

Physics at the LHC



- Introduction, Detector Aspects
- Search for the Higgs Boson
 - Vector boson fusion mode
 - Measurement of Higgs boson parameters
- Standard Model Physics
 - W-mass measurement
 - Top Quark Physics
- Physics Beyond the Standard Model
 - SUSY Signatures
 - Search for Signals from Extra Dimensions

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Revised LHC Schedule



Dec. 2006	Ring closed and cold
Jan. - Mar. 2007	Machine commissioning
Spring 2007	First collisions , pilot run $L=5 \times 10^{32}$ to 2×10^{33} , $\leq 1 \text{ fb}^{-1}$ Start detector commissioning $\sim 10^5 Z \rightarrow ll, W \rightarrow lv, tt$ events
June - Dec. 2007	Complete detector commissioning, Physics run
→ 2009	$L=1-2 \times 10^{34}$, 100 fb^{-1} per year (high luminosity LHC)

low luminosity:	$L = 1 \times 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$	10 fb^{-1} / year
high luminosity:	$L = 1 \times 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$	100 fb^{-1} / year

Cross sections and production rates

$$\mathcal{L} = 1.0 \cdot 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$$

Process	σ	Events/s	Events/year
$W \rightarrow e\nu$	15 nb	15	10^8
$Z \rightarrow ee$	1.5 nb	1.5	10^7
$t\bar{t}$	800 pb	0.8	10^7
$b\bar{b}$	500 μb	10^5	10^{12}
QCD jets ($P_T > 200 \text{ GeV}$)	100 nb	10^2	10^9
$g\bar{g}$ ($m_{\tilde{g}} = 1 \text{ TeV}$)	1 pb	0.001	10^4
Higgs ($m_H = 0.2 \text{ TeV}$)	10 pb	0.01	10^5
($m_H = 0.8 \text{ TeV}$)	1 pb	0.001	10^4

Large production rates



- **Precision measurements**
at initial low luminosity
(W physics, top physics)
precision will be limited by systematic uncertainties.
- **Discoveries** (at low and high luminosity)
Mass reach for new particles up to $\sim 2 \text{ TeV}$
- **Disadvantages:**
 $\sigma_{\text{inelastic}} \sim 70 \text{ mb} \Rightarrow 700 \text{ Mio events / sec}$ at high L
Pile-up: 23 minimum bias events/bunch crossing at high L
2.3 minimum bias events/bunch crossing at low L

Detector Requirements

- Good measurement of **leptons** and **photons**

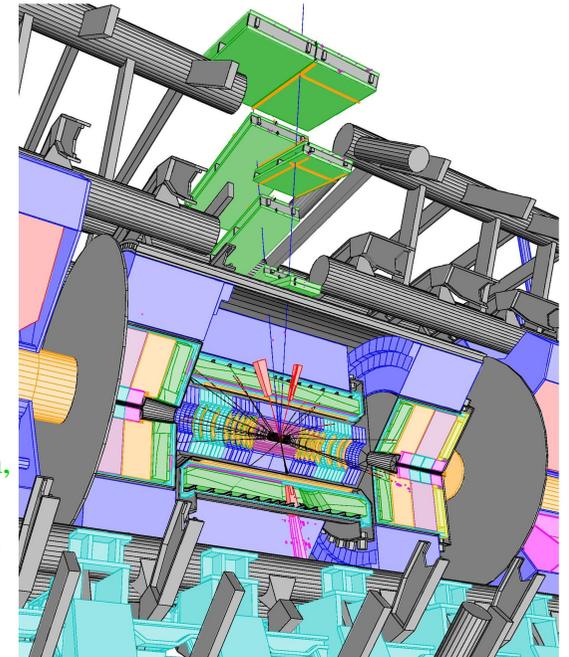
momentum range:

$$\sim \text{GeV} \quad (b \rightarrow l \nu c)$$

$$\sim \text{TeV} \quad (W \rightarrow l \nu)$$

lepton energy / momentum
scale: 0.1 % \rightarrow 0.02%

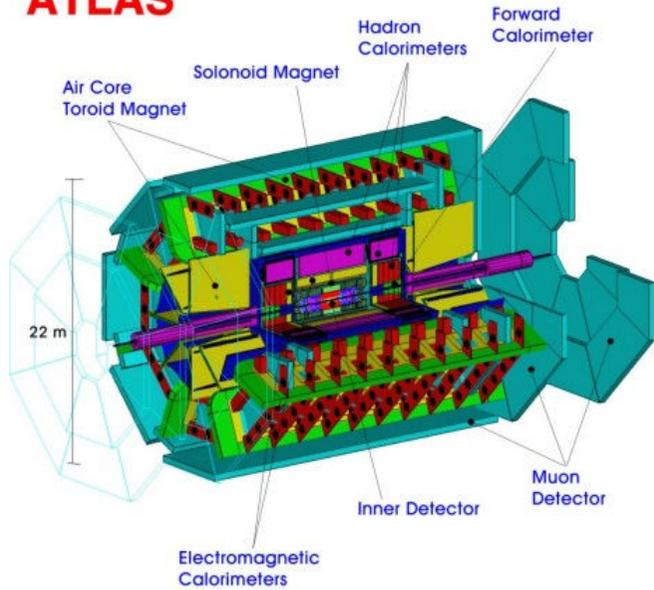
(large statistics for calibration,
 $Z \rightarrow ll$,
 m_Z is close to m_W and m_H (?)



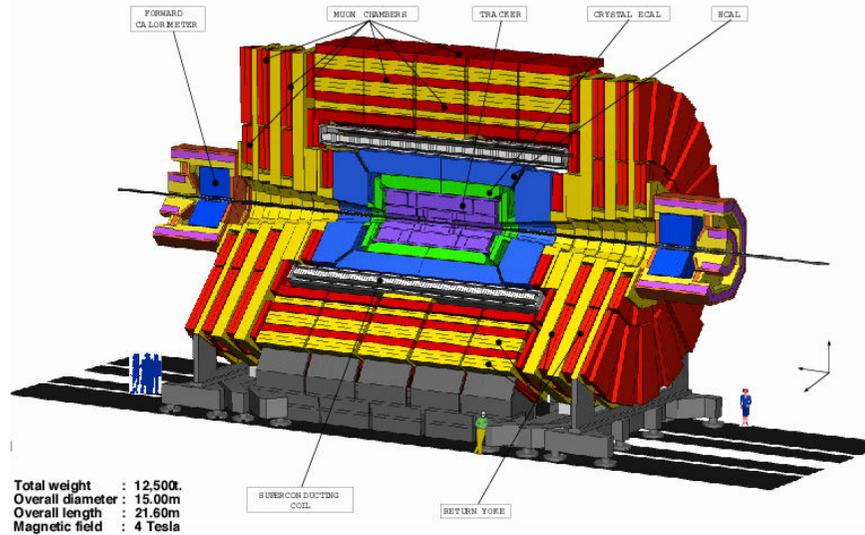
- Good measurement of **missing transverse energy** (E_T^{miss})
and
Jet energy measurements and jet-tagging in forward region
 \Rightarrow calorimeter coverage down to $\eta \sim 5$
Jet energy scale: 1% (relevant for m_{top} , SUSY)
- Efficient **b-tagging** and **τ identification**
(silicon strip and pixel detectors)
- **Fast (25 ns bunch crossing) and rad. hard detectors and electronics**

