

## 3. Wechselwirkung von Photonen mit Materie

3.1 Absorption über Photoeffekt

3.2 Compton-Streuung

3.3 Paarerzeugung

3.4 Abschwächung, Absorptionskoeffizienten

3.5 Anwendungen

# Overview of interaction processes of electrons and photons

Energy loss due to excitation and ionisation

Bethe Bloch formula

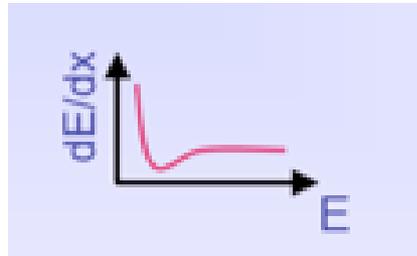
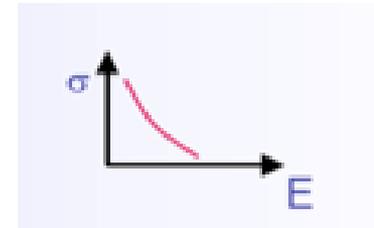
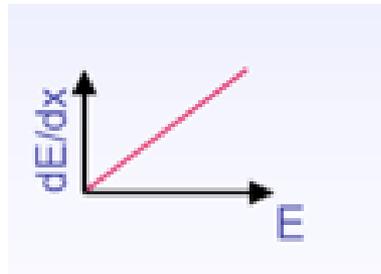


Photo effect

(dominant in  $\sim$  keV energy range)

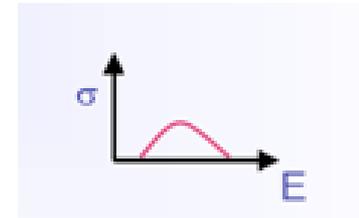


Bremsstrahlung



Compton effect

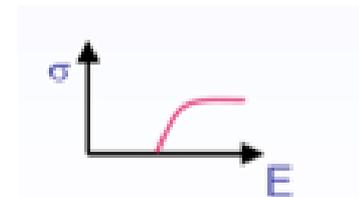
(dominant in MeV energy range)



Cherenkov radiation

Pair creation

(threshold energy =  $2 m_e = 1,022$  MeV)



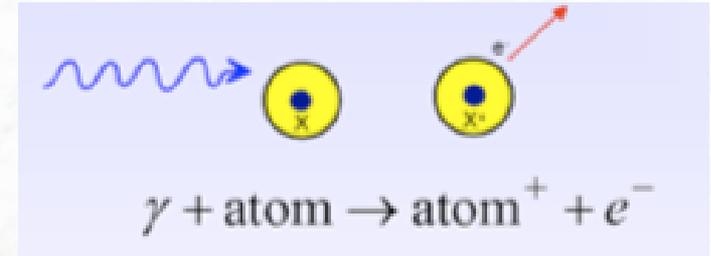
In order to be detected, photons must transfer their energy to charged particles

- Photo electric effect
- Compton scattering
- Pair creation

Photons “disappear” via these reactions. The intensity of a photon beam is exponentially attenuated in matter:

$$I(x) = I_0 e^{-\mu x}$$

## 3.1 Photo electric effect:



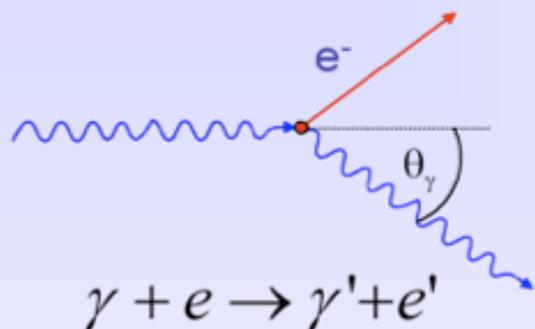
- Release of electrons from the inner shells (K, L, ..) of atoms  
(Only possible in the close neighbourhood of a third collision partner)
- The cross section shows a strong modulation if  $E_\gamma \approx E_{\text{bin}}$  (binding energy)

$$\sigma_{photo}^K = \left(\frac{32}{\epsilon^7}\right)^{\frac{1}{2}} \alpha^4 Z^5 \sigma_{Th}^e \quad \epsilon = \frac{E_\gamma}{m_e c^2} \quad \sigma_{Th}^e = \frac{8}{3} \pi r_e^2 \quad (\text{Thomson})$$

At high energies ( $\epsilon \gg 1$ )

$$\sigma_{photo}^K = 4\pi r_e^2 \alpha^4 Z^5 \frac{1}{\epsilon} \quad \boxed{\sigma_{photo} \propto Z^5}$$

## 3.2 Compton scattering



$$E'_\gamma = E_\gamma \frac{1}{1 + \varepsilon(1 - \cos\theta_\gamma)}$$

$$E_e = E_\gamma - E'_\gamma$$

Assume electron as quasi-free.

