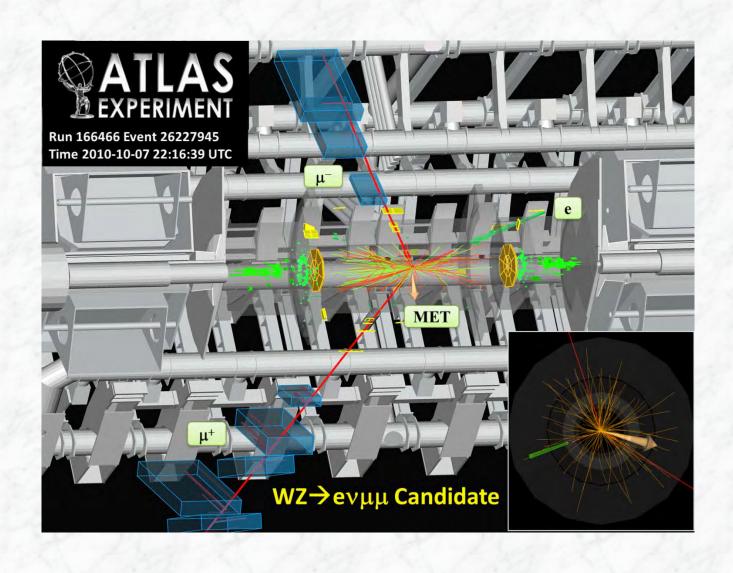
Physics at Hadron Colliders



1. Introduction, Particle Accelerators

- 1.1 Why hadron collider?
- 1.2 Principles on particle accelerators
- 1.3 The Large Hadron Collider (LHC) at CERN
- 1.4 The Tevatron Collider at Fermilab

1.1 Why Hadron Collider?

Key questions investigated in particle physics:

(i) Structure of matter; fundamental constituents

$$\lambda = h/p$$

(ii) Search for new particles, new types of matter

$$E = m c^2$$

For the investigation of both questions, high energies and thereby particle accelerators are needed

High Energy Particle Accelerators (last 20 years):

Accelerator	type, laboratory	energy √s	years of operation
LEP-I	e+e- collider, CERN	91 GeV	1989 - 1994
LEP-II	e+e- collider, CERN	209 GeV	1995 - 2000
HERA-I	ep collider, DESY	27 + 800 GeV	1992 - 2000
HERA-II	ep collider, DESY	27 + 920 GeV	2002 - 2007
TeVatron Run I	ppbar collider, Fermilab	1.8 TeV	1987 - 1996
TeVatron Run II	ppbar collider, Fermilab	1.96 TeV	2002 - 2011
LHC, phase I	pp collider, CERN	7 TeV	2010- 2012
LHC, phase II	pp collider, CERN	14 TeV	2015

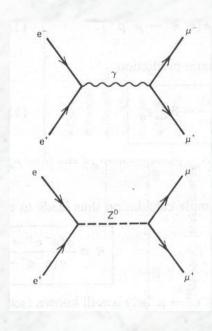
Important parameters of accelerators:

- •Beam energy, centre-of-mass energy √s
- Type of particles (ee, ep, or pp) and form of accelerator (circular or linear accelerator)
- •Luminosity L, or integrated Luminosity (measured in units of cm⁻² s⁻¹)

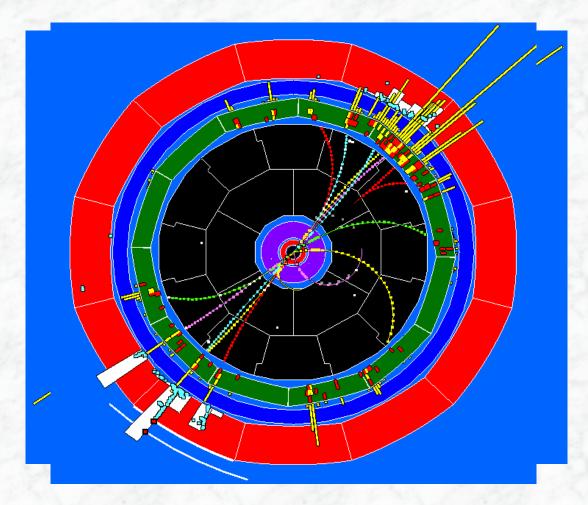
Why a hadron collider?

e⁺e⁻ colliders are excellent machines for precision physics!!

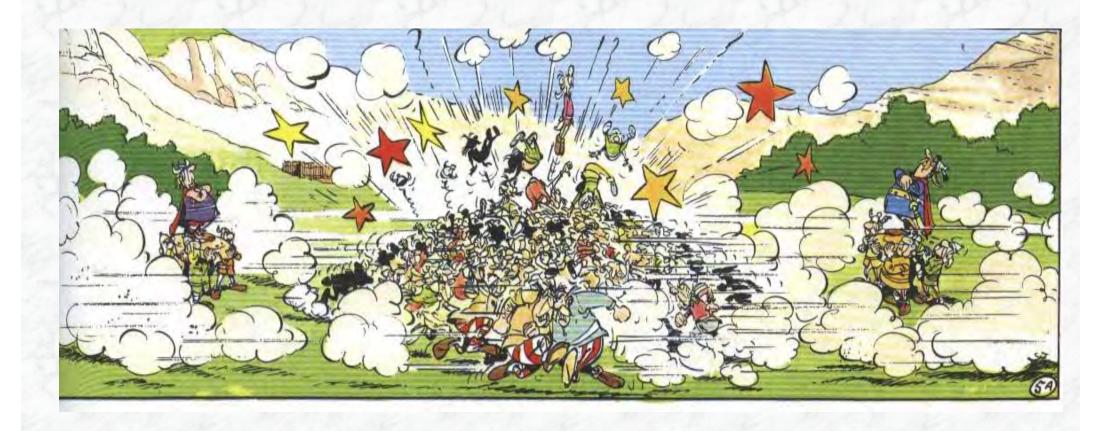
- e+ e⁻ are point-like particles, no substructure → clean events
- complete annihilation, centre-of-mass system, kinematic fixed



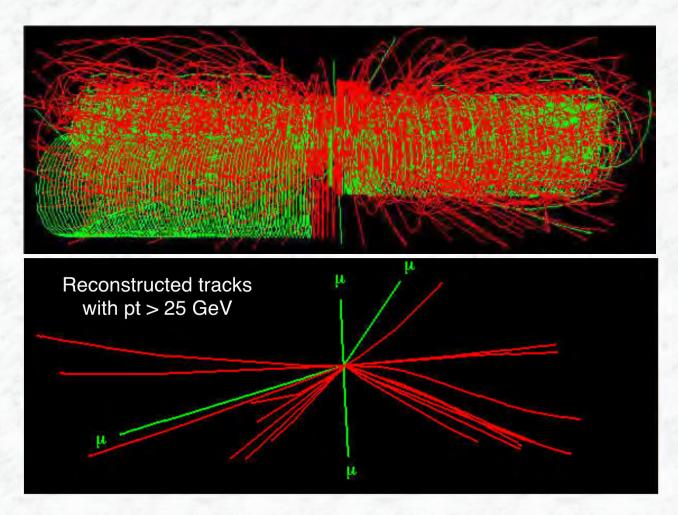
K. Jakobs



Proton proton collisions are more complex



Simulation of a pp collision at the LHC: $\sqrt{s} = 14 \text{ TeV}$, $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$



Reconstruction of particles with high transverse momentum reduces the number of particles drastically (interesting object largely kept, background from soft inelastic pp collisions rejected)

Main drawbacks of e⁺e⁻ circular accelerators:

- Energy loss due to synchrotron radiation
 (basic electrodynamics: accelerated charges radiate, x-ray production via bremsstrahlung, synchrotron radiation.....)
 - Radiated power (synchrotron radiation):
 Ring with radius R and energy E
 - Energy loss per turn:(2 GeV at LEP-II)
 - Ratio of the energy loss between protons and electrons:

$$P = \frac{2 e^2 c}{3 R^2} \left(\frac{E}{mc^2}\right)^4$$

$$-\Delta E \approx \frac{4 \pi e^2}{3 R} \left(\frac{E}{mc^2}\right)^4$$

$$\frac{\Delta E(e)}{\Delta E(p)} = \left(\frac{m_p}{m_e}\right)^4 \sim 10^{13}$$

Future accelerators:

- pp ring accelerators (LHC, using existing LEP tunnel)
- or e⁺e⁻ linear accelerators, International Linear Collider ILC or CLIC (under study / planning)

Limiting factors:

e+ e- accelerators:

- -Energy loss is circular rings
- -Acceleration gradient in linear accelerators (ILC design, 40 MV / m)
- Fixed centre-of-mass energy....

pp accelerators:

- -More complex interactions due to proton substructure
- -Only part of the pp centre-of-mass energy available in the hard scattering process (see later)
 - however: higher mass values can be reached with longer running times
- -Magnetic field in bending magnets (8.3 T in LHC magnets)

Accelerators at the energy frontier

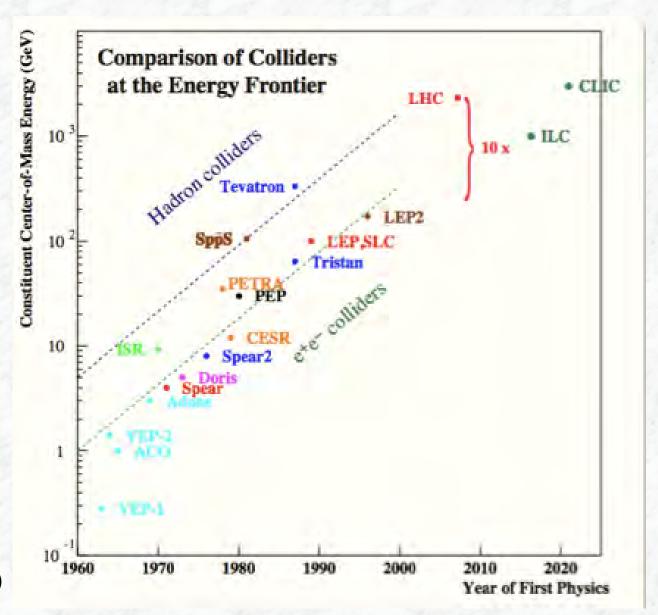
Livingston plot

Exponential growth of √s with time (at least in the past)

Factor 4 every 10 years

pp: discovery machines (W/Z, top, Higgs,....)

e+e-: precision (LEP, QCD and el.weak)



Both required!

Today's open questions in particle physics





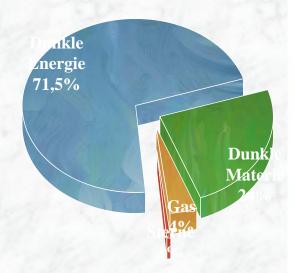
Key questions of particle physics

1. Mass

What is the origin of mass?