The pp experiments

ATLAS



CMS A Compact Solenoidal Detector for LHC



CMS detector construction

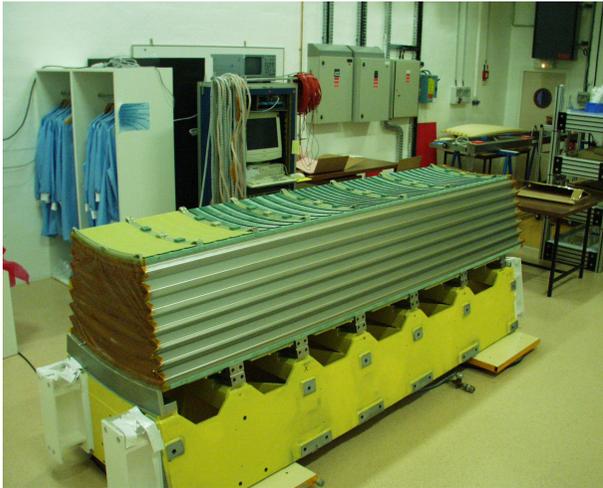


HCAL assembly

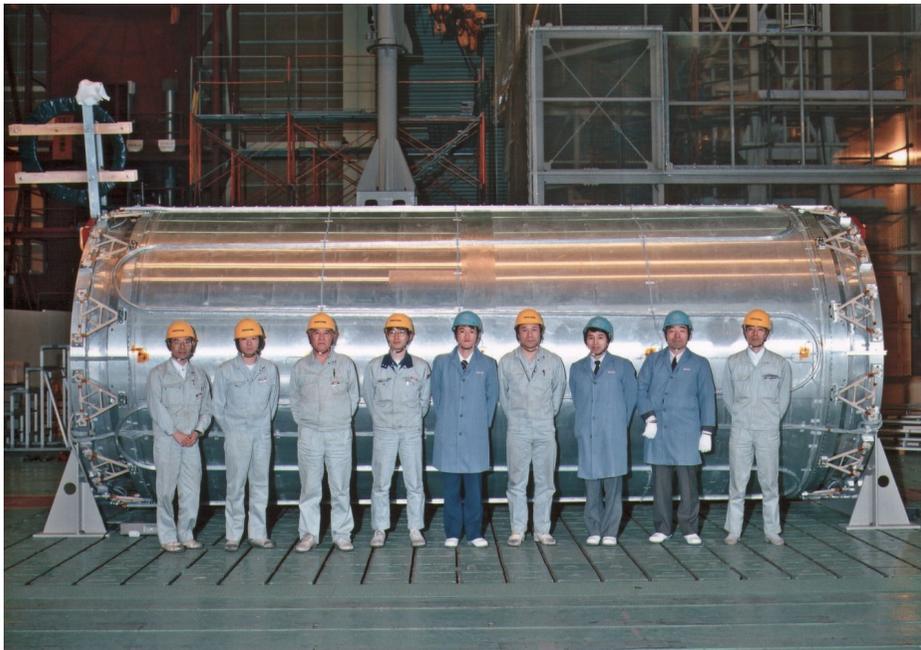


ATLAS detector construction

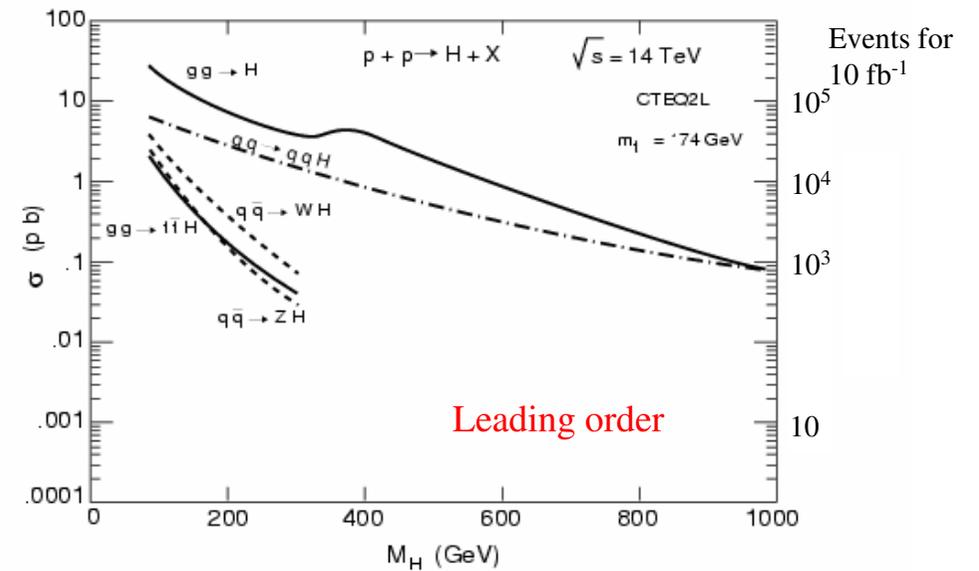
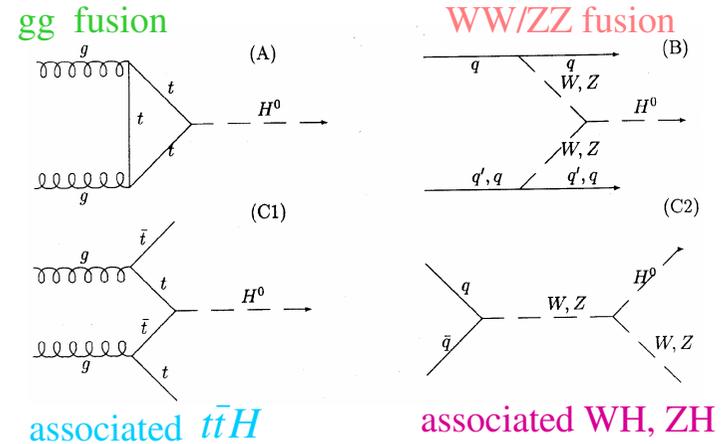
Liquid Argon Calorimeter



Superconducting solenoid ready

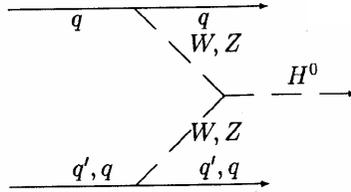


Search for the Higgs boson



- **K-factors** (\equiv higher-order corrections) = 1.6 – 1.9 $gg \rightarrow H$
- **Residual uncertainties** on NLO cross-sections (PDF, NNLO, etc.) $\leq 20\%$

Higgs production via Vector Boson Fusion



Motivation:

- Additional potential for Higgs boson discovery at low mass
- Important for the measurement of Higgs boson parameters (couplings to bosons, fermions (taus), total width)

proposed by D.Rainwater and D.Zeppenfeld et al.:
(hep-ph/9712271, hep-ph/9808468 and hep-ph/9906218)

Distinctive Signature of:

- two high P_T forward jets
- little jet activity in the central region
- ⇒ Jet Veto

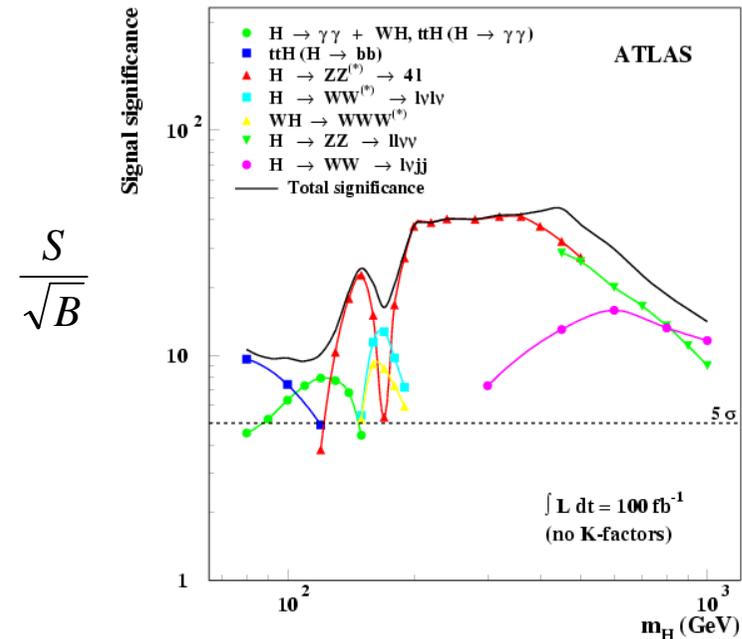
⇒ Experimental Issues:

- Forward jet reconstruction
- Jets from pile-up in the central/forward region

Channels studied:

- qqH → WW* → lνlν
- qqH → ττ → lννlνν
- lνν hadν

Main search channels at the LHC



$\underline{m_H} < 2 m_Z$: $t\bar{t}H \rightarrow \ell b\bar{b} + X, H \rightarrow \gamma\gamma,$
 $H \rightarrow ZZ^* \rightarrow 4\ell, H \rightarrow WW^{(*)} \rightarrow \ell\nu\ell\nu$

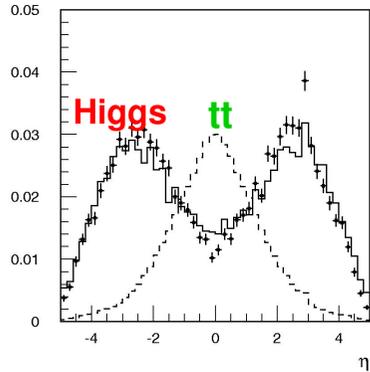
$\underline{m_H} > 2 m_Z$: $H \rightarrow ZZ \rightarrow 4\ell$
 $qqH \rightarrow ZZ \rightarrow \ell\ell\nu\nu$
 $qqH \rightarrow ZZ \rightarrow \ell\ell jj$
 $qqH \rightarrow WW \rightarrow \ell\nu jj$ } $m_H > 300 \text{ GeV}$
forward jet tag

10 fb⁻¹: Discovery possible over the full mass range, however, needs combination of ATLAS + CMS

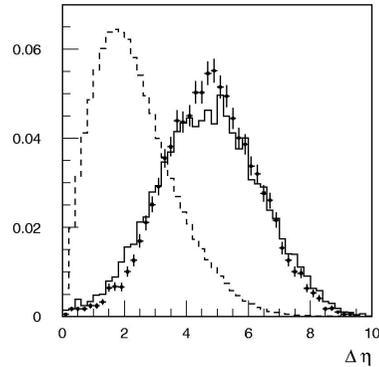
$M_H = 115 \text{ GeV: } S/\sqrt{B} = 4.7$

Forward tag Jets

Rapidity distribution of tag jets
VBF Higgs events vs. tt-background



Rapidity separation

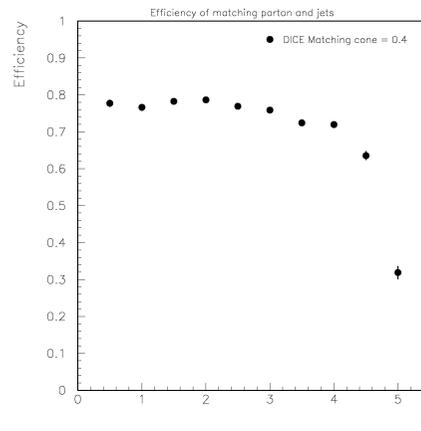


Forward tag jet reconstruction has been studied in full simulation in ATLAS

Results are consistent with TDR-results

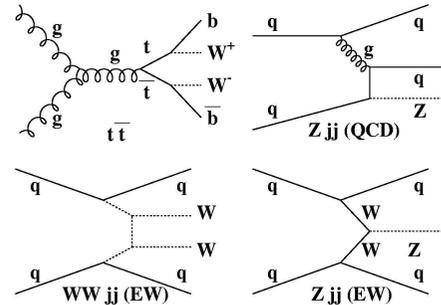
kin. eff. for tag jets = 51.9%
($P_T > 40/20$ GeV, $\Delta\eta > 3.6$)

tag eff. per jet: around 75%



Physics studies based on a fast simulation have been corrected for efficiency losses

Background:



QCD backgrounds:

tt production
Z + 2 jets
(PYTHIA MC)

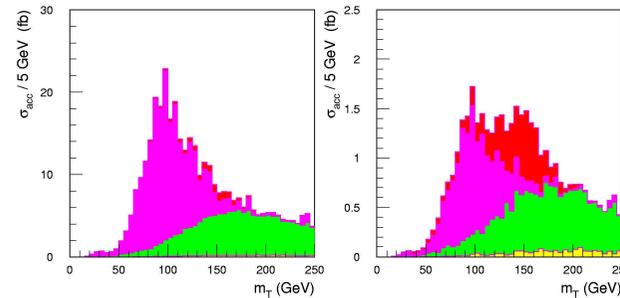
el.weak background:

WW jj production
Z + 2 jets
(matrix elements interfaced to PYTHIA)

