

6. Physics of W and Z bosons

6.1 The theory of electroweak interactions (a short repetition)

6.2 Summary of precision tests at LEP

6.3 W and Z boson production in hadron colliders

6.4 Test of QCD in W/Z (+jet) production

6.5 W mass measurement

Weak Isospin and Hypercharge Quantum

Lepton	T	T^3	Q	Y
ν_e	$\frac{1}{2}$	$\frac{1}{2}$	0	-1
e_L^-	$\frac{1}{2}$	$-\frac{1}{2}$	-1	-1
e_R^-	0	0	-1	-2

Numbers of Leptons and Quarks

Quark	T	T^3	Q	Y
u_L	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{1}{3}$
d_L	$\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{3}$
u_R	0	0	$\frac{2}{3}$	$\frac{4}{3}$
d_R	0	0	$-\frac{1}{3}$	$-\frac{2}{3}$

Important Milestones towards Electroweak Unification

1961	S. Glashow proposes an electroweak gauge theory, Introduction of massive W^\pm and Z^0 bosons, to explain the large difference in strength of electromagnetic and weak interactions. Key question: how acquire W and Z bosons mass?
1964	R. Brout, F. Englert and P. Higgs demonstrate that mass terms for gauge bosons can be introduced in local gauge invariant theories via spontaneous symmetry breaking
1967	S. Weinberg and A. Salam use Brout-Englert-Higgs mechanism to introduce mass terms for W and Z bosons in Glashow's theory → GSW theory (Glashow, Salam, Weinberg) → mass terms for W, Z bosons, γ remains massless → Higgs particle (see chapter 7)
1973	G. t'Hooft and M. Veltman show that GSW theory is renormalizable
1979	Nobel price for S. Glashow, A. Salam and S. Weinberg
1983	Experimental discovery of the W and Z bosons by UA1 and UA2 experiments at the CERN ppbar collider ($\sqrt{s} = 540$ GeV)
1990-2000	Precise test of the electroweak theory at LEP
1999	Nobel price for G. t'Hooft and M. Veltman
2012	Discovery of a Higgs particle by the ATLAS and CMS experiments at the LHC

W and Z vertex factors

$$\left. \begin{aligned} & -i\frac{g}{\sqrt{2}}(\bar{\chi}_L\gamma^\mu\tau_+\chi_L)W_\mu^+ \\ & = -i\frac{g}{\sqrt{2}}(\bar{\nu}_L\gamma^\mu e_L)W_\mu^+ \end{aligned} \right\} \begin{array}{c} e^+ \\ \leftarrow W^+ \text{---} \\ \nu \end{array}$$

$$\boxed{-i\frac{g}{\sqrt{2}}\gamma^{\mu\frac{1}{2}}(1-\gamma^5)}$$

$$\left. \begin{aligned} & -i\frac{g}{\sqrt{2}}(\bar{\chi}_L\gamma^\mu\tau_-\chi_L)W_\mu^- \\ & = -i\frac{g}{\sqrt{2}}(\bar{e}_L\gamma^\mu\nu_L)W_\mu^- \end{aligned} \right\} \begin{array}{c} e^- \\ \leftarrow W^- \text{---} \\ \bar{\nu} \end{array}$$

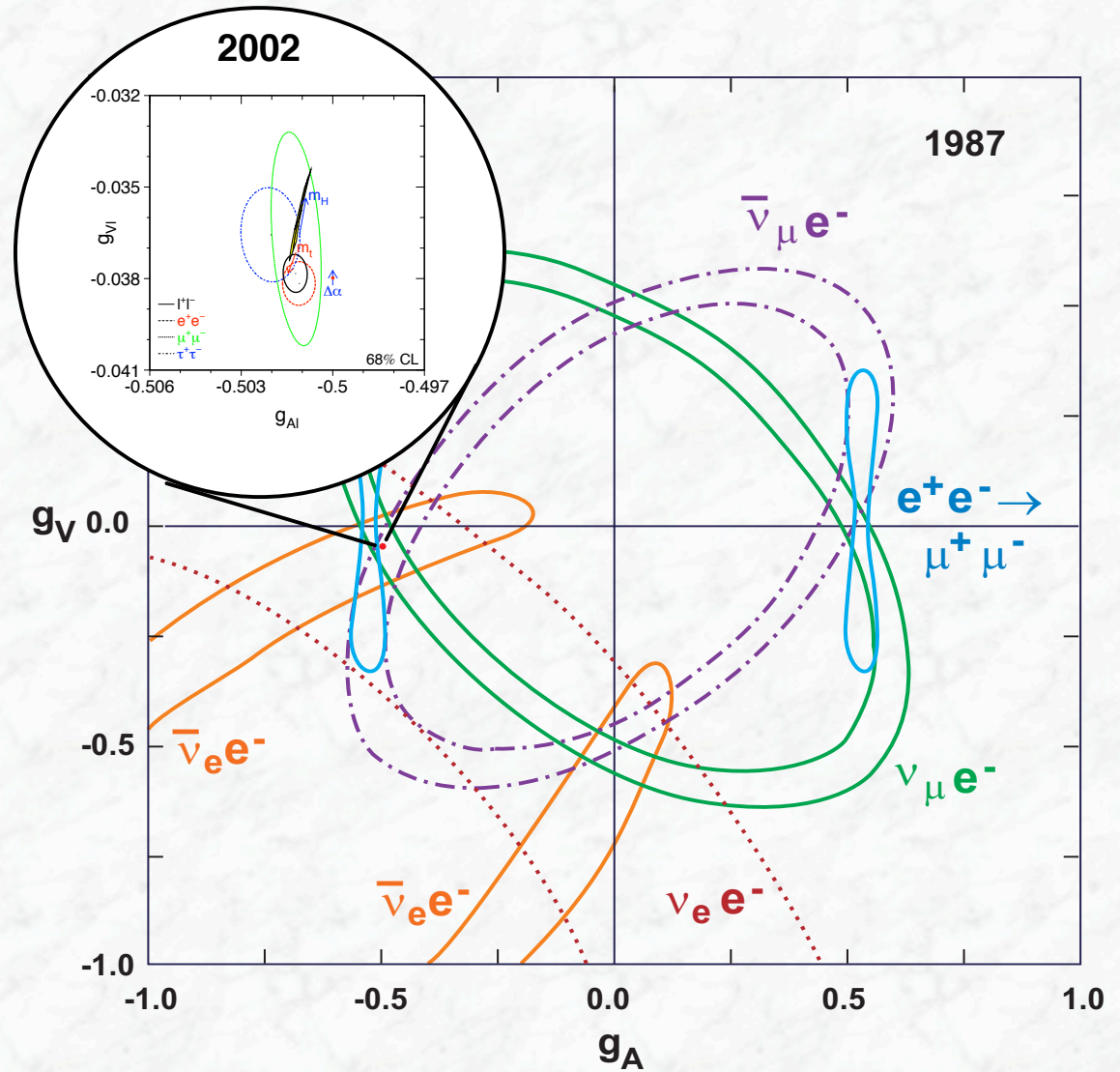
$$\Rightarrow \begin{array}{c} f \\ \leftarrow Z^0 \text{---} \\ \bar{f} \end{array} \quad \boxed{-i\frac{g}{\cos\theta_W}\gamma^{\mu\frac{1}{2}}(c_V^f - c_A^f\gamma^5)}$$