Does the Higgs particle exist?

Much progress in the past 2 years



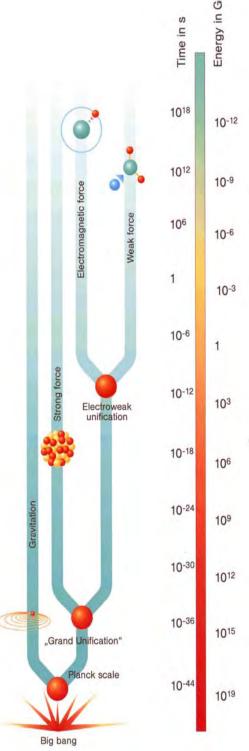
2. Unification

- Can the interactions be unified?
- Are there new types of matter,
 e.g. supersymmetric particles?
 Are they responsible for the Dark Matter in the universe?

3. Flavour

- Why are there three generations of particles?
- What is the origin of the matter-antimatter asymmetry (Origin of CP violation)

Answers to some of these questions are expected on the TeV energy scale, i.e. at the LHC



The role of the LHC

1. Explore the TeV mass scale

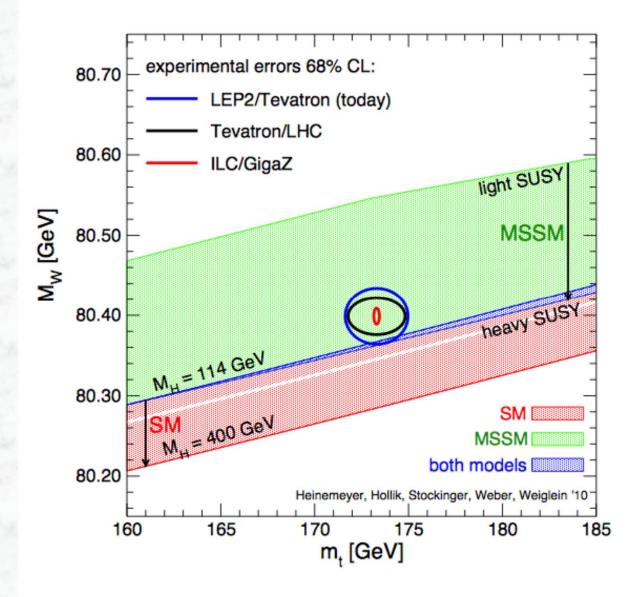
- What is the origin of the electroweak symmetry breaking?
 Does the Higgs boson exist?
- The search for "low energy" supersymmetry
 Can a link between SUSY and dark matter be established?
- Other scenarios beyond the Standard Model
-

Look for the "expected", but we need to be open for surprises

→ perform as many searches (inclusive, exclusive...) for as many final states as possible

2. Precise tests of the Standard Model

- There is much sensitivity to physics beyond the Standard Model in the precision area
- Many Standard Model measurements can be used to test and to tune the detector performance

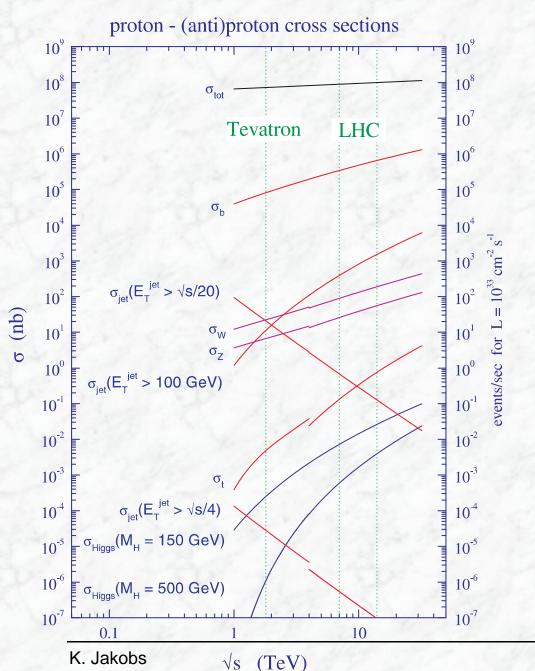


Predictions for future precision (including LHC), compared to the Standard Model and its Minimal Supersymmetric Extension (MSSM)

Ultimate test of the Standard Model: compare direct prediction of Higgs mass with direct observation



Production Rates and Cross Sections at the LHC



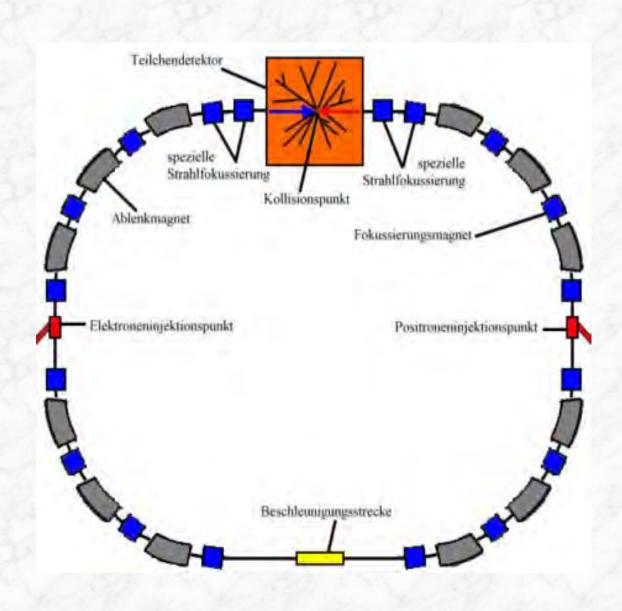
$$N = \sigma \cdot L$$

Rates for the design luminosity: $\sqrt{s} = 7 \text{ TeV}$, L = $10^{33} \text{ cm}^{-2} \text{ s}^{-1}$:

Higgs (150 GeV)Gluino, Squarks (1 TeV)	0.02 /s 0.003 /s
• W \rightarrow e v • Z \rightarrow e e	15 /s 1.5 /s
bb pairstt pairs	5 10 ⁵ /s 1 /s
 Inelastic proton-proton collisions: 	10 ⁸ /s

Hadron Collider Physics, Freiburg, SS 2013

1.2 Principles of particle accelerators



Circular accelerator principles

• Cyclotron: constant RF

magnetic field radius ρ increases with energy

used for smaller machines



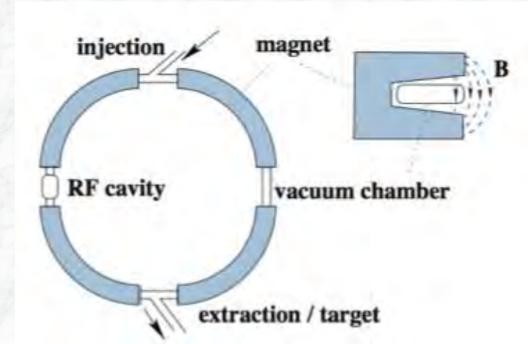
Synchrotron: $\rho = const$

B increases with energy

RF frequency adjusted slightly ($\beta = 0.999 \dots 1.0$)

Most High Energy accelerators and all CERN ring accelerators

(PS, SPS, LEP, LHC) are of this type



Basic parameters, Lorentz Force

$$\mathbf{F} = \mathbf{q} \left(\mathbf{E} + \mathbf{v} \times \mathbf{B} \right)$$

charge q, normally q = e; q = Ze for ions

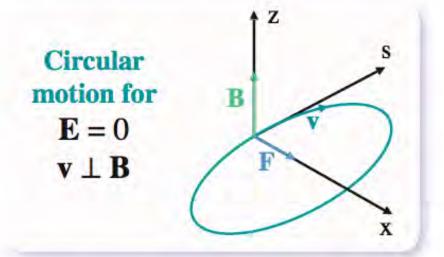
- Electric field E provides the acceleration or rather energy gain
- The magnetic field B keeps the particles on their path

 ρ is the radius of curvature for motion perpendicular to the static magnetic field. Often called

- gyromagnetic or Larmor radius in astroparticle physics
- bending radius for accelerators
 Bp known as magnetic rigidity, units Tm

LHC

- Momentum p = 7 TeV/c
- LHC bending radius $\rho = 2804 \text{ m}$
- Bending field B = 8.33 Tesla
- magnets at 1.9 K, super-fluid He



$$B = \frac{p}{q \rho} \quad \begin{array}{l} \text{for } q = e \text{ numerically} \\ B [T] = p [GeV/c] \quad 3.336 \text{ m} / \rho \\ \text{high energy, } v = c \text{ "p} = E" \\ E < E_H = q B \rho \text{ Hillas criterion} \end{array}$$

Astroparticle

units
$$10^{-4}T = 1Gauss$$
; a.u. = $1.5 \times 10^{11}m$
Solar system B = $10\mu G$ E = $5 \text{ TeV } \rho = 11 \text{ a.u.}$
Intergalactic B = $1nG$ E = 5 PeV (knee)
 $\rho = 1.7 \times 10^{19}m$ (4 % of galaxy-radius)

Luminosity and collision rates

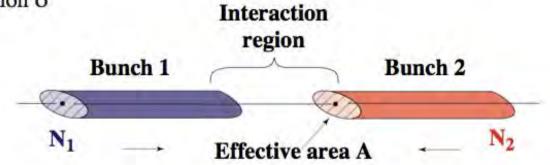
Event rate for process with cross section of

$$\dot{n} = \mathcal{L}\sigma$$

Luminosity from bunch

crossings at frequency $f = f_{rev} n_b$

$$\mathcal{L} = \frac{N_1 N_2 f}{A}$$



for Gaussian bunches with rms sizes $\sigma_x \sigma_y$ $A = 4 \pi \sigma_x \sigma_y$

High luminosity: Large number of particles (N₁, N₂)

Small beam dimensions (A) in the interaction point

Large f (large number of bunches,

→ small time difference between bunch crossings)

LHC: $N = 1.15 \ 10^{11}$, $n_b = 2808$ (number of bunches)

bunch separation: 25 ns (corresponds to 7.5 m)

A: beams squeezed (using strong, large aperture quadrupoles close to the interaction region) from σ = 0.2 mm to 16 μ m