Background rejection: $qqH \rightarrow WW^* \rightarrow l\nu l\nu$

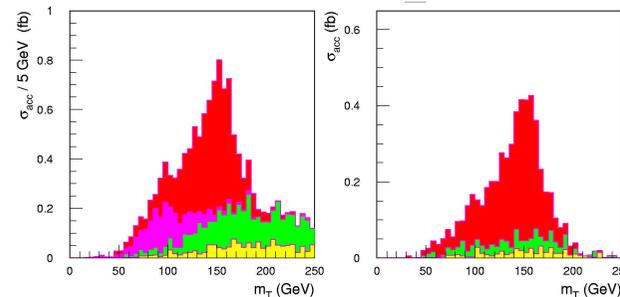
- Lepton P_T cuts and tag jet requirements ($\Delta\eta, P_T$)
- Require large mass of tag jet system, tau rejection
- Jet veto
- Lepton angular and mass cuts

$$M_T = \sqrt{(E_T^{l1} + E_T^{\nu\nu})^2 - (\vec{p}_T^{l1} + \vec{p}_T^{miss})^2} \quad qqH \rightarrow WW^* \rightarrow l\nu l\nu$$



tt background

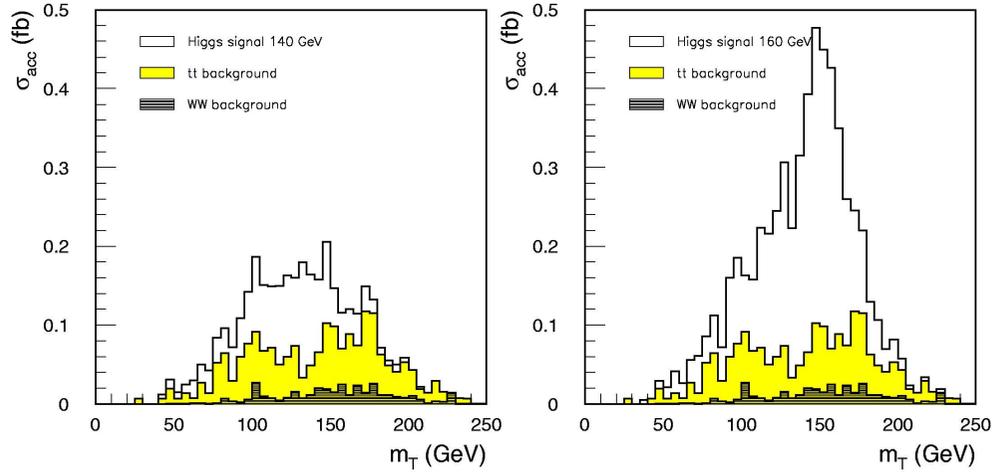
$\gamma^*/Z + \text{jets}$



el.weak WW jj

Higgs boson
 $m_H = 160 \text{ GeV}$

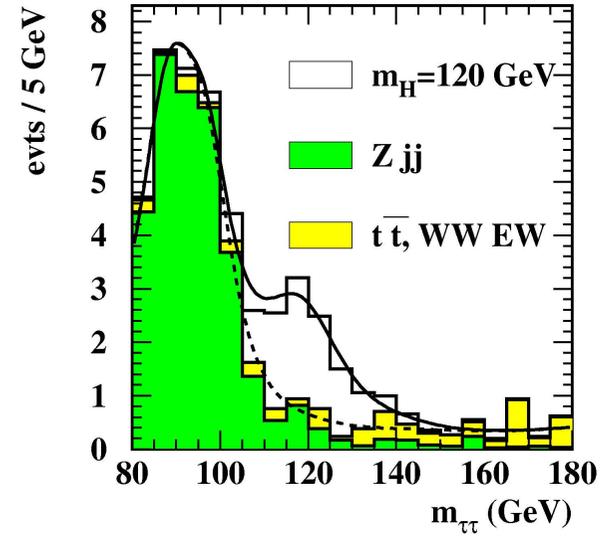
$qq H \rightarrow qq WW^* \rightarrow qq l \nu l \nu$



Number of expected events and signal significance for 5 fb⁻¹:

m_H	(GeV)	130	140	150	160	170	180
$H \rightarrow WW^{(*)} \rightarrow e\mu + X$							
Signal	(5 fb ⁻¹)	4.7	8.3	13.3	21.6	21.7	18.1
Background	(5 fb ⁻¹)	3.1	3.8	4.3	5.5	6.2	6.9
Stat. significance	(5 fb ⁻¹)	2.1	3.3	4.7	6.5	6.3	5.2
$H \rightarrow WW^{(*)} \rightarrow ee/\mu\mu + X$							
Signal	(5 fb ⁻¹)	4.4	8.3	14.1	20.4	22.8	18.3
Background	(5 fb ⁻¹)	4.2	4.7	5.5	6.4	7.3	7.9
Stat. significance	(5 fb ⁻¹)	1.8	3.0	4.6	6.0	6.2	5.1

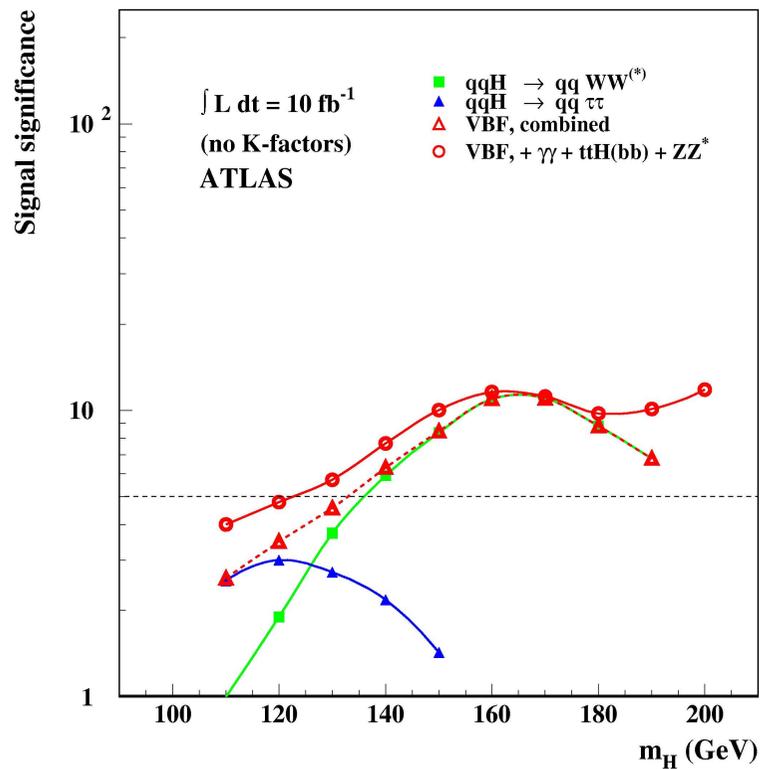
$qq H \rightarrow qq \tau \tau$



Number of expected events and signal significance for 30 fb⁻¹:

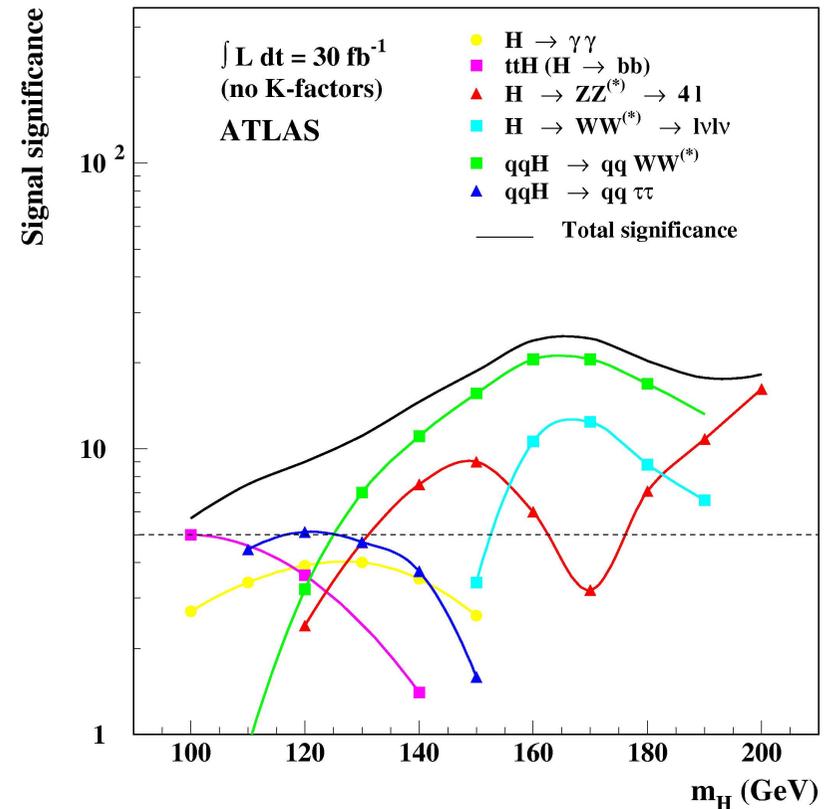
m_H	(GeV)	110	120	130	140	150
$H \rightarrow \tau\tau \rightarrow e\mu P_T^{miss}$						
Signal		7.7	7.0	5.1	3.3	1.5
Background		10.1	3.7	3.3	2.7	2.2
Stat. significance		2.1	2.8	2.2	1.6	-
$H \rightarrow \tau\tau \rightarrow ee/\mu\mu P_T^{miss}$						
Signal		9.2	7.2	5.7	3.1	1.5
Background		15.4	7.6	5.6	4.6	3.4
Stat. significance		2.1	2.2	2.0	1.2	-
$H \rightarrow \tau\tau \rightarrow l had P_T^{miss}$						
Signal		19	15.6	13	10	5
Background		27.0	11.7	10.6	7.4	6.7
Stat. significance		3.3	3.8	3.4	3.0	1.6
combined						
Stat. significance		4.3	5.1	4.4	3.6	2.1