The $Z \rightarrow ff$ vertex factors in the Standard Model
 ($\sin^2 \theta_W$ is assumed to be 0.234)

f	Q_f	c_A^f	c_V^f
ν_e, ν_μ, \dots	0	$\frac{1}{2}$	$\frac{1}{2}$
e, μ, \dots	-1	$-\frac{1}{2}$	$-\frac{1}{2} + 2 \sin^2 \theta_W$ 0.03
u, c, \dots	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{2} - \frac{4}{3} \sin^2 \theta_W$ 0.19
d, s, \dots	$-\frac{1}{3}$	$-\frac{1}{2}$	$-\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W$ 0.34

6.2 Summary of electroweak precision tests at LEP

- Results of 30 years of experimental and theoretical progress
- The electroweak theory is tested at the level of 10^{-4}



LEP am CERN / Genf



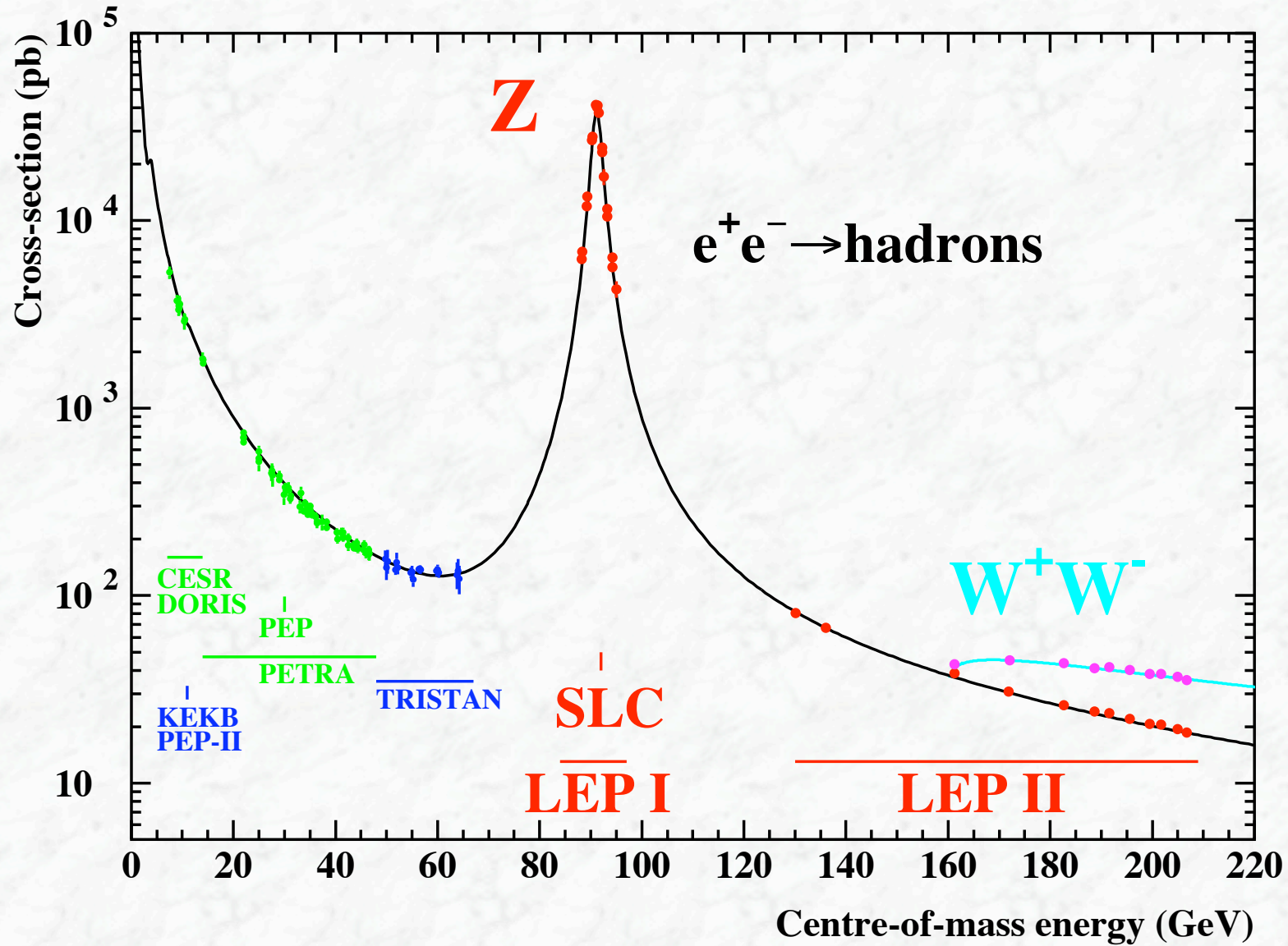
e^+e^- -Beschleuniger, 27 km Umfang

Schwerpunktenergie: LEP-I (1989-1995) 91 GeV

LEP-II (1996-2000) → 208 GeV

4 Experimente: ALEPH, DELPHI, L3, OPAL

Cross sections for W and Z boson production



Precision tests
of the Z sector

Tests of the
W sector

Cross section for $e^+e^- \rightarrow \mu^+\mu^-$ at LEP I

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left[F_\gamma(\cos\theta) + F_{\gamma Z}(\cos\theta) \frac{s(s-M_Z^2)}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} + F_Z(\cos\theta) \frac{s^2}{(s-M_Z^2)^2 + M_Z^2\Gamma_Z^2} \right]$$

γ

γ/Z interference

Z

vanishes at $\sqrt{s} \approx M_Z$