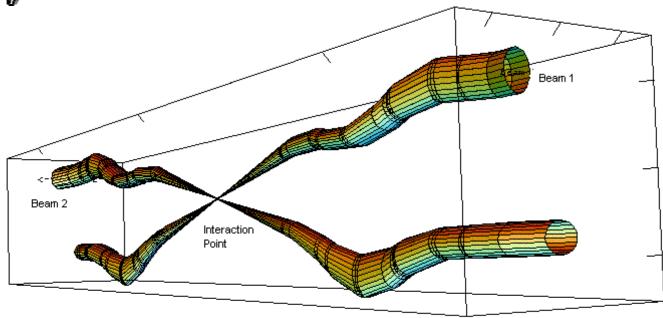
$$L = f \frac{N_1 \ N_2}{4\pi \ \sigma_x \sigma_y} = f \frac{N}{4\pi \ \sigma^2}$$

for N = N₁ = N₂ particles per bunch with transverse r.m.s. beam size σ = σ_x = σ_y and frequency f

Accelerator physicist express this often using the transverse emittance ϵ and the β function:

$$L = f \frac{N_1 N_2}{4 \sqrt{\epsilon_x \beta_x^* \epsilon_y \beta_y^*}}$$

where
$$\epsilon=\pi\sigma^2/eta$$



Relative beam sizes around IP1 (Atlas) in collision

Alternate gradient focusing

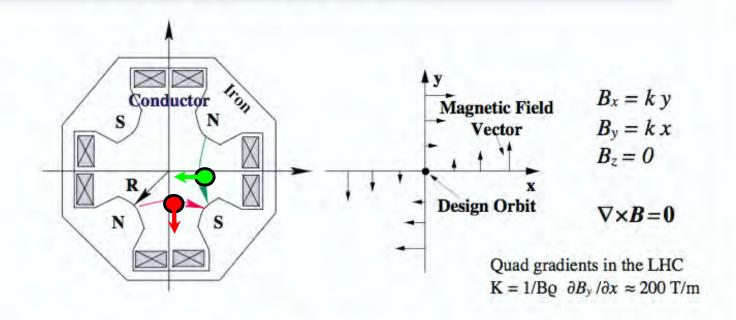
Quadrupole lens focusing in x, defocusing in y or vice versa

$$\mathbf{F} = \mathbf{e} (\mathbf{v} \times \mathbf{B})$$
here
$$\mathbf{F} = \mathbf{e} (0, 0, \mathbf{v}) \times (B_x, B_y, 0)$$

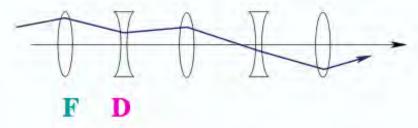
$$= \mathbf{e} (-\mathbf{v} B_y, +\mathbf{v} B_x, 0)$$

Combine F D
Defocusing when at small amplitude
Overall focusing

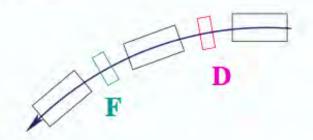
Normal (light) optics: Focal length of two lenses at distance D $1/f = 1/f_1 + 1/f_2 - D/f_1f_2$ is overall focusing with $1/f = D/f^2$ for $f = f_1 = -f_2$



alternate gradient focusing



together with bending magnets FODO lattice



N. C. Christofilos, unpublished manuscript in 1950 and patent Courant, Snyder in 1952, Phys. Rev. 88, pp 1190 - 1196 + longer review in Annals of Physics 3 (1958)

Betatron motion

Equation of motion of particles in a ring (with bending fields) and quadrupoles (field gradients \approx \particles B/\particles r)

In both transverse planes, here written with x for x, y: x' = dx(s)/ds; $x'' = d^2x(s)/ds^2$

x''(s) + k(s) x(s) = 0 known as Hill's equation, derived in 1801 to describe planetary motion

Generalised oscillator equation with position dependent, periodic restoring force k(L+s) = k(s) given by the quadrupole

gradients (+ the small weakly focusing bending term in the ring plane)

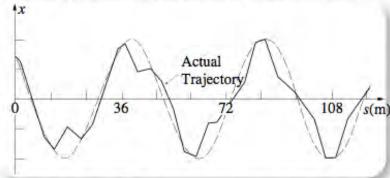
Solution:
$$x(s) = \sqrt{\epsilon \beta(s)} \cos(\mu(s) + \phi)$$

Phase advance

$$\mu(\mathbf{s}) = \int_0^{\mathbf{s}} \frac{d\mathbf{s}}{\beta(\mathbf{s})}$$

Tune # of betatron oscillations

$$Q = \mu / 2\pi$$



motion $x/\sqrt{\beta}$ plotted with phase advance normalised coordinates - becomes simple cos

 $\beta(s)$ beta function, describes the focusing properties of the magnetic lattice invariant, together with $\beta(s)$ amplitude. "single particle emittance"

Motion conveniently described in phase space (x, x') with local slopes or angles $x' = p_x / p$ and linear optics elements as matrices; with simple case for M, applies for IP to IP

$$\begin{pmatrix} x(s) \\ x'(s) \end{pmatrix} = \mathbf{M} \begin{pmatrix} x(s_0) \\ x'(s_0) \end{pmatrix} \qquad \mathbf{M} = \begin{pmatrix} \cos 2\pi Q & \beta \sin 2\pi Q \\ -\frac{1}{\beta} \sin 2\pi Q & \cos 2\pi Q \end{pmatrix}$$

Transverse beam size and emittance

consider : beam of many particles on stable orbit and

simple case: dispersion and slope $\beta' = 0$ by default at IP - relevant for experiments

beam size, r.m.s. $\sigma(s) = \sqrt{\varepsilon \beta(s)}$ beam divergence, r.m.s. $\theta(s) = \sqrt{\varepsilon/\beta(s)}$ product $\varepsilon = \sigma(s)\theta(s)$

 β - function : local machine quantity - focusing of lattice

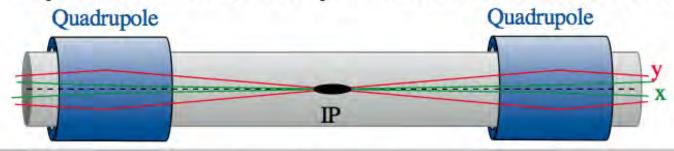
Emittance ε : beam quantity - the average action

related to phase space density or kind of beam temperature given by initial conditions (injected beam)

or equilibrium of quantum excitation and damping - 2nd lecture

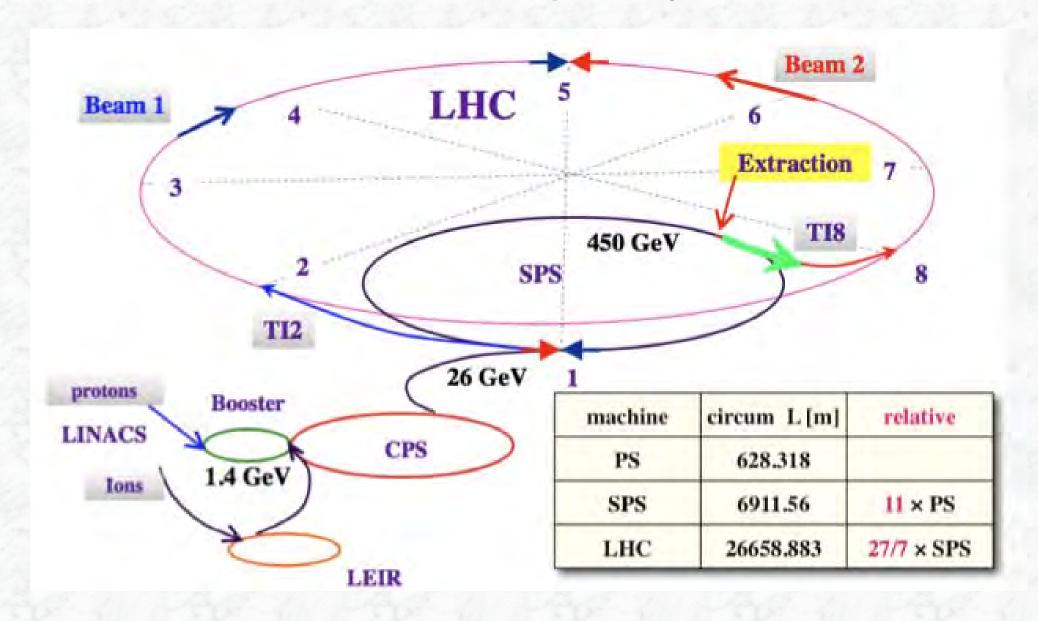
in ideal machine: x, y, z motion uncoupled, 3 emittances ε_x , ε_y , ε_z

IP: squeeze β to a minimum, called $\beta^* \Rightarrow$ maximum of divergence, needs aperture



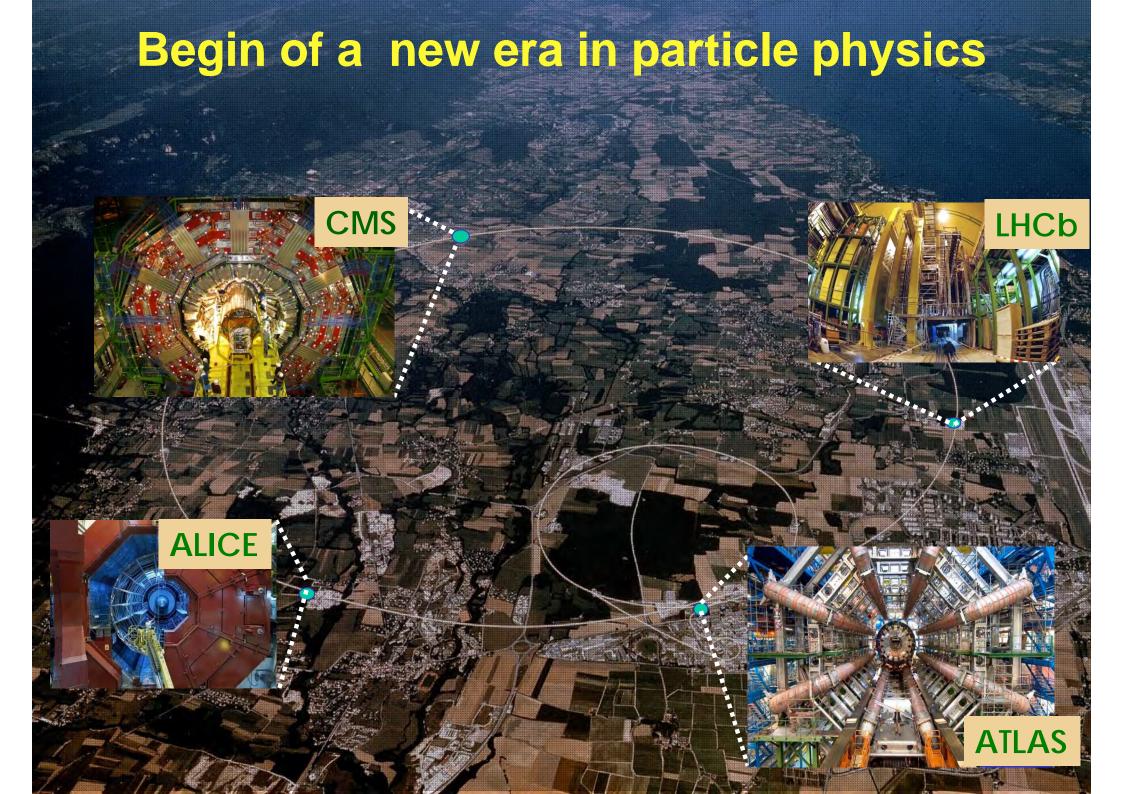
LHC $\epsilon_N = \epsilon \ \beta \gamma = 3.75 \ \mu m$, at top $E_b = 7 \ TeV$: $\epsilon = 0.503 \ nm$, $\beta^* = 0.55 \ m$, $\sigma^* = 16.63 \ \mu m$, $\theta^* = 30 \ \mu rad$