Combined significance of VBF channels for 10 fb^{-1}



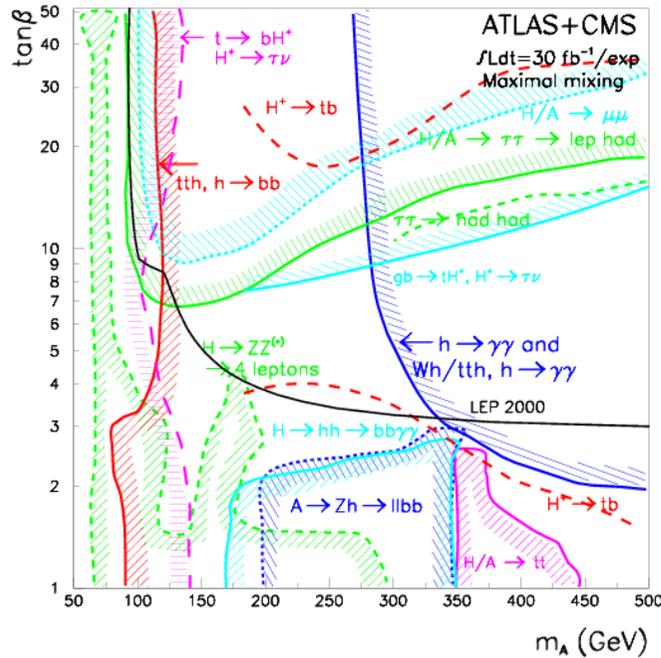
- Vector boson fusion channels (in particular WW^*) are discovery channels at low luminosity
- For 10 fb^{-1} in ATLAS: **5σ significance for $120 \leq m_H \leq 190 \text{ GeV}$**
(after combination with the standard channels)

ATLAS Higgs discovery potential for 30 fb^{-1}



- Vector boson fusion channels improve the sensitivity significantly in the low mass region
- Several channels available over the full mass range

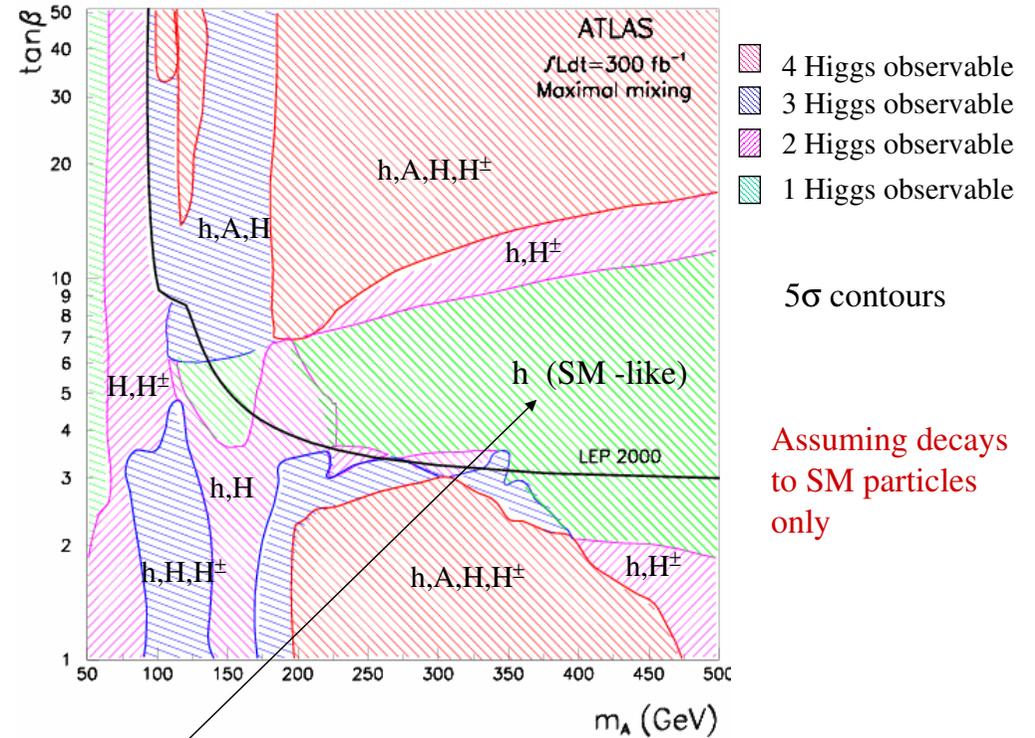
LHC discovery potential for MSSM Higgs bosons



Assuming SUSY particles are heavy

Not all channels shown

- Plane fully covered (no holes) at low L (30 fb^{-1})
- Main channels : $h \rightarrow \gamma\gamma, b\bar{b}, A/H \rightarrow \mu\mu, \tau\tau, H^\pm \rightarrow \tau\nu$
- Two or more Higgs can be observed over most of the parameter space \rightarrow disentangle SM / MSSM
- If LEP excess due to hZ production ($\tan\beta > 2, m_A > 115 \text{ GeV}$), LHC will observe:
 - h for any $\tan\beta$ and m_A
 - A, H, H^\pm for large $\tan\beta$ and moderate m_A

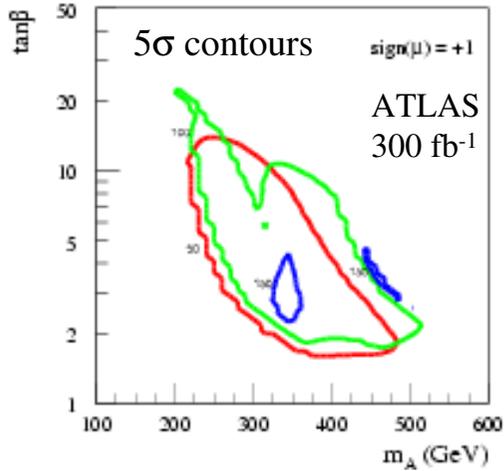


Here only SM-like h observable if SUSY particles neglected.

Higgs decays via SUSY particles

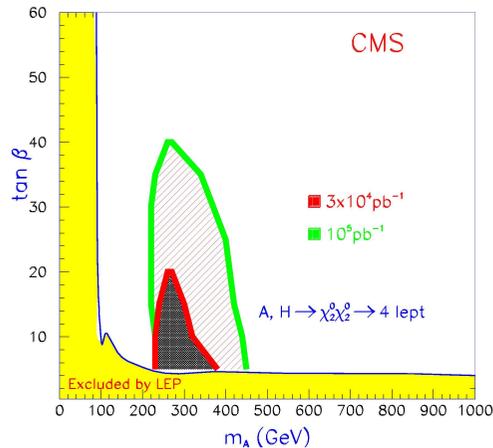
If SUSY exists : search for

$$H/A \rightarrow \chi_2^0 \chi_2^0 \rightarrow \ell \ell \chi_1^0 \ell \ell \chi_1^0$$



ATLAS:
SUGRA scan

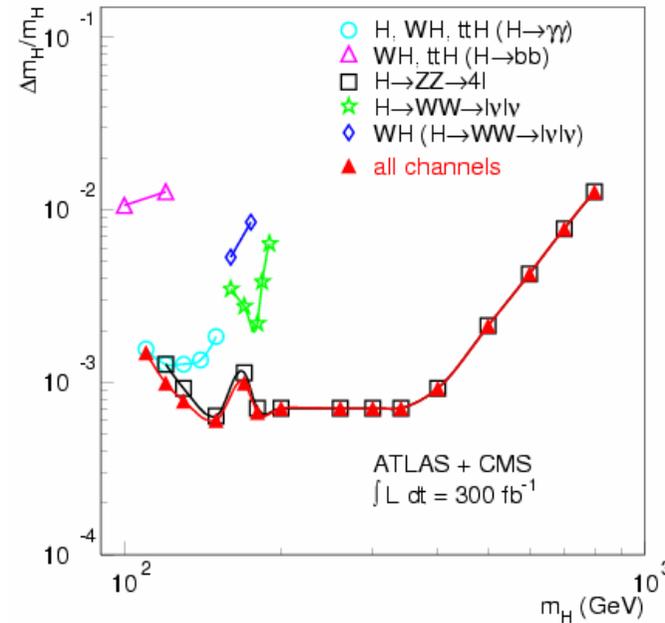
$$\begin{aligned} m_0 &= 50 - 250 \text{ GeV} \\ m_{1/2} &= 100 - 300 \text{ GeV} \\ \tan \beta &= 1.5 - 50 \\ A_0 &= 0 \end{aligned}$$



CMS:
special choice in MSSM
(no scan)
 $M_1 = 60 \text{ GeV}$
 $M_2 = 110 \text{ GeV}$
 $\mu = -500 \text{ GeV}$

Exclusions depend on MSSM parameters (slepton masses, μ)

Measurement of the Higgs boson mass



No theoretical error
e.g. mass shift for
large Γ_H (interference
resonant/non-resonant
production)

Dominant systematic
uncertainty: γ/ℓ E scale.
Assumed 1‰
Goal 0.2‰
Scale from $Z \rightarrow \ell\ell$
(close to light Higgs)

MSSM Higgs	$\Delta m/m$ (%)	300 fb ⁻¹
h, A, H $\rightarrow \gamma\gamma$	0.1–0.4	
H $\rightarrow 4\ell$	0.1–0.4	
H/A $\rightarrow \mu\mu$	0.1–1.5	
h $\rightarrow bb$	1–2	
H $\rightarrow hh \rightarrow bb\gamma\gamma$	1–2	
A $\rightarrow Zh \rightarrow bb\ell\ell$	1–2	
H/A $\rightarrow \tau\tau$	1–10	

Note: present theoretical error $\Delta m_h \sim 3 \text{ GeV}$

Measurements of Higgs boson couplings

i) Ratio between W and Z partial widths

- Direct measurements

$$- \frac{\sigma \times \text{BR}(H \rightarrow WW^*)}{\sigma \times \text{BR}(H \rightarrow ZZ^*)} = \frac{\Gamma_g \Gamma_W}{\Gamma_g \Gamma_Z} = \frac{\Gamma_W}{\Gamma_Z}$$

- QCD corrections cancel

- Indirect measurements (via $H \rightarrow \gamma\gamma$)

ii) Ratio of boson to fermion couplings

- Direct measurement

VBF:
$$- \frac{\sigma \times \text{BR}(qq \rightarrow qqH(H \rightarrow WW))}{\sigma \times \text{BR}(qq \rightarrow qqH(H \rightarrow \tau\tau))} = \frac{\Gamma_W \Gamma_W}{\Gamma_W \Gamma_\tau} = \frac{\Gamma_W}{\Gamma_\tau}$$