$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2})(1 + \cos^2\theta) + 8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

$\alpha = \alpha(m_Z)$: running el.magnetic coupling [$\alpha(M_Z) = \alpha / (1 - \Delta\alpha)$ mit $\Delta\alpha \approx 0.06$]

$g_V, g_A = c_V, c_A$: effective coupling constants (vector and axial vector)

Cross section for $e^+e^- \rightarrow ff$ on resonance ($\sqrt{s} = m_Z$)

- On resonance, $\sqrt{s} = m_Z$:
 - γ^*/Z interference terms vanishes
 - γ term contributes $\sim 1\%$
 - **Z contribution dominates !**
- Contribution of the γ^*/Z interference term at $s = (M_Z - 3 \text{ GeV})^2$: $\sim 0.2\%$

Total cross section for $e^+e^- \rightarrow \mu^+\mu^-$ (integration over $\cos \theta$)

$$\sigma_{\text{tot}} \approx \sigma_Z = \frac{4\pi}{3s} \frac{\alpha^2}{16 \sin^4 \theta_W \cos^4 \theta_W} \cdot [(g_V^e)^2 + (g_A^e)^2][(g_V^\mu)^2 + (g_A^\mu)^2] \cdot \frac{s^2}{(s - M_Z^2)^2 + (M_Z \Gamma_Z)^2}$$

$$\sigma_Z(\sqrt{s} = M_Z) = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2} \quad \text{Peak cross section}$$

$$\Gamma_f = \frac{\alpha M_Z}{12 \sin^2 \theta_W \cos^2 \theta_W} \cdot [(g_V^f)^2 + (g_A^f)^2]$$

Partial width

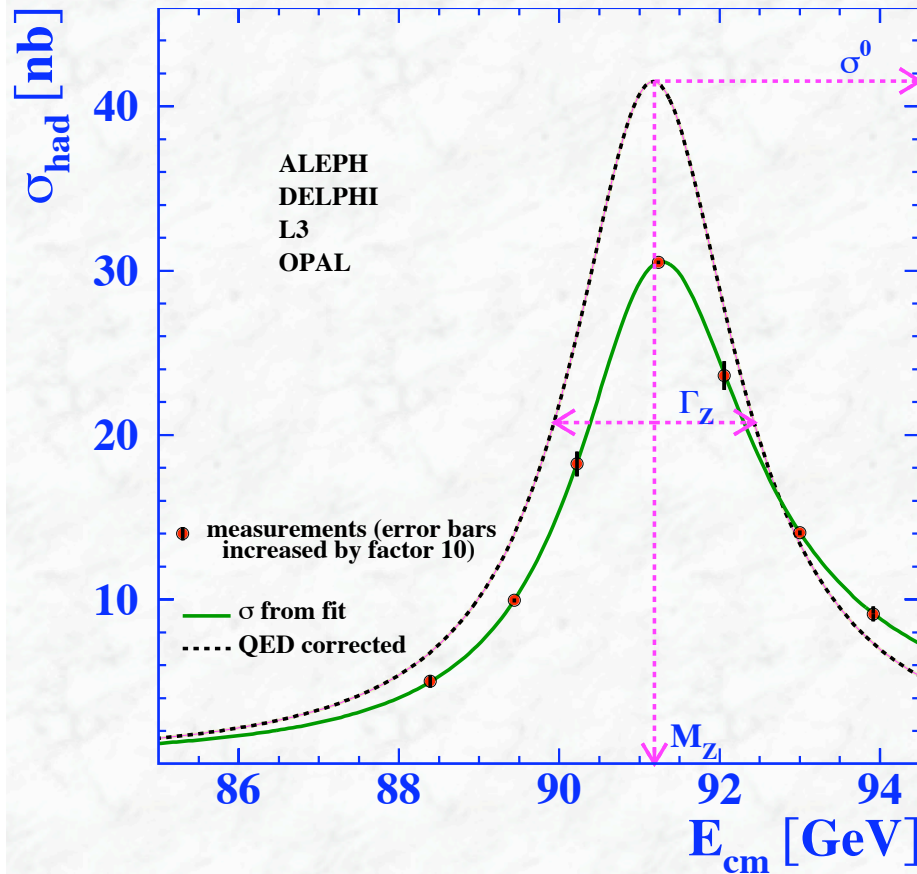
$$\Gamma_Z = \sum_i \Gamma_i \quad \text{Total width}$$

From the energy dependence of the total cross section (for various fermions f) the parameters

M_Z, Γ_Z, Γ_f

can be determined.