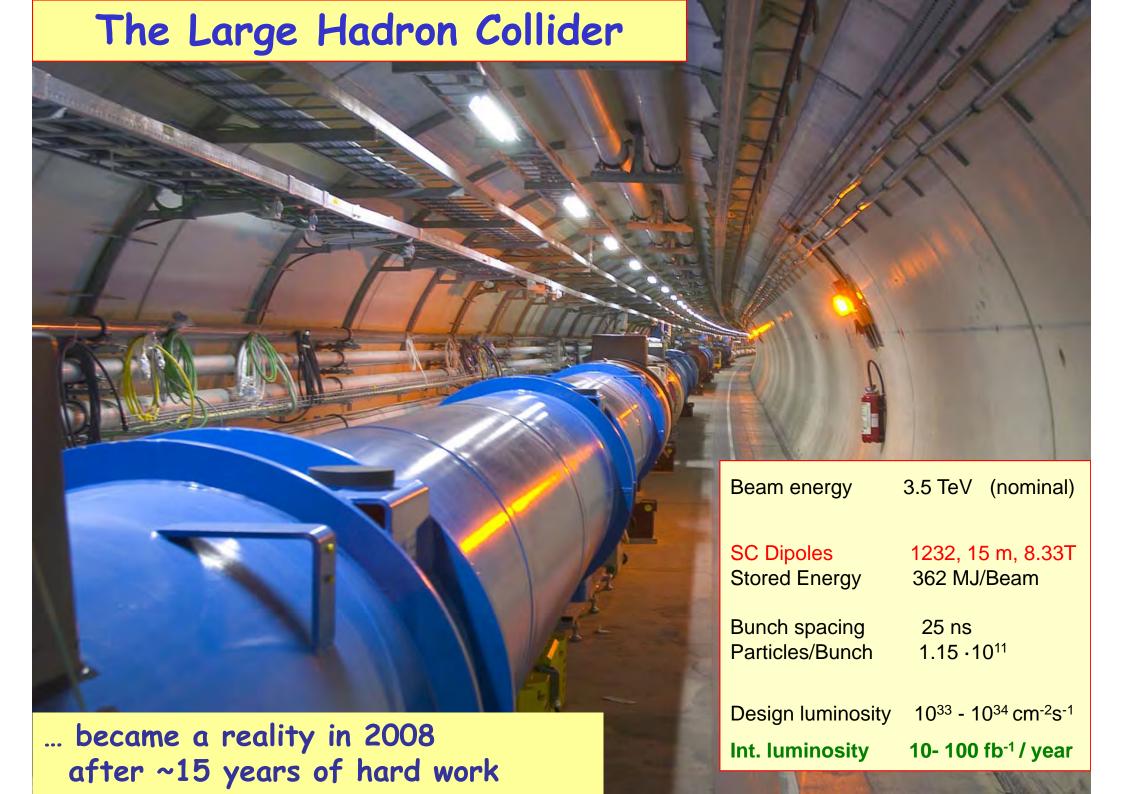
The CERN accelerator complex: injectors and transfer



1.3 The Large Hadron Collider (LHC)







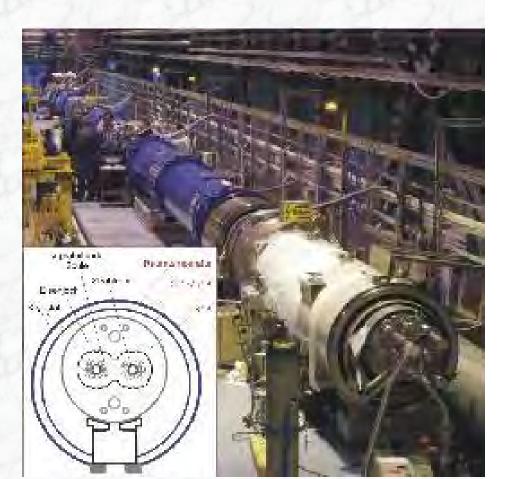
Important components of the accelerator

- Superconducting dipole magnets (the largest challenge)
 - Magnetic field of 8.33 Tesla
 - in total 1232, 15 m long
 - Operation temperature of 1.9 K (helium cooling)
 - Eight acceleration structures,
 Field gradient of 5 MV/m

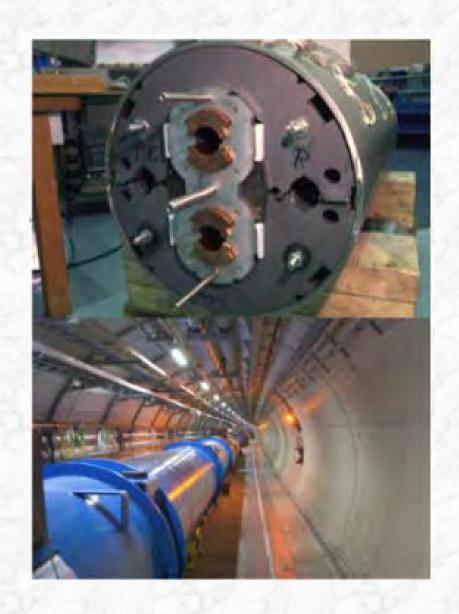
Unprecedented complexity

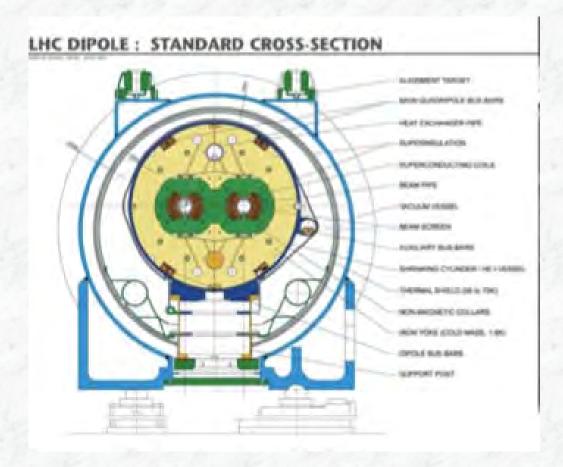
(in total: ~10.000 magnets powered by 1.700 electrical circuits, large stored energy, complex protection systems)





LHC dipole magnets

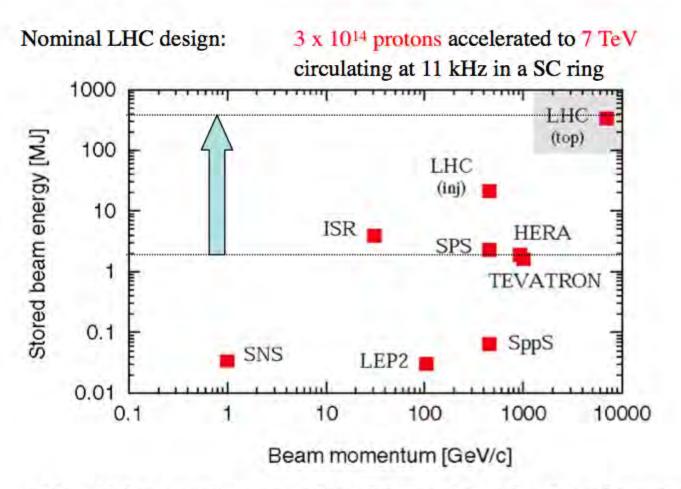




Two beams in one dipole magnet, 8.33 Tesla (opposite magnetic dipole fields, protons circulating in opposite directions)

15 m long, mass of 30 tons

The total stored energy of the LHC beams



LHC: > 100 x higher stored energy and small beam size: ~ 3 orders of magnitude in energy density and damage potential. Active protection (beam loss monitors, interlocks) and collimation for machine and experiments essential.

Only the specially designed beam dump can safely absorb this energy.

Beam parameters, LHC compared to LEP

	LHC	LEP2
Momentum at collision, TeV/c	7	0.1
Nominal design Luminosity, cm ⁻² s ⁻¹	1.00E+34	1.00E+32
Dipole field at top energy, T	1	1
Number of bunches, each beam	2808	4
Particles / bunch	1.15E+11	4.20E+11
Typical beam size in ring, μm	200-300	1800/140 (H/V)
Beam size at IP, µm	16	200/3 (H/V)

• Energy stored in the magnet system:

10 GJoule

Kin. energy of Airbus A380, 560t at 700 km/h.

Energy stored in one (of 8) dipole circuit: 1.1 GJ

Energy stored in one beam:

362 MJ

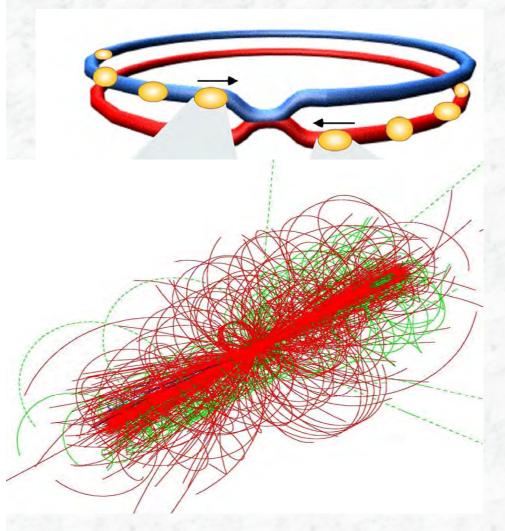
Energy to heat and melt one kg of copper: 0.7 MJ

the LEP2 total stored beam energy was about 0.03 MJ

LHC: From first ideas to realisation



Proton-proton collisions at the LHC



Proton-proton:

2808 x 2808 bunches Separation: 7.5 m (25 ns)

 10^{11} protons / bunch crossing rate of p bunches: 40 Mio / s Luminosity: $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

~10⁹ pp collisions / s (superposition of > 20 pp interactions per crossing: pile-up)

- ~1600 charged particles in the detector
- ⇒ high particle densities, high requirements on detectors

An excellent LHC start: first beams - Sept 10, 2008



Incident on 19th Sep. 2008, repair, comeback.....

