2)}_{\alpha\beta} + m_{V}^{2} \left( 2g_{5}^{(2)}t_{\mu\nu}\epsilon^{*\mu}_{1}\epsilon^{*\nu}_{2} + 2g_{6}^{(2)}\frac{\tilde{q}^{\mu}q_{\alpha}}{\Lambda^{2}}t_{\mu\nu}\left(\epsilon^{*\nu}_{1}\epsilon^{*\alpha}_{2} - \epsilon^{*\alpha}_{1}\epsilon^{*\nu}_{2}\right) + g_{7}^{(2)}\frac{\tilde{q}^{\mu}\tilde{q}^{\nu}}{\Lambda^{2}}t_{\mu\nu}\epsilon^{*}_{1}\epsilon^{*}_{2} \right) \\ &+ g_{8}^{(2)}\frac{\tilde{q}_{\mu}\tilde{q}_{\nu}}{\Lambda^{2}}t_{\mu\nu}f^{*(1)\alpha\beta}\tilde{f}^{*(2)}_{\alpha\beta} + m_{V}^{2} \left( g_{9}^{(2)}\frac{t_{\mu\alpha}\tilde{q}^{\alpha}}{\Lambda^{2}}\epsilon_{\mu\nu\rho\sigma}\epsilon^{*\nu}_{1}\epsilon^{*\rho}_{2}q^{\sigma} + \frac{g_{10}^{(2)}t_{\mu\alpha}\tilde{q}^{\alpha}}{\Lambda^{4}}\epsilon_{\mu\nu\rho\sigma}q^{\rho}\tilde{q}^{\sigma}\left(\epsilon^{*\nu}_{1}\left(q\epsilon^{*}_{2}\right) + \epsilon^{*\nu}_{2}\left(q\epsilon^{*}_{1}\right)\right) \right) \right], \end{aligned}$$

where  $t_{\mu\nu}$  is the X wave function given by a symmetric traceless tensor [20]. This amplitude can be re-written as

TABLE I: List of scenarios chosen for the analysis of the production and decay of an exotic X particle with quantum numbers  $J^P$ . The subscripts m (minimal couplings) and h (couplings with higher-dimension operators) distinguish different scenarios, as discussed in the last column. The spin-zero and spin-one X production parameters do not affect the angular and mass distributions, and therefore are not specified.

scenario	X production	$X \to VV$ decay	comments
$0_m^+$	$gg \to X$	$g_1^{(0)} \neq 0$ in Eq. (9)	SM Higgs boson scalar
$0_h^+$	$gg \to X$	$g_2^{(0)} \neq 0$ in Eq. (9)	scalar with higher-dimension operators
0-	$gg \to X$	$g_4^{(0)}  e 0$ in Eq. (9)	pseudo-scalar
1+	$q \bar{q}  o X$	$b_2 \neq 0$ in Eq. (16)	exotic pseudo-vector
$1^{-}$	$q\bar{q}  o X$	$b_1 \neq 0$ in Eq. (16)	exotic vector
$2_m^+$	$g_{_1}^{(2)}  e 0$ in Eq. (18)	$g_1^{(2)} = g_5^{(2)} \neq 0$ in Eq. (18)	graviton-like tensor with minimal couplings
$2_h^+$	$g_4^{(2)}  e 0$ in Eq. (18)	$g_4^{(2)}  e 0$ in Eq. (18)	tensor with higher-dimension operators
$2_h^-$	$g_8^{(2)}  e 0$ in Eq. (18)	$g_8^{(2)}  e 0$ in Eq. (18)	"pseudo-tensor"