Measurement of the Z line-shape



Line shape (resonance curve):

$$\sigma(s) = 12\pi \frac{\Gamma_e \Gamma_\mu}{M_Z^2} \cdot \frac{s}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

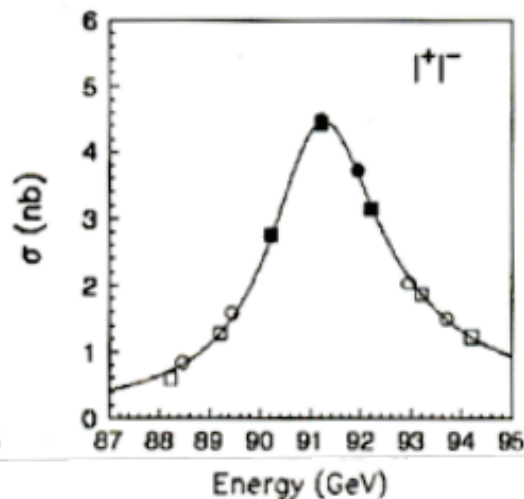
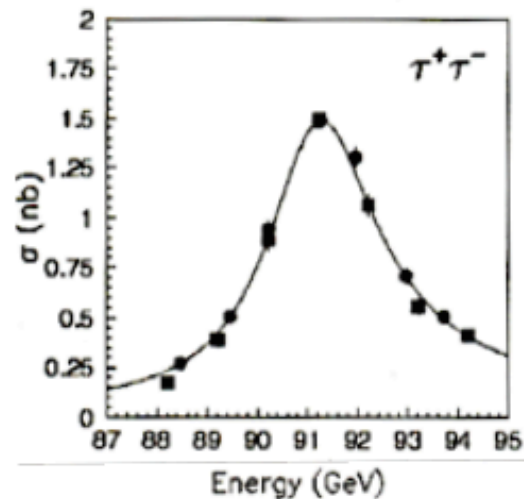
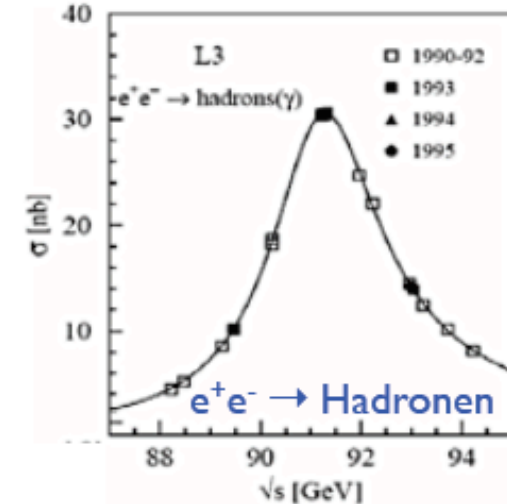
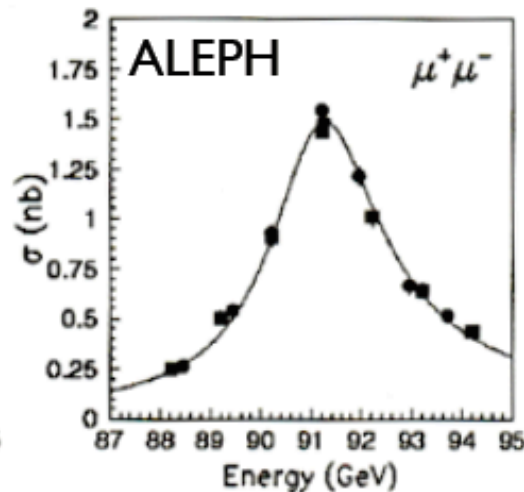
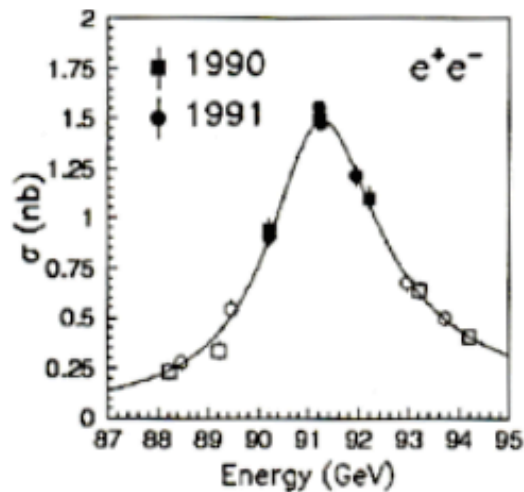
Peak:
$$\sigma_0 = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$$

- Position of maximum $\rightarrow M_Z$
- Full width at half maximum $\rightarrow \Gamma_Z$
- Peak cross section $\sigma_0 \rightarrow \Gamma_e \Gamma_\mu$

Radiative corrections (photon radiation)
important

- with ISR (initial state radiation)
- - - without ISR

Measurement of the Z line-shape (cont.)



Quark-Flavor i.a. nicht exp. trennbar
(Ausnahme: $c, b \rightarrow$ Lebensdauer)
 \Rightarrow had. Breite: $\Gamma_{had} = \Gamma_u + \Gamma_d + \Gamma_s + \Gamma_c + \Gamma_b$

Messe Verhältnisse der Pol-WQ:

$$R_l^0 \equiv \frac{\Gamma_{had}}{\Gamma_{ll}} \quad l = e, \mu, \tau$$

$$R_q^0 \equiv \frac{\Gamma_{qq}}{\Gamma_{had}} \quad q = b, c$$

- Keine Unterschiede für verschiedene Leptonarten \Rightarrow **Leptonuniversalität**
- Form der Resonanzkurve für alle Endzustände gleich (gleicher Propagator!)