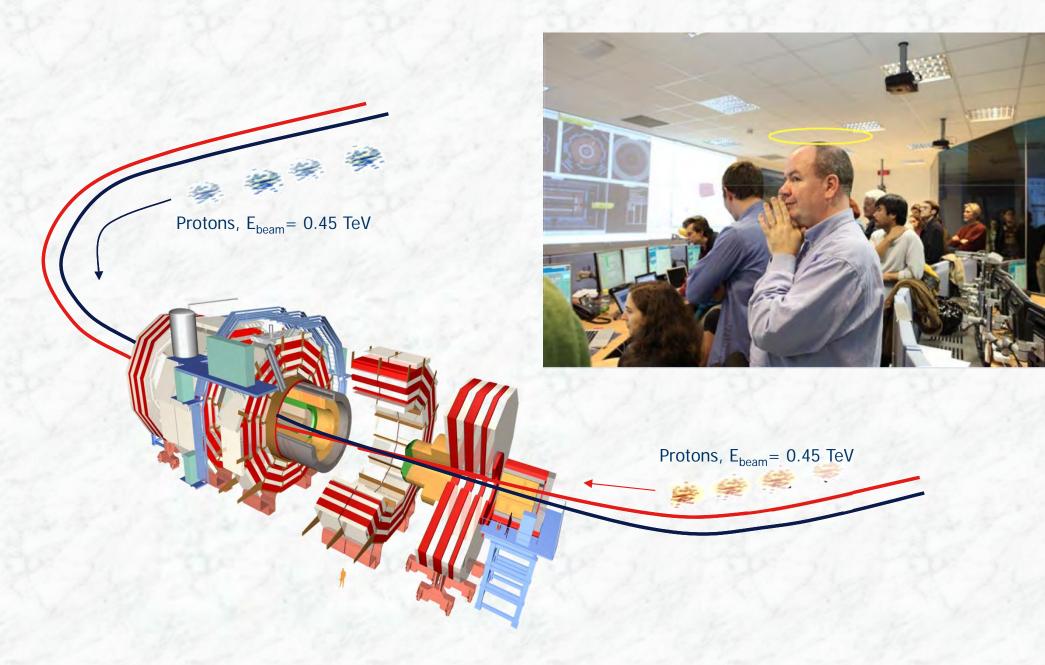
- A resistive zone developed in an electrical bus bar connection
- Electrical arc → punctured the helium enclosure
- Helium release under high pressure
- Relief discs unable to maintain the pressure rise below 0.15 MPa
 - → large pressure forces



- Lot of repair work during 2009
 - (14 quadrupole and 39 dipole magnets replaced, electrical interconnections repaired, larger helium pressure release ports installed,....)
- A very successful re-start in Nov. 2009



LHC re-start in Nov. 2009



The first signals in the ATLAS experiment, 20. Nov 2009

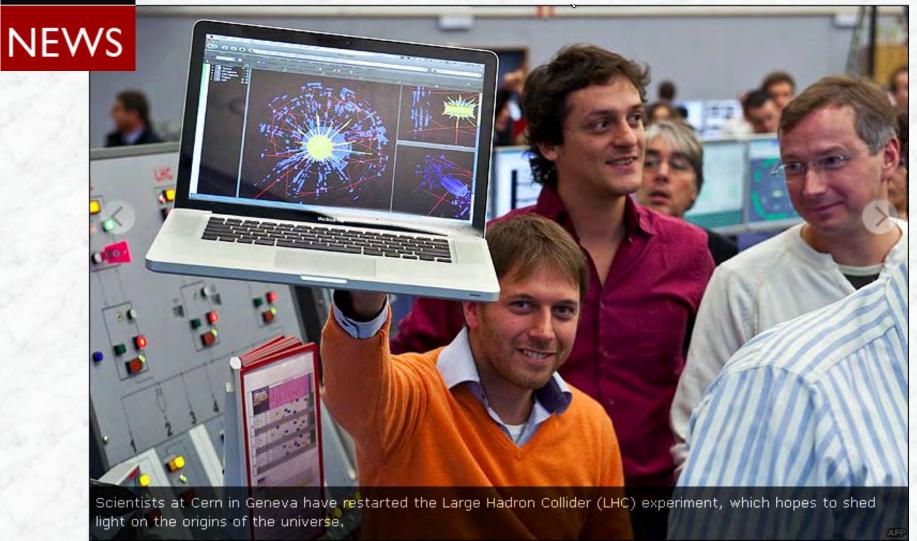




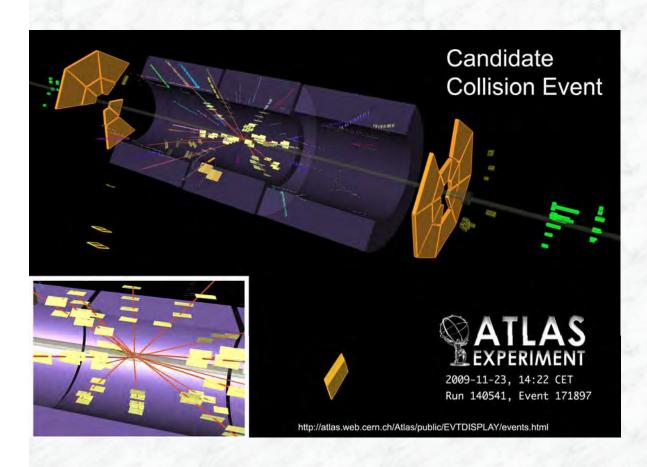
CMS in the BBC news

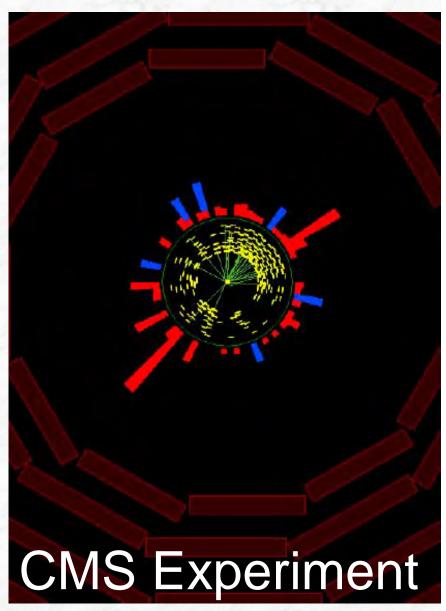
November 21, 2009





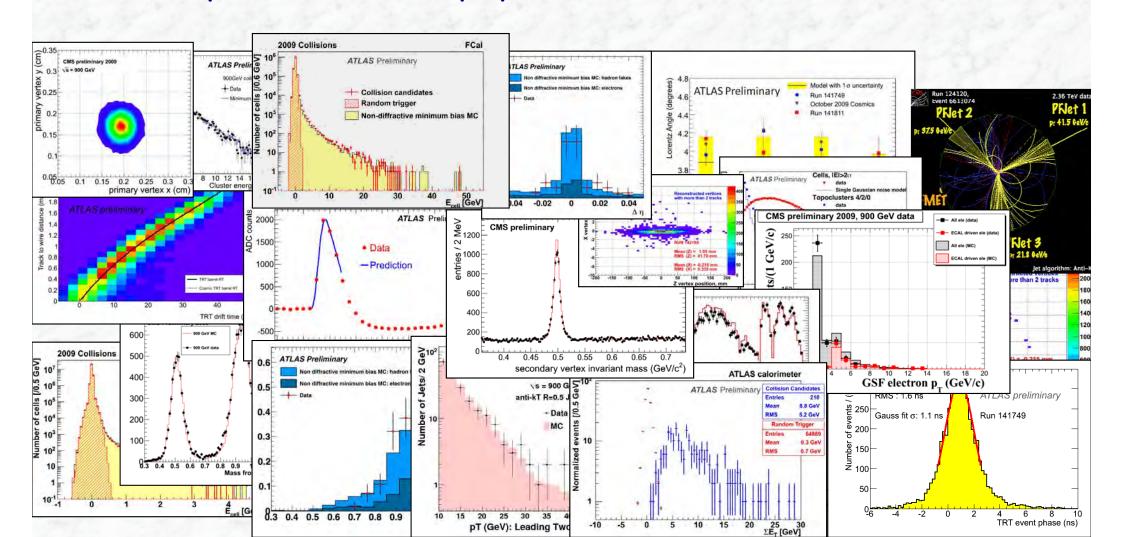
23. Nov 2009: First collisions at 900 GeV





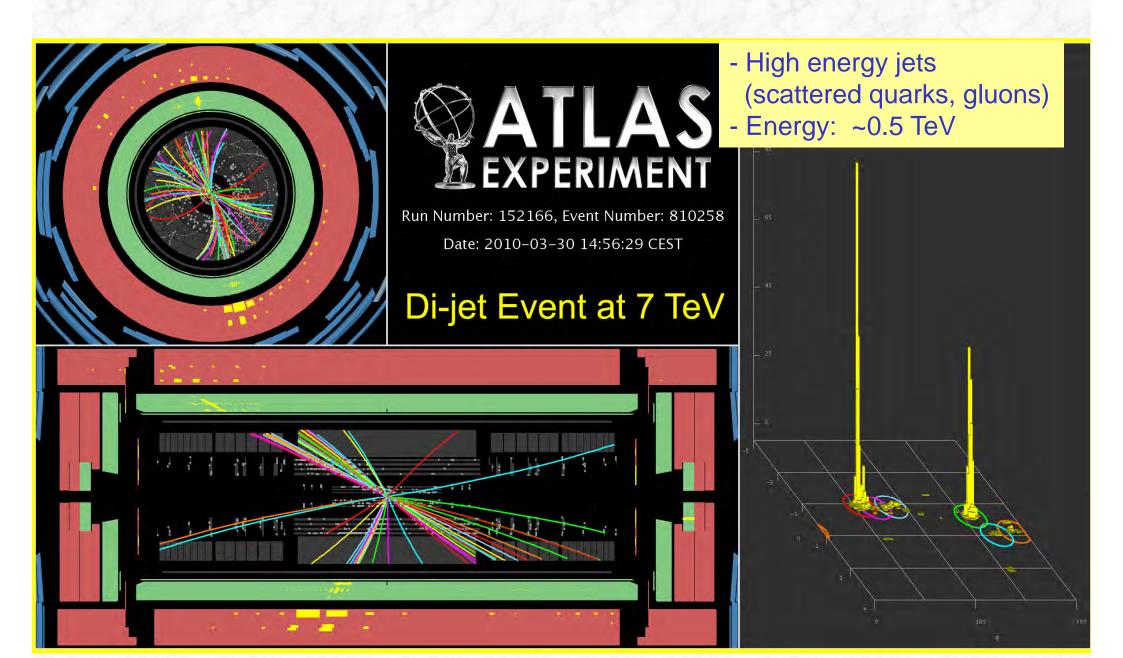
First results on detector performance after only a few days / weeks