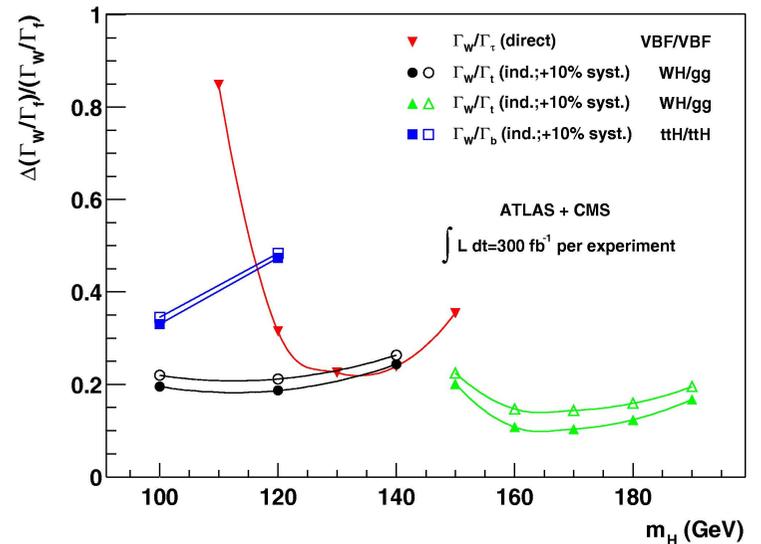
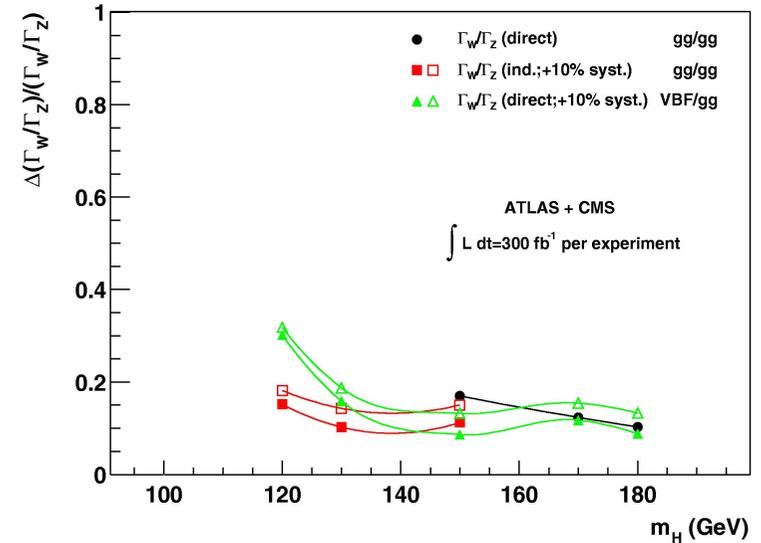
- Indirect measurement

$$- \frac{\sigma \times \text{BR}(WH(H \rightarrow \gamma\gamma))}{\sigma \times \text{BR}(H \rightarrow \gamma\gamma)} = \frac{\Gamma_W \Gamma_\gamma}{\Gamma_g \Gamma_\gamma} \sim \frac{\Gamma_W}{\Gamma_t} * C_{QCD}$$

$$- \frac{\sigma \times \text{BR}(WH(H \rightarrow WW))}{\sigma \times \text{BR}(H \rightarrow WW^*)} = \frac{\Gamma_W \Gamma_W}{\Gamma_g \Gamma_W} \sim \frac{\Gamma_W}{\Gamma_t} * C_{QCD}$$

$$- \frac{\sigma \times \text{BR}(ttH(H \rightarrow bb))}{\sigma \times \text{BR}(ttH(H \rightarrow \gamma\gamma))} = \frac{\Gamma_t \Gamma_b}{\Gamma_t \Gamma_\gamma} \sim \frac{\Gamma_b}{\Gamma_W}$$

* Uncertainties on the ratio arising through different production processes are not included



W-mass measurement

Physics motivation:

Test of the Standard Model: $m_Z, m_W, m_{\text{top}} \Rightarrow m_H$

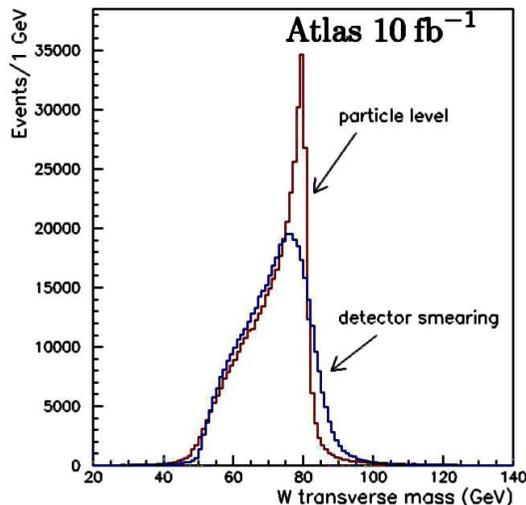
Year 2007: $\Delta m_W < 30 \text{ MeV}$ (LEP2 + Tevatron)

LHC goal: $\Delta m_W \sim 15 \text{ MeV}$ to match the precision on the top quark mass measurement

Experimental numbers:

- L dt = 10 fb⁻¹: 60 Mio. well measured $W \rightarrow \ell \nu$ decays
- Background conditions from pile-up events at low luminosity (2 events / bunch crossing) similar to Tevatron today
- Standard **transverse mass** technique can be used:

$$M_w^T = \sqrt{2 \cdot P_T^l \cdot P_T^\nu \cdot (1 - \cos \Delta\phi^{l,\nu})}$$



Estimate of Δm_W

Source of syst.	CDF Run 1b	ATLAS	Comments
Lepton scale	75 MeV	15 MeV	<40MeV at Run II
Lepton resolution	25 MeV	5 MeV	Known to <1.5%
$P_T(W)$	15 MeV	5 MeV	Constrain with $P_T(Z)$
Recoil model	37 MeV	5 MeV	Constrain with Z data
W width	10 MeV	7 MeV	
PDFs	15 MeV	< 10 MeV	Constraints from the LHC
Radiative decays	20 MeV	< 10 MeV	Theor. calculations
Total	92 MeV	< 25 MeV	per lepton species

- Total error per lepton species and per experiment is estimated to be **$\pm 25 \text{ MeV}$**
- Main uncertainty: lepton energy scale (goal is an uncertainty of $\pm 0.02 \%$)
- Many systematic uncertainties can be controlled in situ, using the $Z \rightarrow \ell\ell$ sample ($P_T(W)$, recoil model, resolution)

Combining both experiments (ATLAS + CMS), both lepton species and assuming a scale uncertainty of $\pm 0.02 \%$

$$\Rightarrow \Delta m_W \sim \pm 15 \text{ MeV}$$

Measurement of the Top Quark Mass

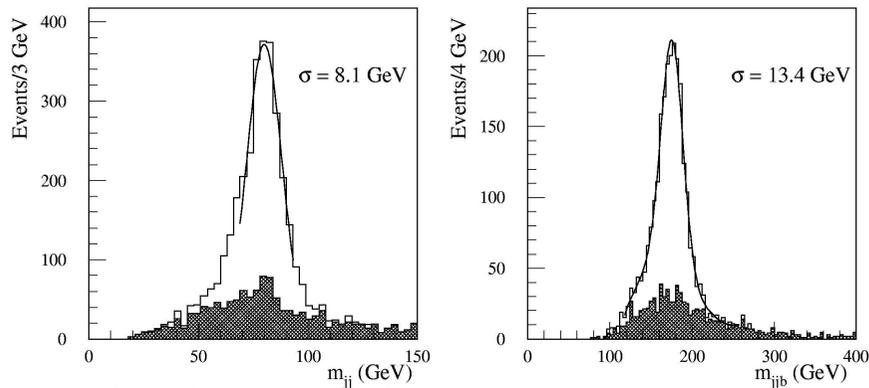
Year 2007: $\Delta m_{\text{top}} \sim 2\text{-}3 \text{ GeV}$ (Tevatron)

Best channel for mass measurement:

$tt \rightarrow Wb \quad Wb \rightarrow \ell \nu b \quad \text{jet jet } b$
 (trigger) (mass measurement)

Experimental numbers:

- Production cross section: 590 pb
- After exp. cuts: 130.000 tt events in 10 fb^{-1} S/B ~ 65



results from full detector simulation

Contribution	Δm_{top} (GeV)
statistics	< 0.07
u,d,s jet scale	0.3
b-jet scale	0.7
b-fragmentation	0.3
initial state rad.	0.3
final state rad.	1.2
background	0.2
Total	~ 1.5 GeV

Syst. uncertainties dominated by final state radiation

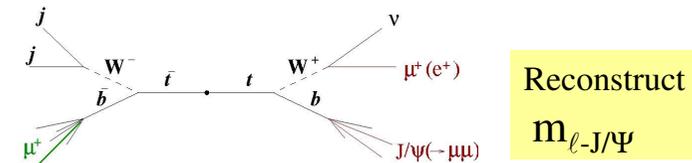
Additional Methods

- Full reconstruction applying kinematical constraints

$$m_{jj} = m_{\ell\nu} = m_W \quad \text{and} \quad m_{jbb} = m_{\ell\nu b}$$

Precision of $\sim \pm 1 \text{ GeV}$ can be reached

- Using $\ell\text{-}J/\psi$ final states:



- BR = 10^{-5} : low rate, but clean signature
- Statistical error: $\pm 0.9 \text{ GeV}$ (for 500 fb^{-1})
- Different systematic uncertainties (dominated by b-fragmentation: $\sim 0.4 \text{ GeV}$)

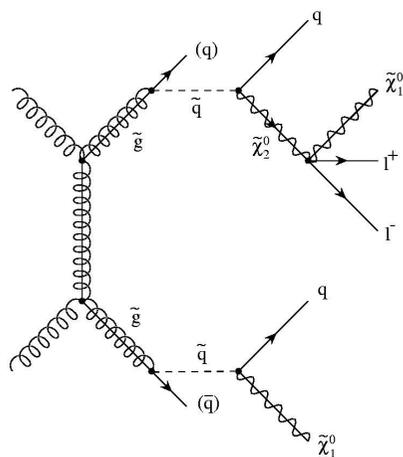
combination of various methods:

$$\Delta m_{\text{top}} < \sim \pm 1 \text{ GeV}$$

Search for Supersymmetry

- If **SUSY** exists at the electroweak scale, a discovery at the LHC should be easy
- **Squarks** and **Gluginos** are strongly produced