Results on Z line-shape parameters

$$M_Z = 91.1876 \pm 0.0021 \text{ GeV} \quad 23 \text{ ppm} (*)$$

$$\begin{aligned} \Gamma_Z &= 2.4952 \pm 0.0023 \text{ GeV} \\ \Gamma_{\text{had}} &= 1.7458 \pm 0.0027 \text{ GeV} \\ \Gamma_e &= 0.08392 \pm 0.00012 \text{ GeV} \\ \Gamma_\mu &= 0.08399 \pm 0.00018 \text{ GeV} \\ \Gamma_\tau &= 0.08408 \pm 0.00022 \text{ GeV} \end{aligned}$$

3 lepton flavours
treated independently



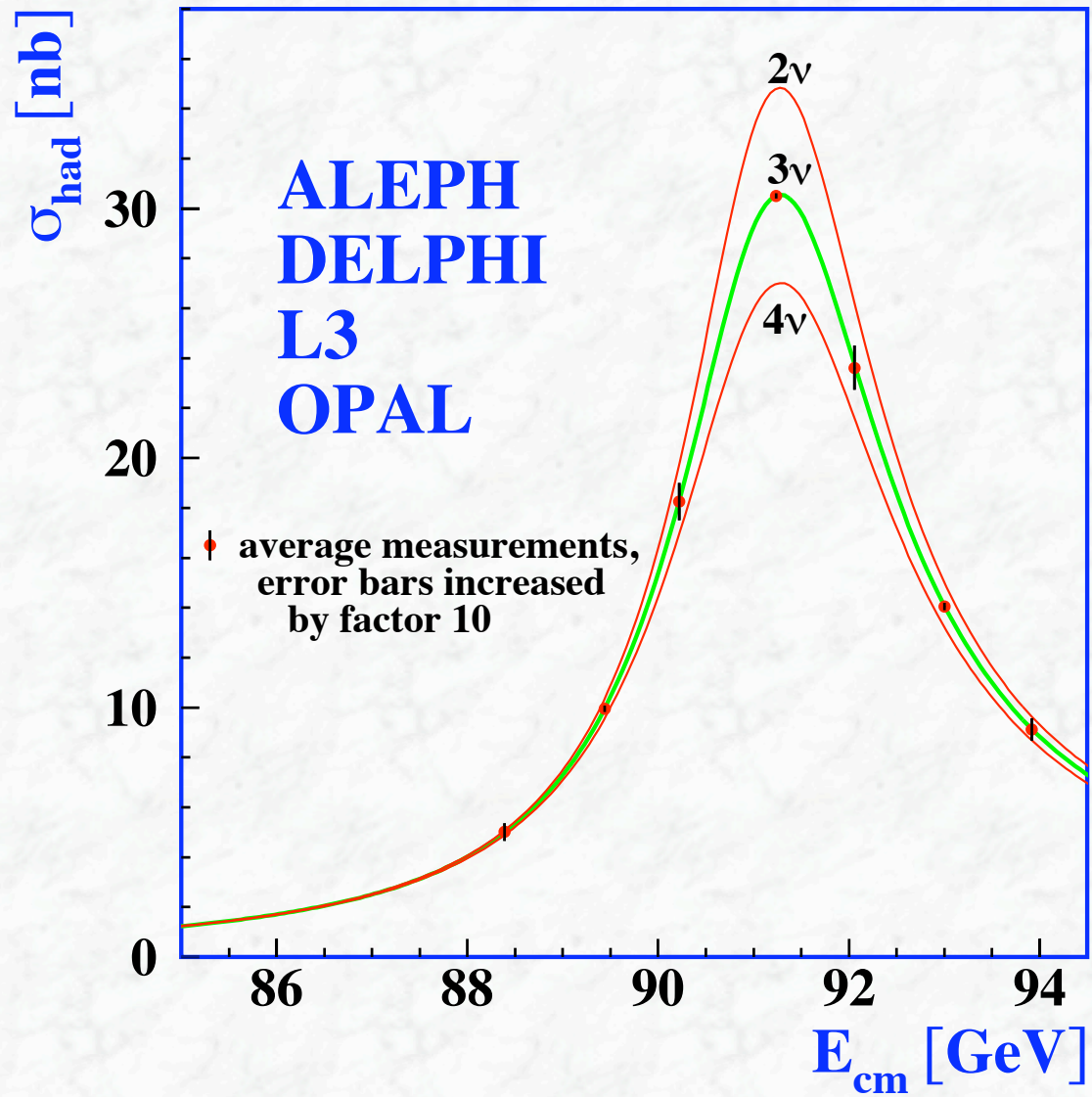
Test of lepton
universality

$$\begin{aligned} \Gamma_Z &= 2.4952 \pm 0.0023 \text{ GeV} \\ \Gamma_{\text{had}} &= 1.7444 \pm 0.0022 \text{ GeV} \\ \Gamma_e &= 0.083985 \pm 0.000086 \text{ GeV} \end{aligned}$$