First publications of physics results in Feb/March 2010

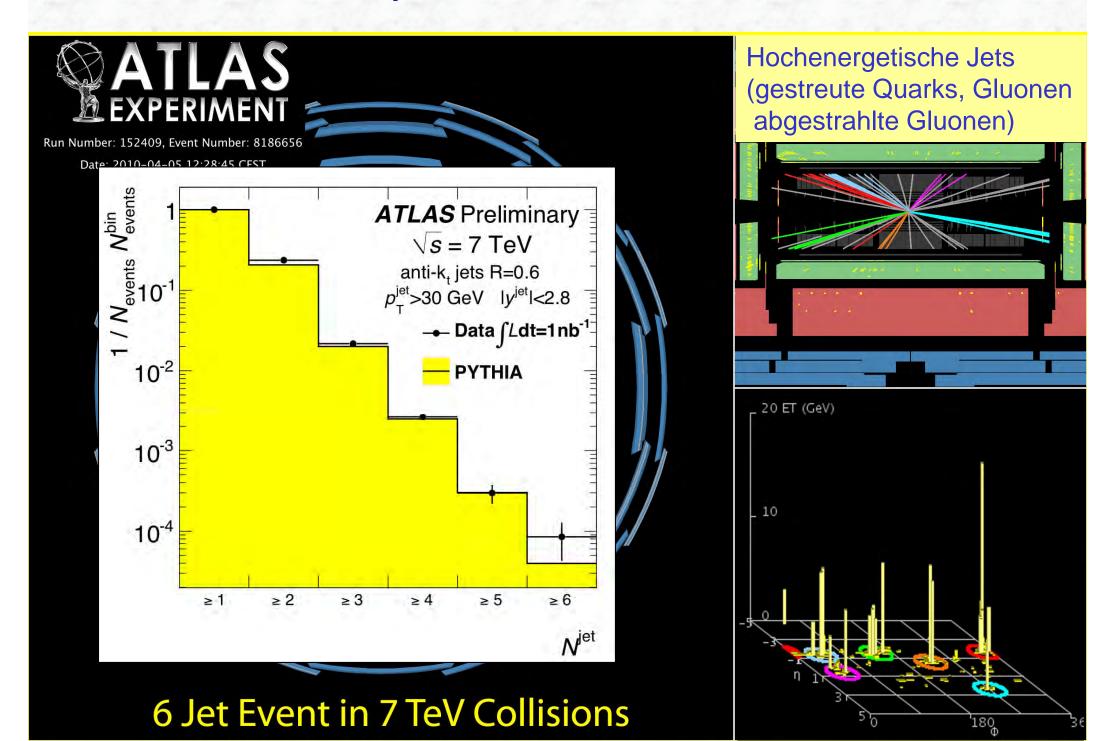


Since 30. March 2010: collisions at 7 TeV

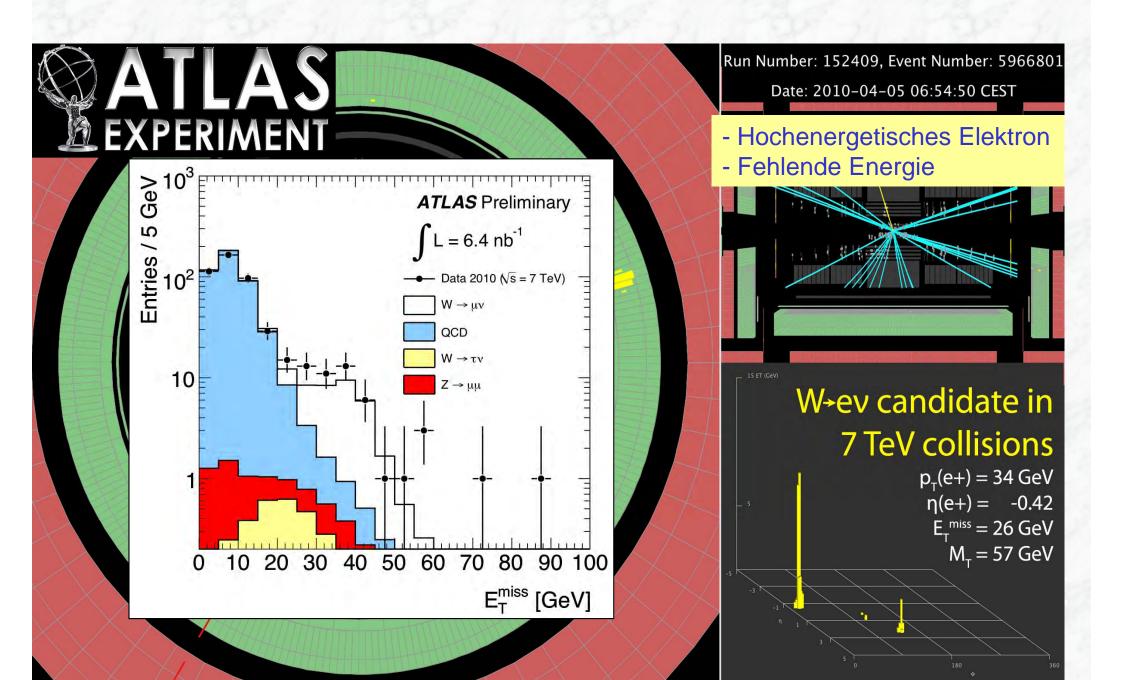
(.... first interesting results appeared soon)



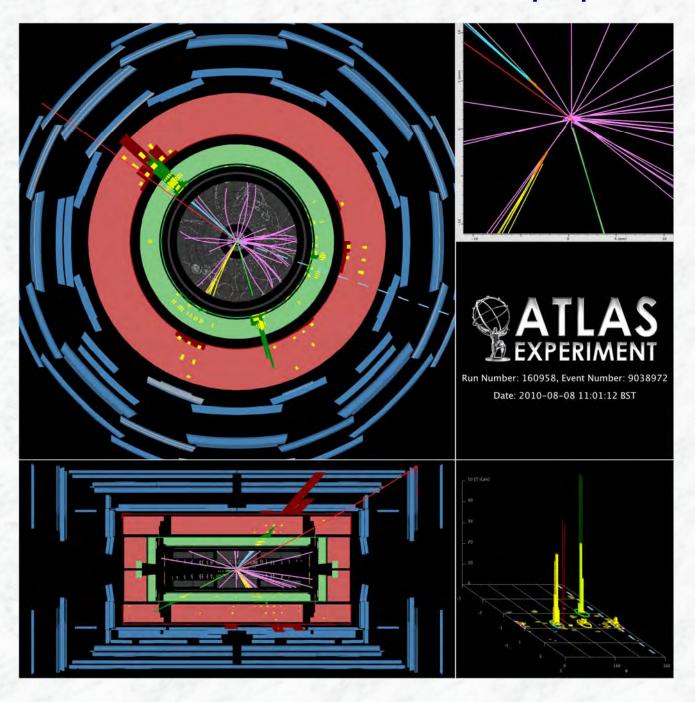
A six-jet event at 7 TeV

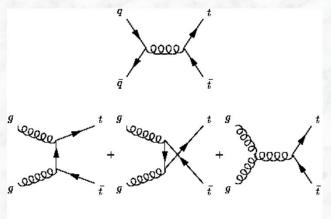


Production of W and Z bosons



Production of the first top quarks in Europe





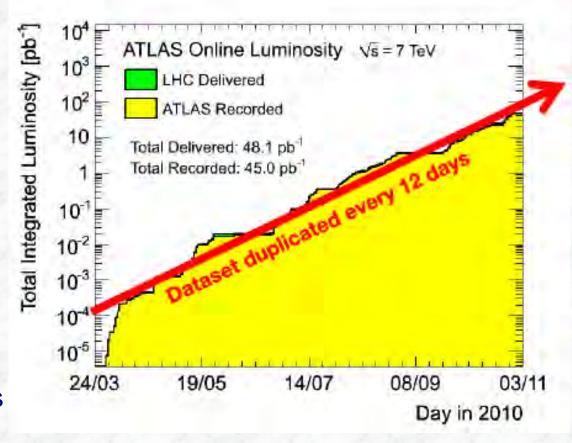
 $tt \rightarrow Wb Wb \rightarrow evb \mu vb$

The fragmentation products of b-quarks (B-Hadrons) have a life time of 1.5 ps

= decay distance of ~2.5 mm

Collected data in 2010:

~40 pb⁻¹ recorded ~36 pb⁻¹ used in analysis (good quality)

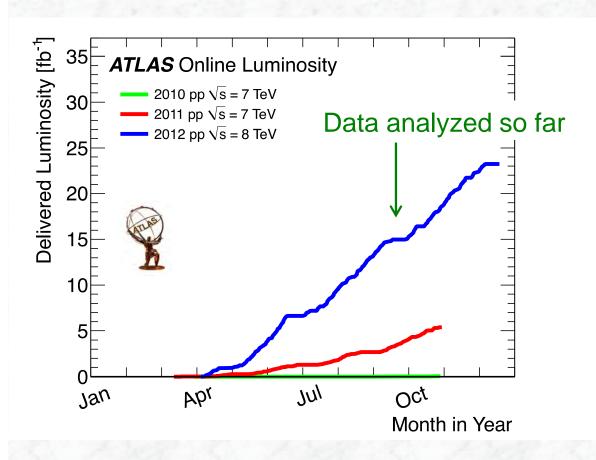


Running again since a few weeks

collected so far: $> 100 \text{ pb}^{-1}$ expected for 2011: $1 - 2 \text{ fb}^{-1}$

World record in instantaneous luminosity on 22. April 2011: 4.67 10³² cm⁻² s⁻¹ so far: Tevatron record: 4.02 10³² cm⁻² s⁻¹

Data taking in 2011/2012



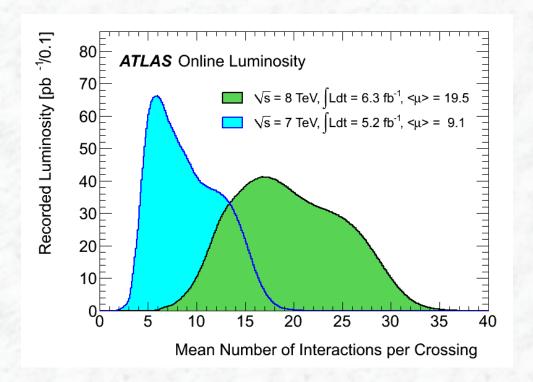
Until end 2012:

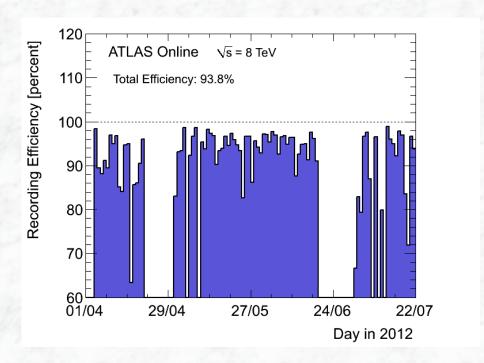
 $> 10^{15}$ pp collisions

~10¹⁰ pp collisions recorded

25 ⋅10⁶ Z \rightarrow µµ decays produced

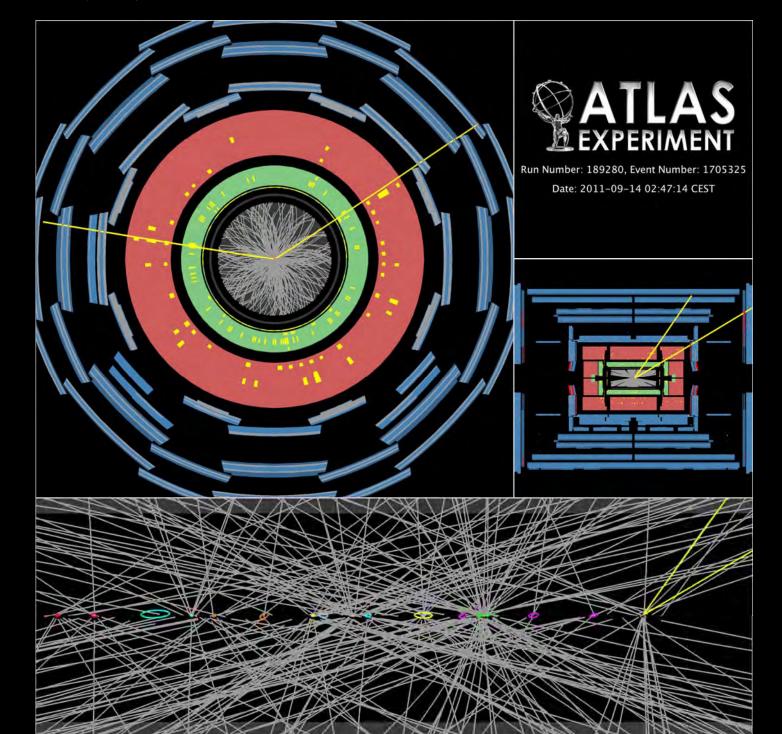
- Excellent LHC performance in 2011and 2012 (far beyond expectations)
- Peak luminosity seen by ATLAS:
 7.7 10³³ cm⁻² s⁻¹ (world record, 2012)
- Excellent performance of the experiments in recording the data (efficiency ~93.5%, working detector channels >99%, speed of data analysis,...)



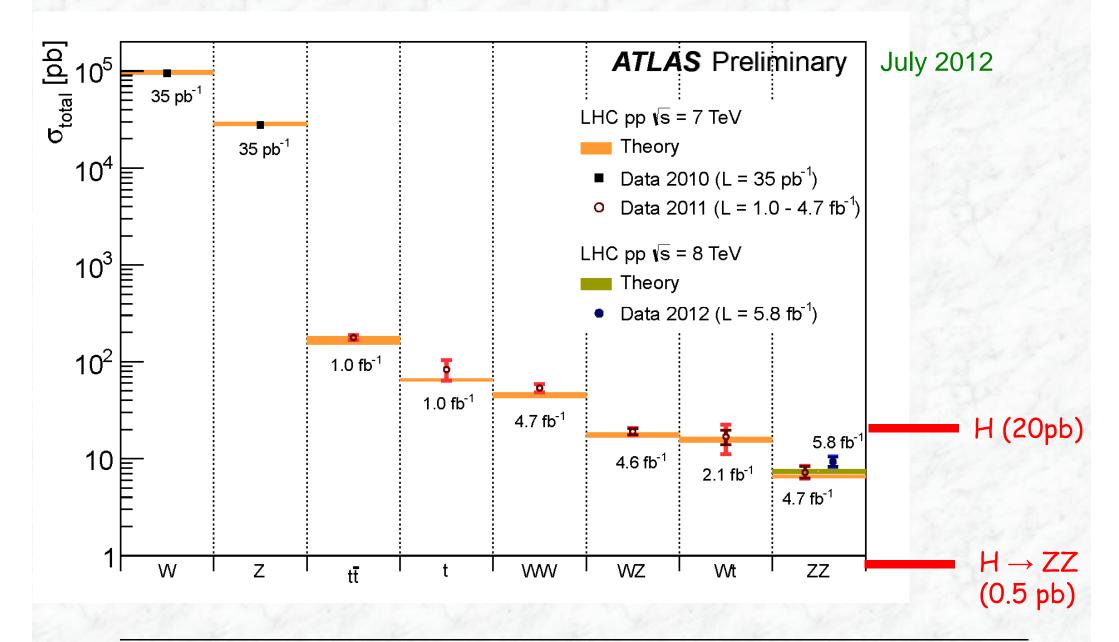


 High efficiencies of the ATLAS and CMS experiments to collect and analyze the data

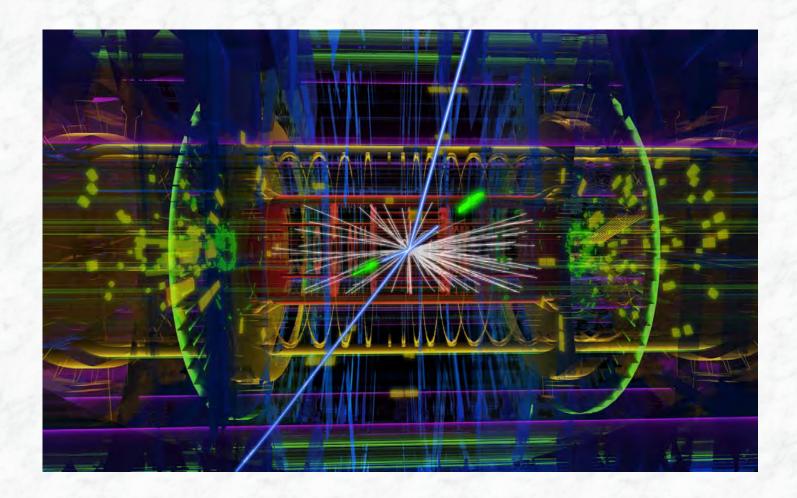
$Z \rightarrow \mu^+ \mu^-$ with 20 superimposed events



The Standard Model at the LHC



Discovery of a Higgs-like particle





K. Jakobs



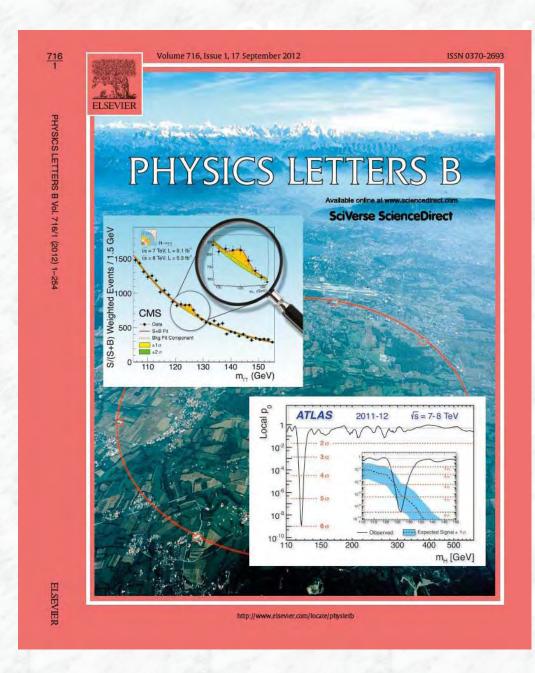
4th July 2012





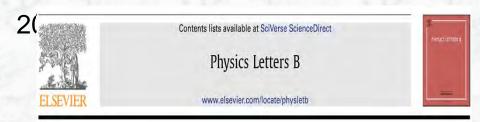
From the editorial:

The top Breakthrough of the Year — the discovery of the Higgs boson — was an unusually easy choice, representing both a triumph of the human intellect and the culmination of decades of work by many thousands of physicists and engineers



a New Particle

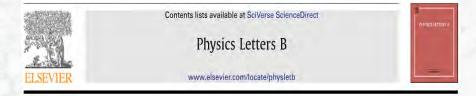
Submission to PLB on 31st July



Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC $^{\diamondsuit}$

ATLAS Collaboration *

This paper is dedicated to the memory of our ATLAS colleagues who did not live to see the full impact and significance of their contributions to the experiment.



CMS Collaboration*

CERN, Switzerland

This paper is dedicated to the memory of our colleagues who worked on CMS but have since passed away. In recognition of their many contributions to the achievement of this observation.

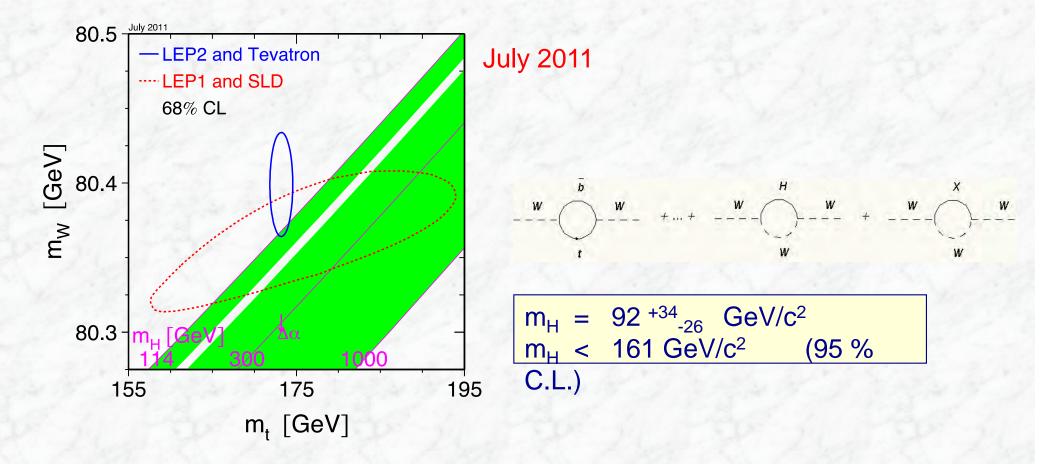
Decay observed into particles with same spin and electric charge sum = 0