They decay through cascades to the lightest SUSY particle (LSP)



⇒ combination of

Jets, Leptons, E_T^{miss}

1. Step: Look for **deviations from the Standard Model**

Example: Multijet + E_T^{miss} signature

2. Step: Establish the **SUSY mass scale**

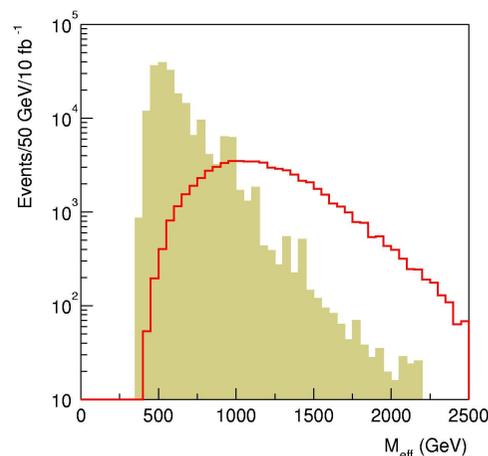
use inclusive variables, e.g. effective mass distribution

3. Step: Determine **model parameters** (difficult)

Strategy: select particular decay chains and use kinematics to determine mass combinations

Squarks and Gluginos

- strongly produced, cross sections comparable to QCD cross sections at same Q^2
- If R-parity conserved, cascade decays produce distinctive events: **multiple jets, leptons, and E_T^{miss}**
- Typical selection: $N_{\text{jet}} > 4$, $E_T > 100, 50, 50, 50$ GeV
 $E_T^{\text{miss}} > 100$ GeV
- Define: $M_{\text{eff}} = E_T^{\text{miss}} + P_T^1 + P_T^2 + P_T^3 + P_T^4$
(effective mass)



example: mSUGRA
 $m_0 = 100$ GeV
 $m_{1/2} = 300$ GeV
 $\tan \beta = 10$
 $A_0 = 0, \mu > 0$

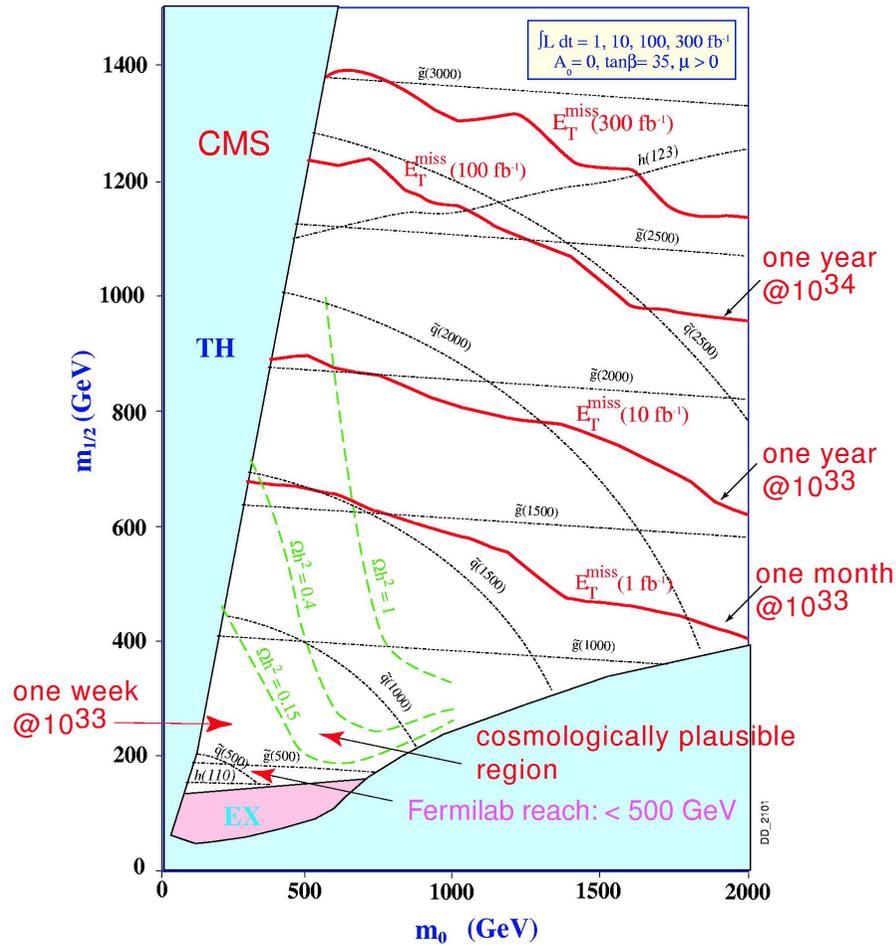
- LHC reach for Squark- and Gluino masses:

$$\begin{aligned} 1 \text{ fb}^{-1} &\Rightarrow M \sim 1500 \text{ GeV} \\ 10 \text{ fb}^{-1} &\Rightarrow M \sim 1900 \text{ GeV} \\ 100 \text{ fb}^{-1} &\Rightarrow M \sim 2500 \text{ GeV} \end{aligned}$$

TeV-scale SUSY can be found quickly !

SUSY mass scale

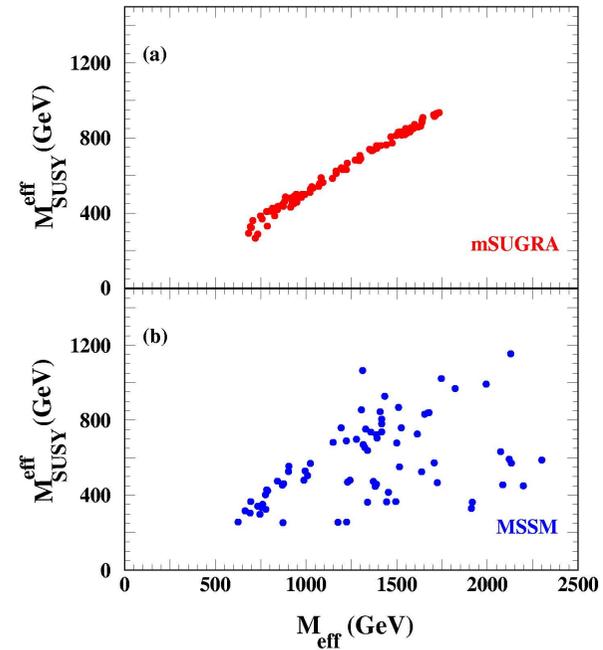
LHC reach in $m_0 - m_{1/2}$ mSUGRA plane:



- define average produced SUSY mass

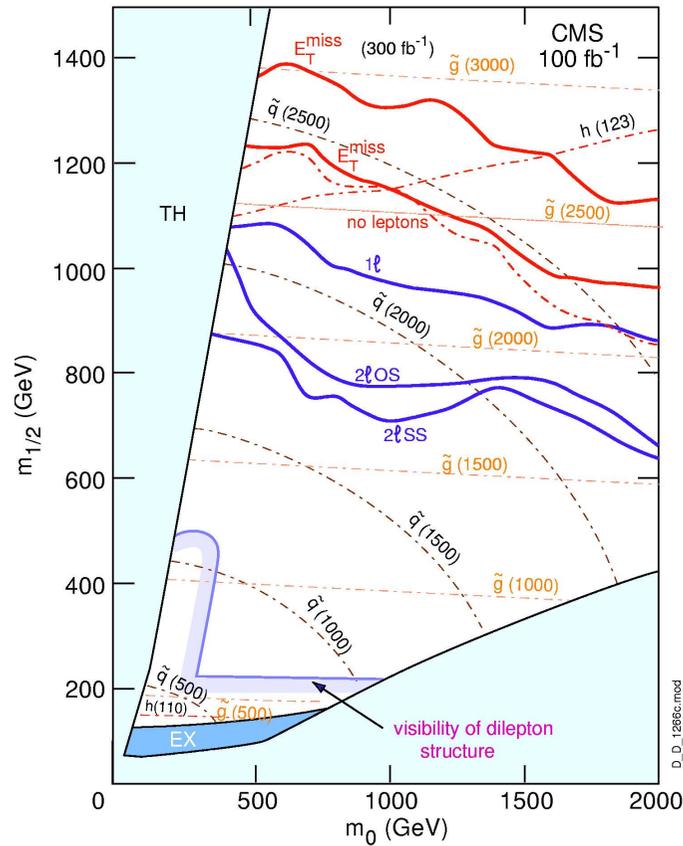
$$M_{\text{SUSY}} \equiv \frac{\sum_i M_i \sigma_i}{\sum_i \sigma_i}$$

$$M_{\text{SUSY}}^{\text{eff}} \equiv M_{\text{SUSY}} - \frac{M^2(\tilde{\chi}_1^0)}{M_{\text{SUSY}}}$$



- Good correlation with M_{eff} for mSUGRA
- Not bad even for MSSM

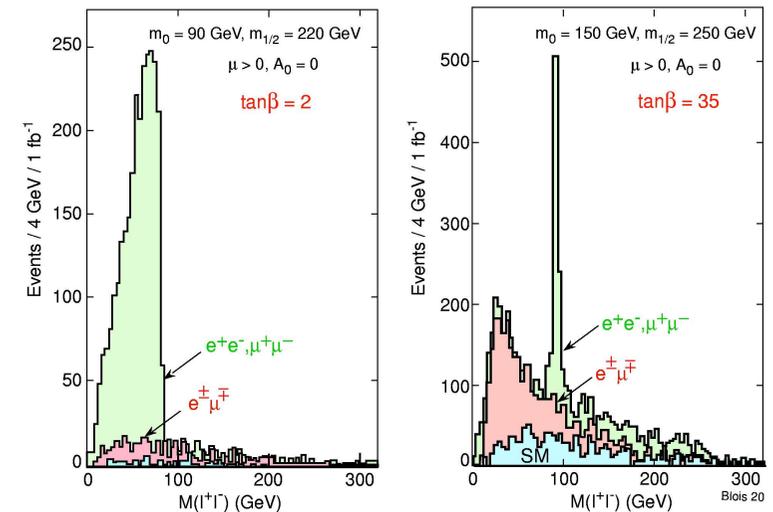
SUSY cascade decays give rise to many inclusive signatures: leptons, b-jets, τ 's