lepton universality
assumed:
 $\Gamma_e = \Gamma_\mu = \Gamma_\tau$

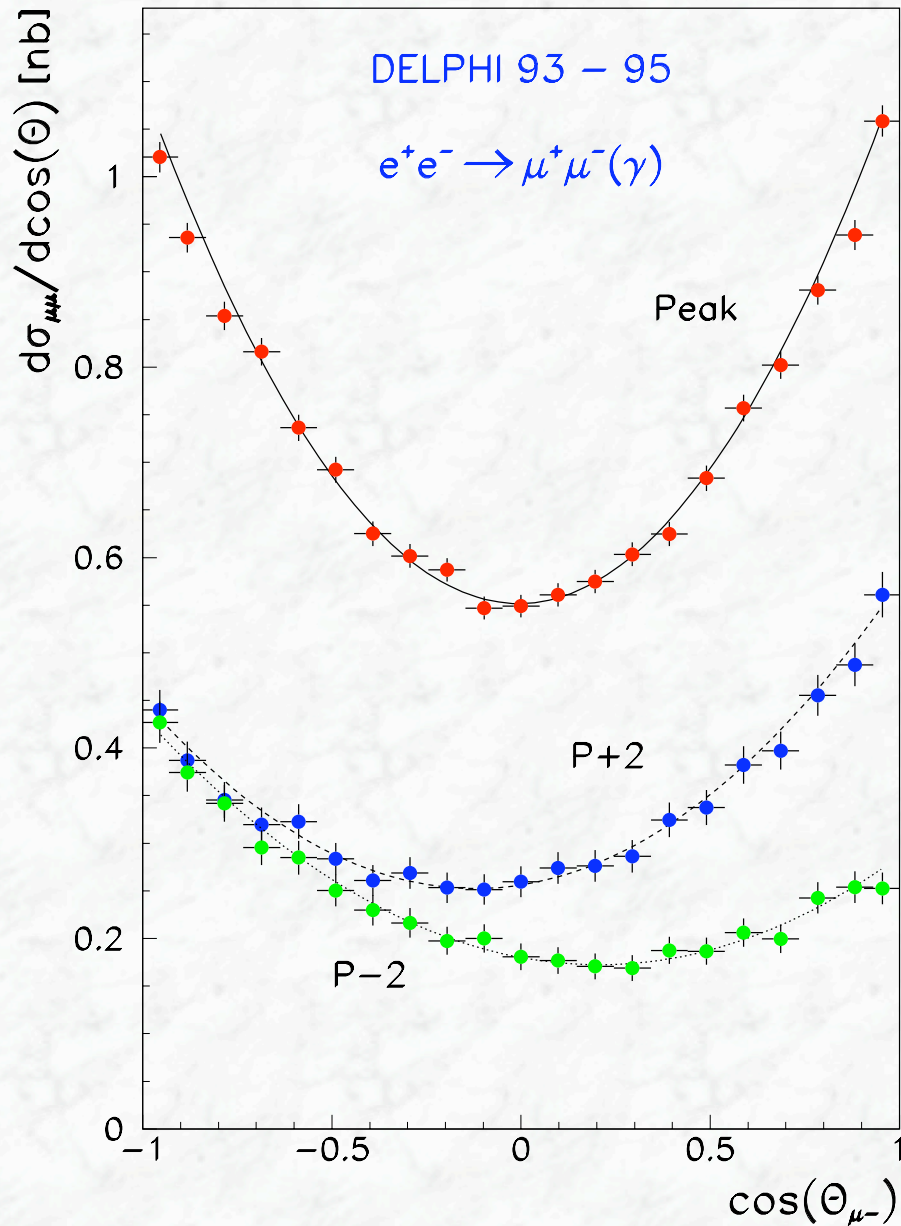
*) Uncertainty on LEP energy measurement: $\pm 1.7 \text{ MeV}$ (19 ppm)