→ a new neutral boson has been discovered

Constraints on the Higgs boson mass (before LHC)

• $m_H > 114.4 \text{ GeV/c}^2$

- from direct searches at LEP
- $m_H < 156 \text{ GeV/c}^2$.or. $m_H > 177 \text{ GeV/c}^2$ from direct searches at the Tevatron

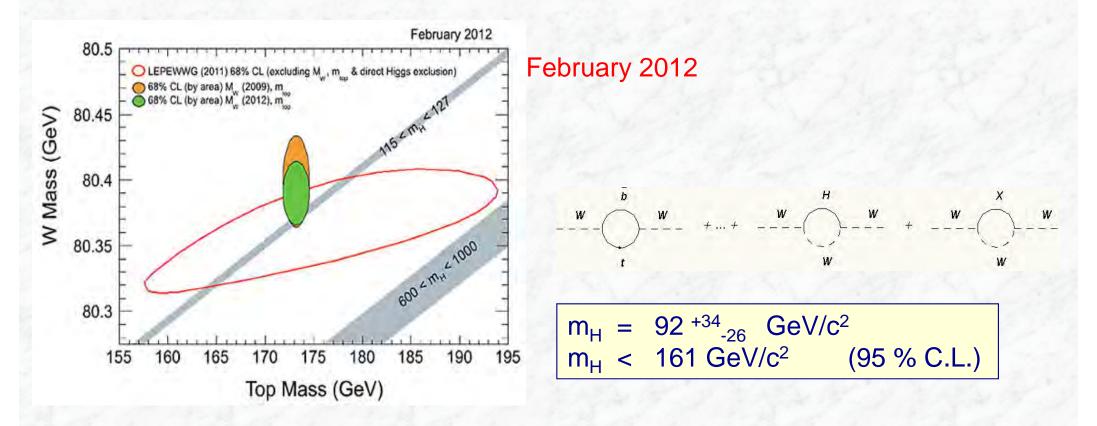


Indirect constraints from precision measurements (quantum corrections)

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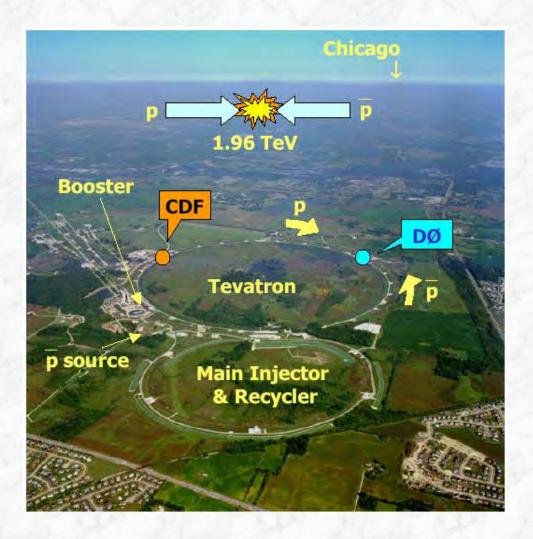
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Indirect constraints from precision measurements (quantum corrections)

1.4 The Fermilab Tevatron collider

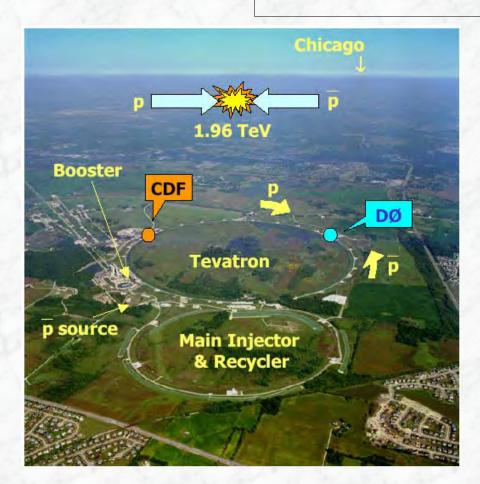




The Tevatron Collider at Fermilab

0

- Proton antiproton collider
 - 6.5 km circumference
 - Beam energy 0.98 TeV, $\sqrt{s} = 1.96 \text{ TeV}$
 - 36 bunches, 396 ns separation (time between crossings)
- 2 Experiments: CDF and DØ
- Main challenges:
 - Antiproton production and storage
 - → luminosity, stability of operation



Collider is running in so called Run II (since 2001)

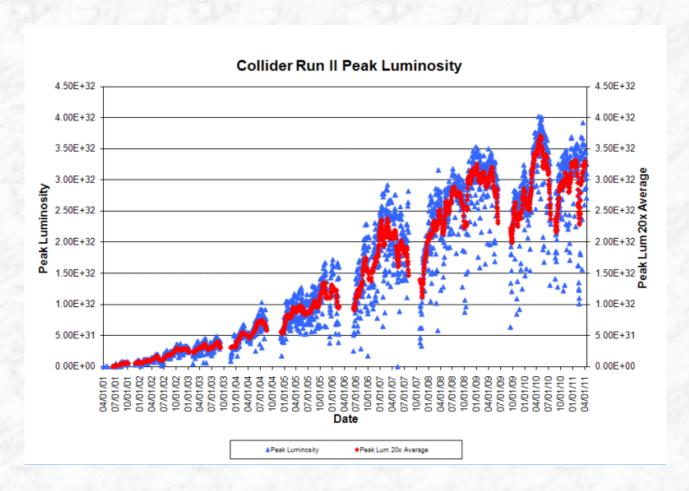
[Run I from 1990 – 1996, int. luminosity: 0.125 fb⁻¹, Top quark discovery]

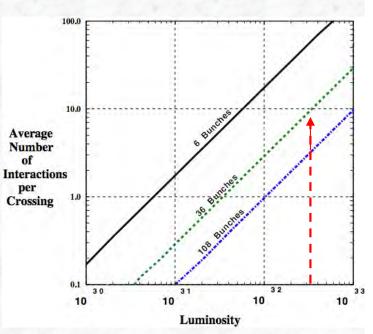
* March 2001 – Feb 2006: Run II a, $\int L dt = 1.2 \text{ fb}^{-1}$

* July 2006 - 2011: Run II b, $\int L dt = 10 - 12 \text{ fb}^{-1}$

Tevatron performance

Peak luminosities of the machine as a function of time

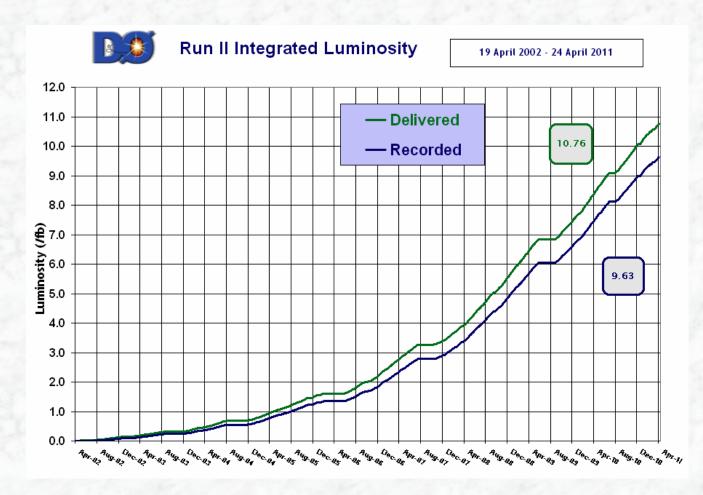




- Peak luminosity of 4 · 10³² cm⁻² s⁻¹
- Corresponds to ~10 interactions per bunch crossing (superposition of minimum bias events on hard collision)

The integrated Tevatron luminosity (until April 2011)

- After a slow start-up (2001 2003), the Tevatron accelerator has reached an excellent performance
- Today, Tevatron delivers a data set equal to Run I (~100 pb⁻¹) every 2 weeks
- Integrated luminosity delivered to the experiments so far ~ 10.8 fb⁻¹
- Anticipate an int. luminosity of ~12 fb⁻¹ until end of 2011.

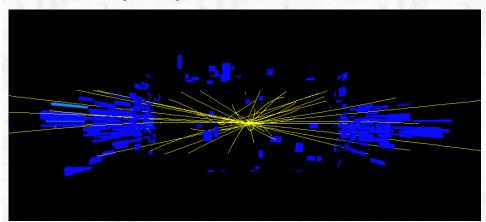


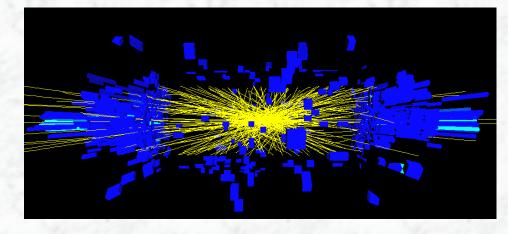
Data corresponding to an int. luminosity of up to ~8 fb⁻¹ analyzed...

Challenges with high luminosity

Min. bias pileup at the Tevatron, at $0.6 \cdot 10^{32} \text{ cm}^2\text{s}^{-1}$

... and at 2.4 · 10³² cm²s⁻¹

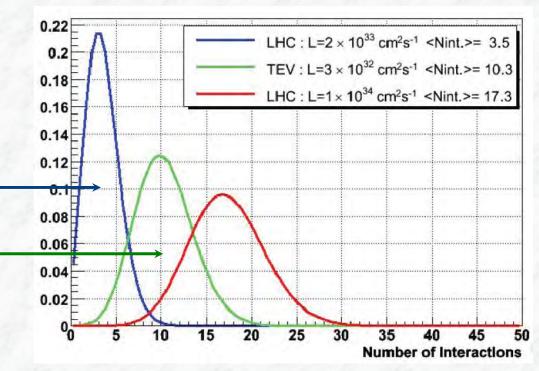




Average number of interactions:

LHC: initial "low" luminosity run $(L=2.10^{33} \text{ cm}^2\text{s}^{-1})$: < N>=3.5

TeV: $(L=3 \cdot 10^{32} \text{ cm}^2\text{s}^{-1})$: < N > = 10



Comparison of the LHC and Tevatron machine parameters

	LHC (design)	Tevatron (achieved)
Centre-of-mass energy	14 TeV	1.96 TeV
Number of bunches	2808	36
Bunch spacing	25 ns	396 ns
Energy stored in beam	360 MJ	1 MJ
Peak Luminosity	10 ³³ -10 ³⁴ cm ⁻² s ⁻¹	3.5 x 10 ³² cm ⁻² s ⁻¹
Integrated Luminosity / year	10-100 fb ⁻¹	~ 2 fb ⁻¹

- 7 times more energy (after initial 3.5 TeV phase)
- Factor 3-30 times more luminosity
- Physics cross sections factor 10-100 larger