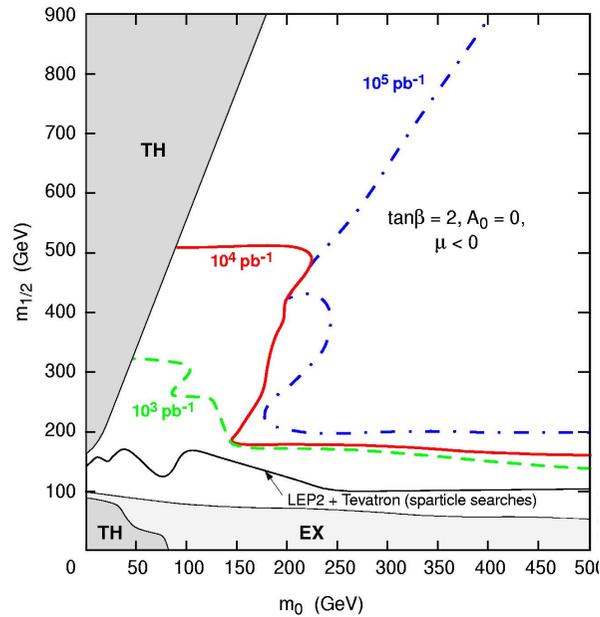
Expect multiple signatures for TeV-scale SUSY

Determination of model parameters

- **Invisible LSP** \Rightarrow no mass peaks, but kinematic endpoints
 \Rightarrow mass combinations
- Simplest case: $\chi^0_2 \rightarrow \chi^0_1 l^+ l^-$
 endpoint: $M_{ll} = M(\chi^0_2) - M(\chi^0_1)$
- Significant mode if no $\chi^0_2 \rightarrow \chi^0_1 Z, \chi^0_1 h, \tilde{l} l$ decays
- Require: 2 isolated leptons, multiple jets, and large E_T^{miss}



- Modes can be distinguished using shape of ll -spectrum



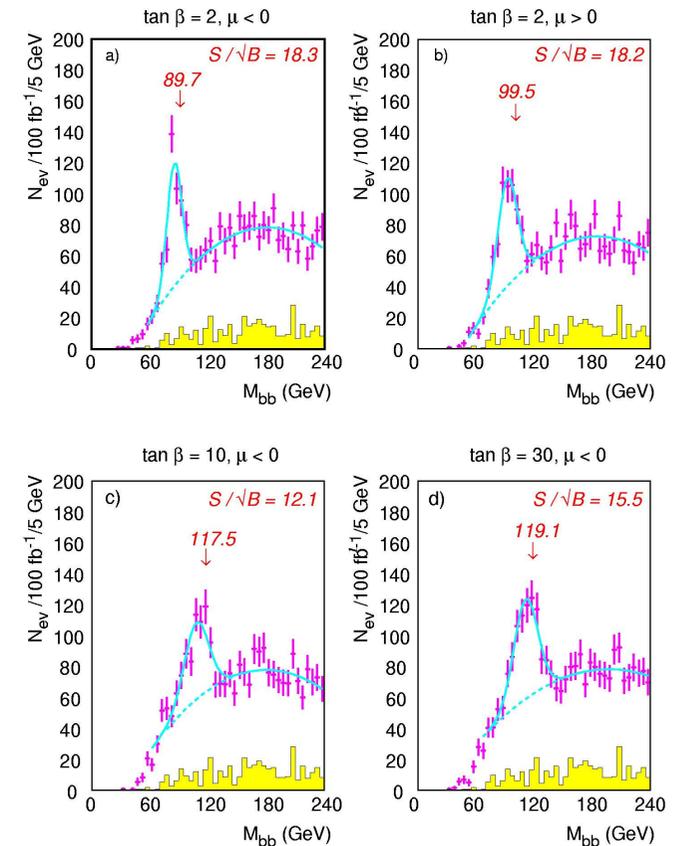
ll - endpoint can be observed over a significant fraction of the parameter space

(covers part of the SUGRA region favoured by cold dark matter (Ellis et al.))

$h \rightarrow bb$:

important if $\chi^0_2 \rightarrow \chi^0_1 h$ is open;

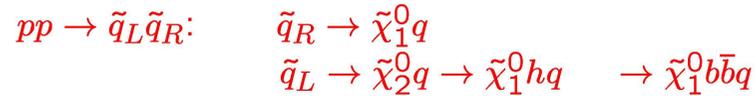
bb peak can be reconstructed in many cases



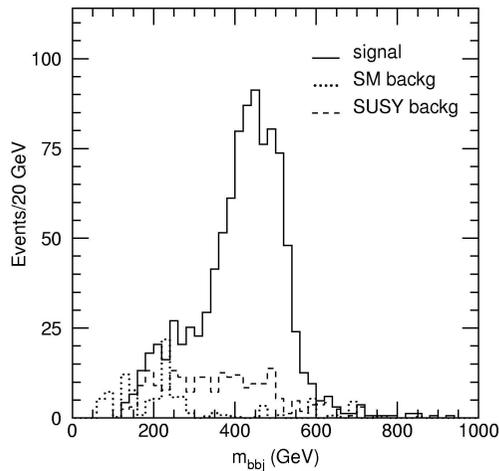
Could be a Higgs discovery mode !

SM background can be reduced by applying a cut on E_t^{miss}

work backwards the decay chain:
example: **SUGRA study point 5**



combine $h \rightarrow bb$ with jets to determine other masses



$\tilde{q} \rightarrow \tilde{\chi}_1^0 h q$ endpoint

Strategy in SUSY Searches at the LHC:

- Search for multijet + E_T^{miss} excess
- If found, select SUSY sample (simple cuts)
- Look for special features (γ 's, long lived sleptons)
- Look for l^\pm , $l^+ l^-$, $l^\pm l^\pm$, b-jets, τ 's
- End point analyses, global fit

Models other than SUGRA

GMSB:

- LSP is light gravitino
- Phenomenology depends on nature and lifetime of the NLSP
- Generally longer decay chains, e.g.



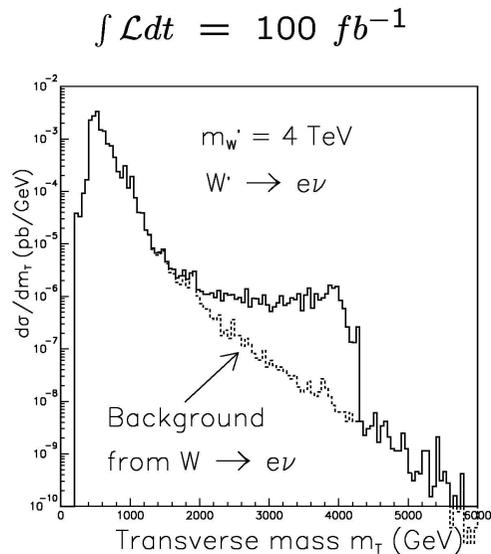
- \Rightarrow models with prompt NLSP decays give add. handles and hence are easier than SUGRA
- NLSP lifetime can be measured:
 - For $\tilde{\chi}_1^0 \rightarrow \tilde{G} \gamma$, use Dalitz decays (short lifetime) or search for non-pointing photons
 - Quasi stable sleptons: muon system provides excellent „Time of Flight“ system

RPV :

- R-violation via $\chi_1^0 \rightarrow ll\nu$ or $qq\ell$, $qq\nu$ gives additional leptons and/or E_T^{miss}
- R-violation via $\chi_1^0 \rightarrow cds$ is probably the hardest case; (c-tagging, uncertainties on QCD N-jet background)

Beyond SUSY, a few examples

Excited quarks: $q^* \rightarrow q\gamma$, up to:	$m \sim 6 \text{ TeV}$
Leptoquarks, up to:	$m \sim 1.5 \text{ TeV}$
Monopoles: $pp \rightarrow \gamma\gamma pp$, up to:	$m \sim 20 \text{ TeV}$
Lepton flavour viol. $\tau \rightarrow \mu\gamma$:	$10^{-6} - 10^{-7}$
Compositeness, up to: from di-jet and Drell-Yan, needs calorimeter linearity better than 2%	$\Lambda \sim 40 \text{ TeV}$
$Z' \rightarrow ll, jj$, up to:	$m \sim 5 \text{ TeV}$
$W' \rightarrow l\nu$, up to:	$m \sim 6 \text{ TeV}$



Search for Signals from Extra Dimensions

- Much recent theoretical interest in models with extra dimensions
- New physics can appear at the TeV-mass scale, i.e. accessible at the LHC
- **Gravitons** propagating in the extra dimensions will appear as massive states

Examples of searches:

(1) Search for direct graviton production

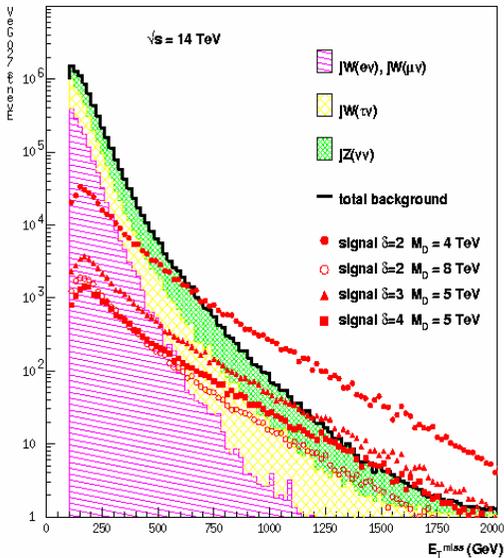
$$gg \rightarrow gG, q\bar{q} \rightarrow q\bar{q}G$$

$$q\bar{q} \rightarrow G\gamma$$

\Rightarrow Jets or Photons with E_T^{miss}

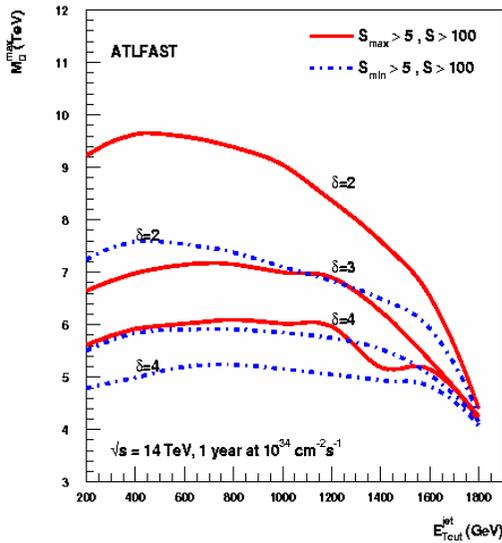
(2) Search for graviton resonances
(Randall Sundrum models)