**Klein-Nishina**  $\frac{d\sigma}{d\Omega}(\theta, \varepsilon)$   $\rightarrow$

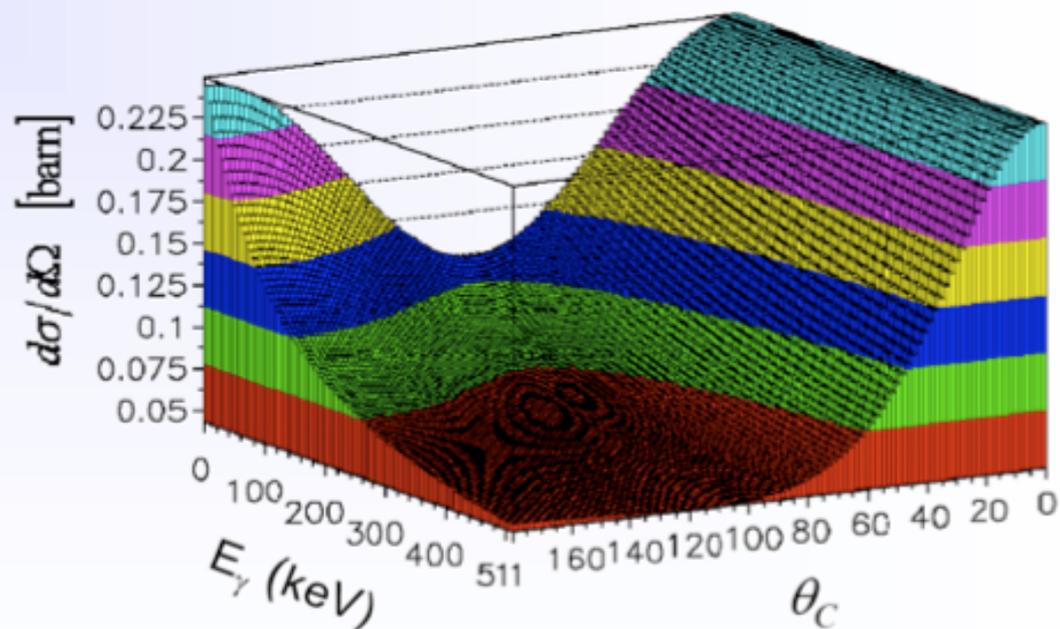
At high energies approximately

$$\sigma_c^e \propto \frac{\ln \varepsilon}{\varepsilon}$$

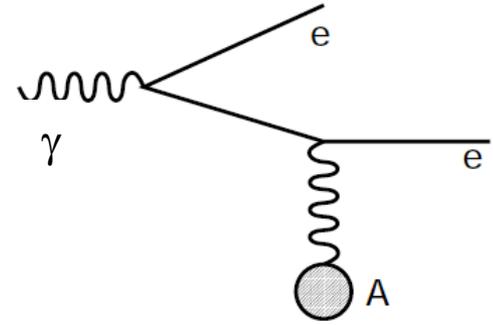
Atomic Compton cross-section:

$$\sigma_c^{atomic} = Z \cdot \sigma_c^e$$

Compton cross-section (Klein-Nishina)



### 3.3 Pair production: $\gamma + (A) \rightarrow e^+ e^- + (A)$



- Only possible in the close neighbourhood of a collision partner (atomic nucleus)
- Threshold energy:  $E_\gamma > 2 m_e c^2 = 1.022 \text{ MeV}$
- Cross section (high energy approximation):

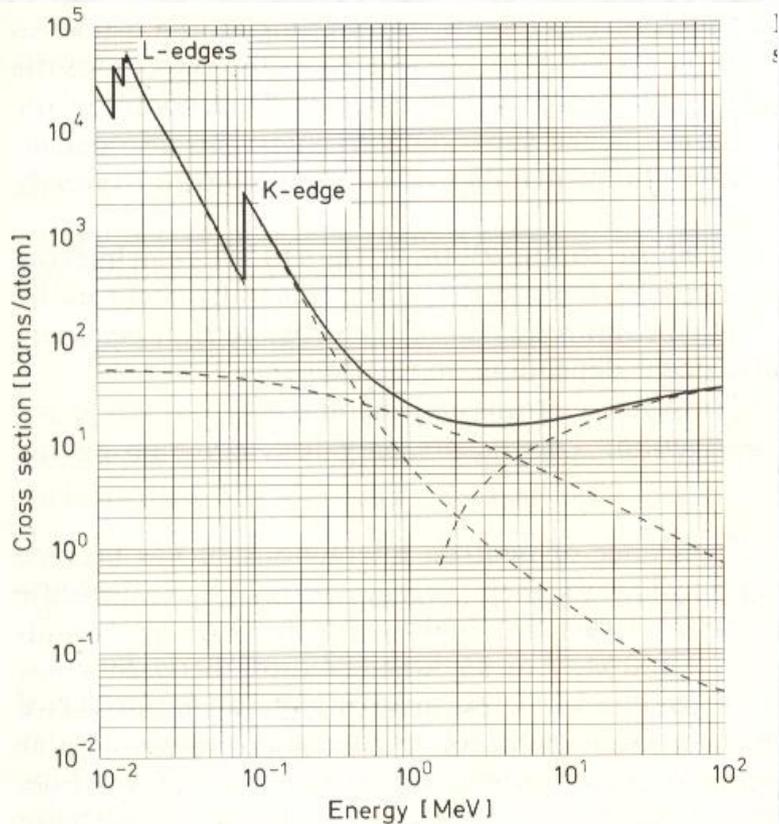
$$\sigma_{pair} = 4\alpha r_e^2 Z^2 \left( \frac{7}{9} \ln \frac{183}{Z^{1/3}} - \frac{1}{54} \right) \approx \frac{7}{9} \frac{A}{N_A} \frac{1}{X_0}$$