Number of neutrinos



$$N_{\nu} = 2.9840 \pm 0.0082$$

Forward-backward asymmetries



$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2})(1 + \cos^2\theta) +$$

$$8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

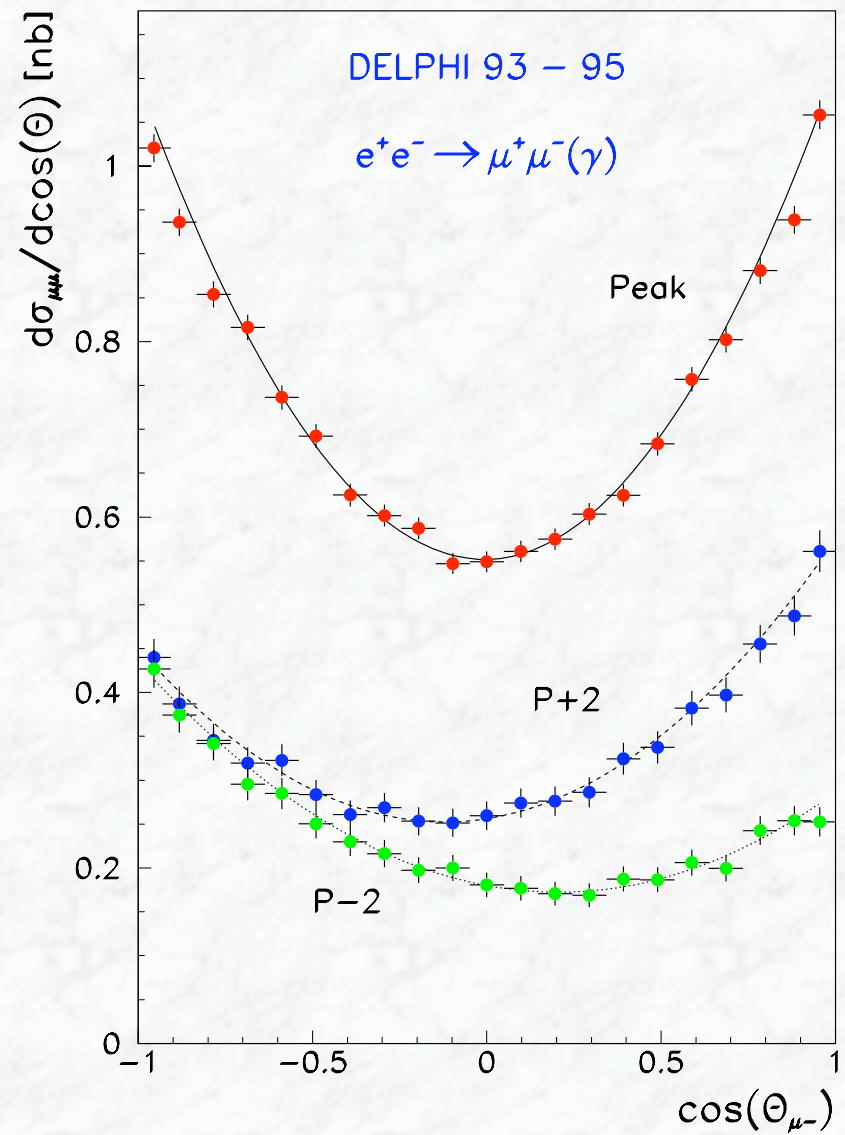
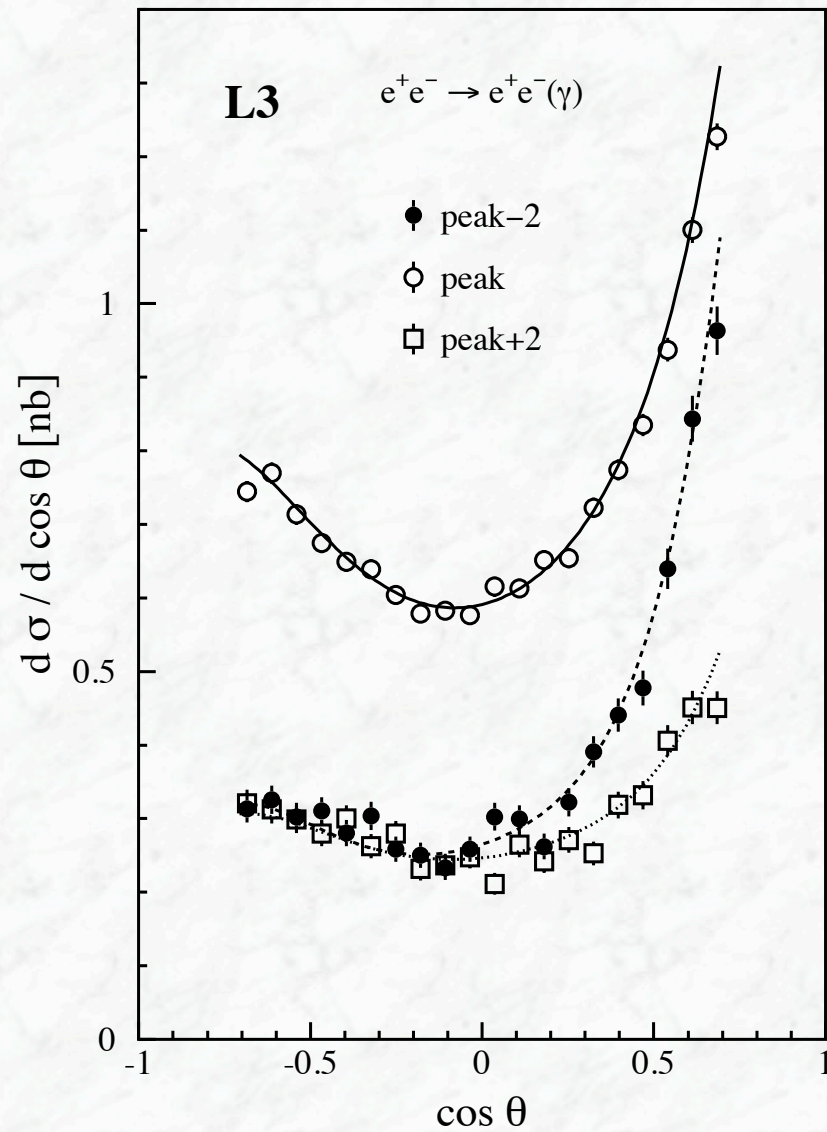
Terms $\propto \cos\theta$ in $d\sigma/d\cos\theta$
 \rightarrow asymmetry

$$\sigma_{F(B)} = \int_{0(-1)}^{1(0)} \frac{d\sigma}{d\cos\theta} d\cos\theta$$

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

Forward-backward asymmetries

-comparison between ee and $\mu\mu$ final states-



Forward-backward asymmetries and fermion couplings

- Asymmetry at the Z pole (no interference) **is small**

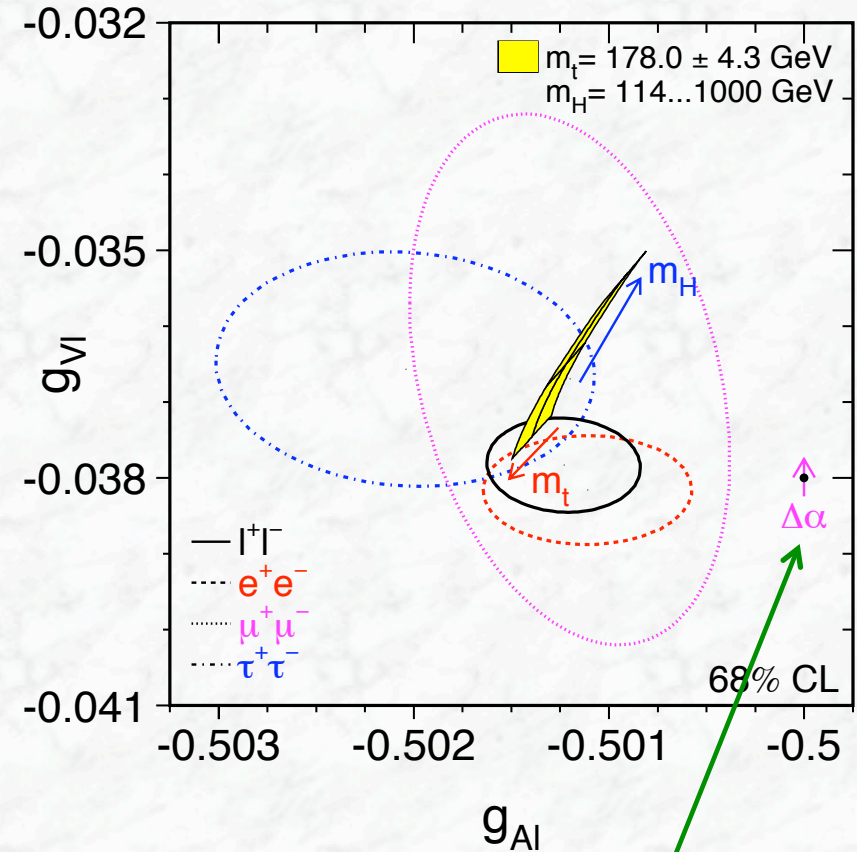
$$A_{\text{FB}} \sim g_A^e g_V^e g_A^f g_V^f$$

since g_V^f is small
(in particular for leptons)