Search for Graviton Production



Jet + E_T^{miss} search:

Main backgrounds:
 $\text{jet} + Z(\rightarrow \nu\nu), \text{jet} + W \rightarrow \text{jet} + (e, \mu, \tau)\nu$



$$G_N^{-1} = 8\pi R^\delta M_D^{2+\delta}$$

δ : # extra dimensions
 M_D = scale of gravity

for 100 fb^{-1} :

$M_D^{\text{max}} = 9.1, 7.0, 6.0 \text{ TeV}$
 for
 $\delta = 2, 3, 4$

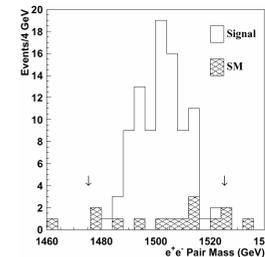
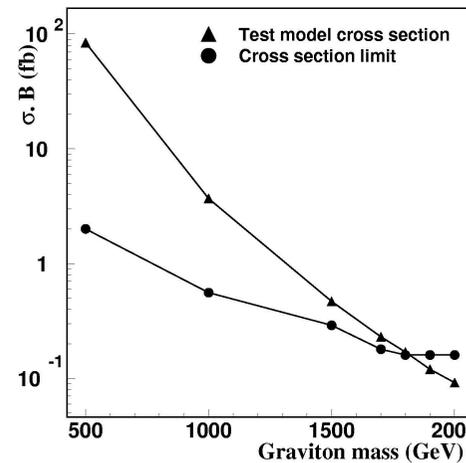
$(\gamma + E_T^{\text{miss}}$ search is less sensitive)

Search for Narrow Graviton Resonances

- use **Randall Sundrum model** as reference model:
- Kaluza-Klein graviton spectrum with a scale
 $\Lambda_\pi = M_{\text{Planck}} \exp(-k\pi r_c)$
- Properties of the model are determined by the ratio k/M_{Planck}

Atlas and CMS studies on sensitivity to narrow resonance states decaying into lepton pairs:

$$gg (qq) \rightarrow G \rightarrow e^+e^-$$



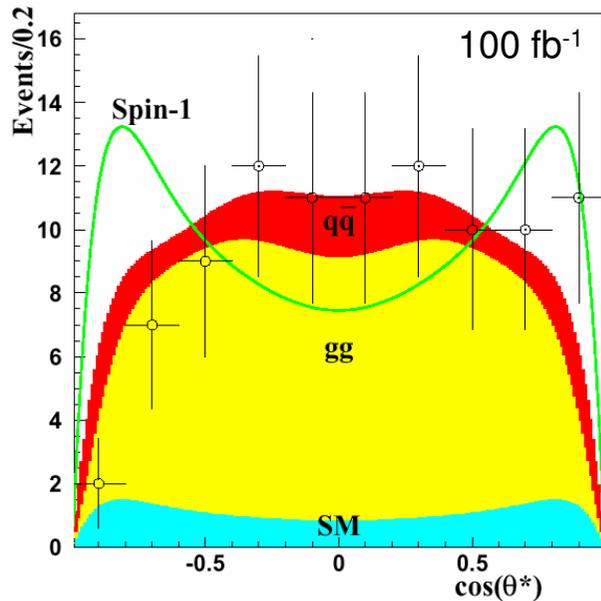
For $k / M_{\text{Planck}} = 0.01$
 (conservative choice)

With 100 fb^{-1} , signal can be seen in the mass range

$0.5 < M < 2.08 \text{ TeV}$

Spin determination:

from di-lepton angular distribution



acceptance effects included;
use likelihood method to discriminate between
spin-1 and spin-2 hypotheses

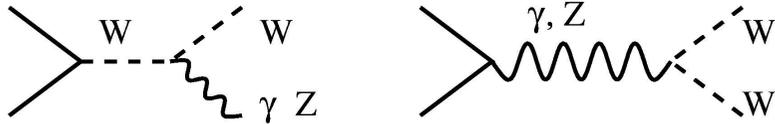
Spin determination possible up to $M \sim 1.7$ TeV
(100 fb^{-1} , 90%CL)

Conclusions

1. The pp experiments at the LHC have a huge discovery potential
 - **SM Higgs**: full mass range, already at low lumi; Vector boson fusion channels improve the sensitivity significantly
 - **MSSM Higgs**: parameter space covered; new benchmark scenarios investigated at present
 - **SUSY**: discovery of TeV-scale SUSY should be easy, determination of model param. is more difficult
 - **Exotics**: experiments seem robust enough to cope with new scenarios
2. Experiments have also a great potential for precision measurements
 - m_W to ~ 15 MeV
 - m_{top} to ~ 1 GeV
 - $\Delta m_H / m_H$ to 0.1% (100 - 600 GeV)

+ gauge couplings and measurements in the top sector

Triple Gauge Boson Couplings



- Probe non-Abelian structure of $SU(2) \times U(1)$ and sensitive to **New Physics**
- general assumptions (Lorentz invariance, P,C inv.):
 $\Rightarrow WW\gamma$ and WWZ couplings specified by five parameters: $g_1^Z, \lambda_\gamma, \lambda_Z, \kappa_\gamma, \kappa_Z$

$WW\gamma$ -vertex: related to

- magnetic moment $\mu_W = \frac{e}{2M_W} (g_1^Z + \kappa_\gamma + \lambda_\gamma)$
- quadrupole moment $Q_W = -\frac{e}{M_W^2} (\kappa_\gamma - \lambda_\gamma)$

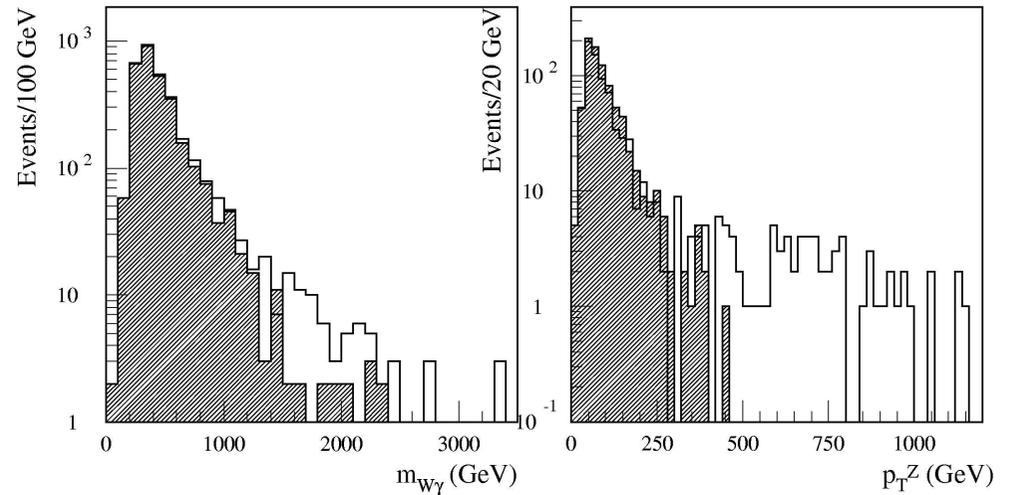
Standard Model: $g_1^Z = \kappa_V = 1$
 $\lambda_V = 0$

year 2005: known to better than 10^{-2} from LEP2+TeVatron

- $W\gamma \rightarrow l\nu\gamma$ studied
- $WZ \rightarrow l\nu ll$ studied
- $WW \rightarrow l\nu l\nu$ large $t\bar{t}$ background
- Sensitivity from:
 - cross section measurements: λ -type, increase with s
 - P_T and angular distributions: constrain κ -type

$W\gamma$ 30 fb⁻¹ : ~3000 events
 $\lambda_\gamma = 0.01$

WZ 30 fb⁻¹ : ~1200 events
 $\Delta g_Z^1 = 0.05$



$\int \mathcal{L} dt = 30 \text{ fb}^{-1}$

Coupling	95% C.L.
Δg_Z^1	0.008
λ_γ	0.0025
λ_Z	0.0060
$\Delta \kappa_\gamma$	0.035
$\Delta \kappa_Z$	0.070

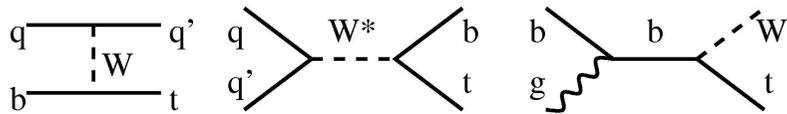
Systematics under study

Other measurements in top physics

- Cross section measurement, $\sigma_{t\bar{t}} < 10\%$
(limited by uncertainty on luminosity)
- Sensitivity to FCNC top couplings:

$BR(t \rightarrow Zq) < 10^{-4}$	$\int \mathcal{L} dt = 100 \text{ fb}^{-1}$
$BR(t \rightarrow \gamma q) < 10^{-4}$	5 σ discovery limit
$BR(t \rightarrow gq) < 7 \cdot 10^{-3}$	5 σ discovery limit
	95% C.L.

- Single Top production: $\sigma \sim 300 \text{ pb}$ (40% of $t\bar{t}$)



- probe $W - tb$ vertex,
→ sensitive to new physics
- measure V_{tb} to $\sim 10\%$ (syst. limited)
- measure W, **top polarisation**
→ anomalous couplings,