- For off-resonance points, the interference term dominates and gives larger contributions

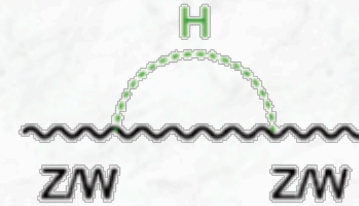
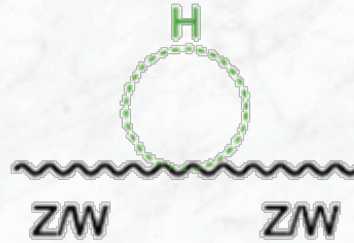
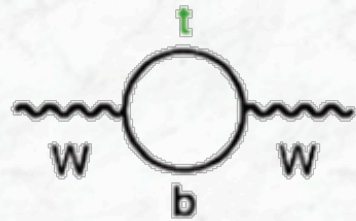
$$A_{\text{FB}} \sim g_A^e g_A^f \cdot \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

- A_{FB} can be used for the determination of the fermion couplings



LO Standard Model prediction:
 $g_A = T_3$
 $g_V = T_3 - 2 Q \sin^2 \theta_W$

Electroweak radiative corrections



Standard Model relations
(lowest order)

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F}$$

$$\alpha(0)$$

Relations including
radiative corrections

$$\vec{\rho} = 1 + \Delta\rho$$

$$\sin^2 \theta_{\text{eff}} = (1 + \Delta\kappa) \sin^2 \theta_W$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} \cdot \frac{1}{(1 - \Delta r)}$$

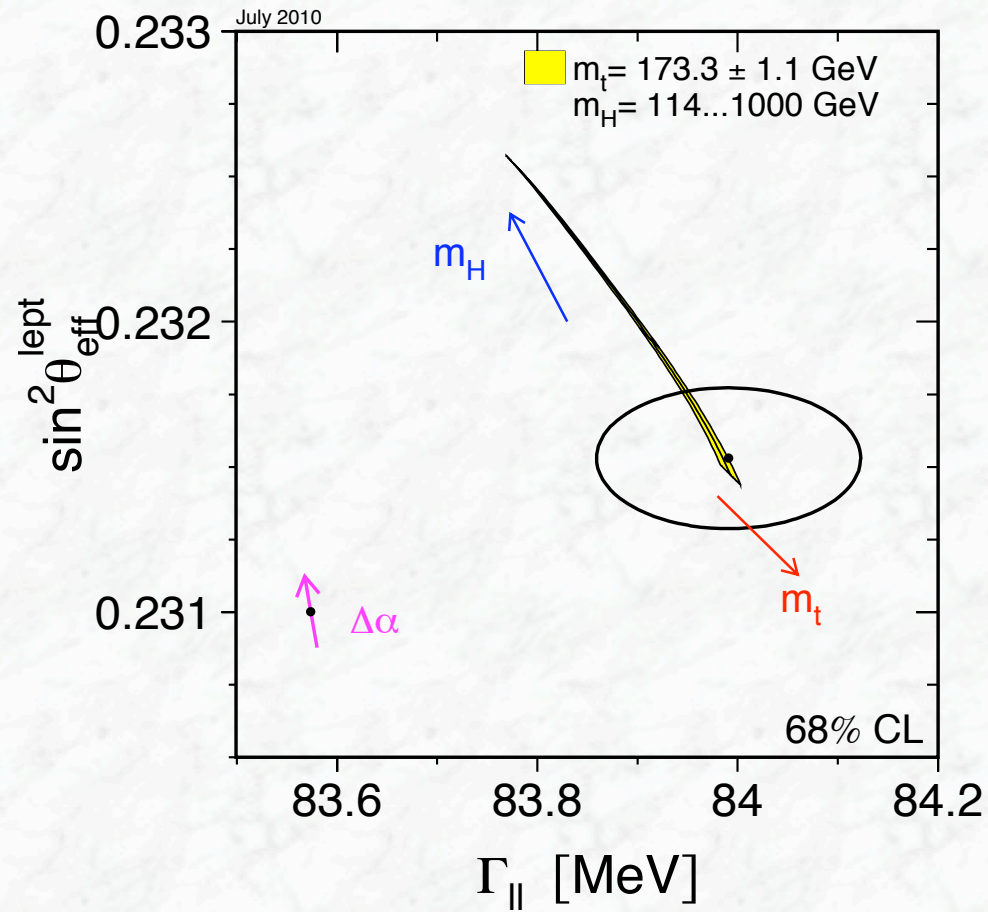
$$\alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta\alpha}$$

$$\Delta\alpha = \Delta\alpha_{\text{lepl}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$$

$$\Delta\rho, \Delta\kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$$

Results of electroweak precision tests at LEP (cont.)

partial decay width versus $\sin^2 \theta_W$:



Results of electroweak precision tests at LEP (cont.)

Summary of results:

- All measurements in agreement with the Standard Model
- They can be described with a limited set of parameters