$$\Rightarrow \mu_{pair} = \frac{7}{9} \frac{1}{X_0}$$

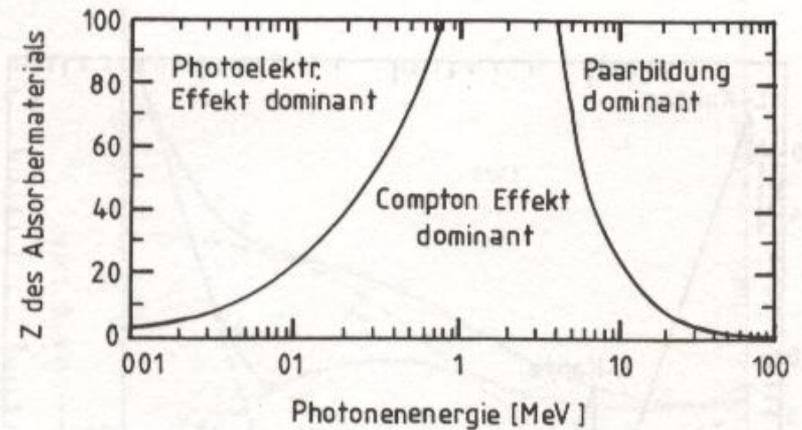
- After traversing a material thickness of  $9/7 X_0$ , the photon intensity –due to pair creation- is decreased by  $1/e$ .  
.or.
- For high photon energies, pair production occurs after traversing a material thickness corresponding to one radiation length with a probability of

$$p = 1 - e^{-7/9} = 0.54$$

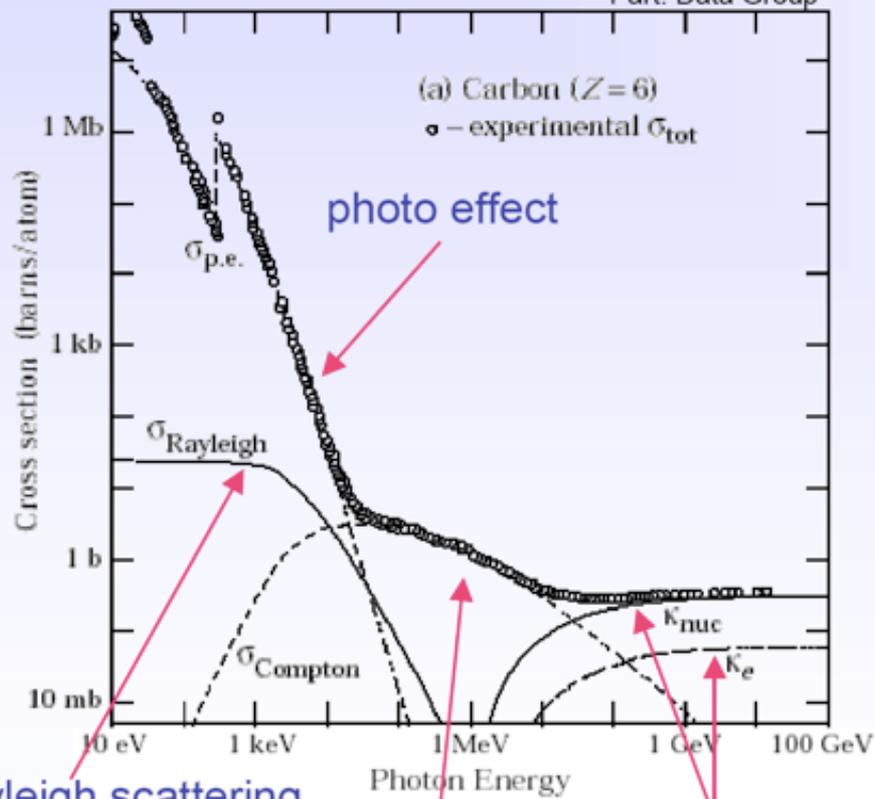
## 3.4 Photon interaction cross sections



Total photon absorption cross section for lead ( $Z=82$ ) as a function of the photon energy. The contributions from photo absorption, Compton scattering and pair creation are shown separately [Ref. Leo]



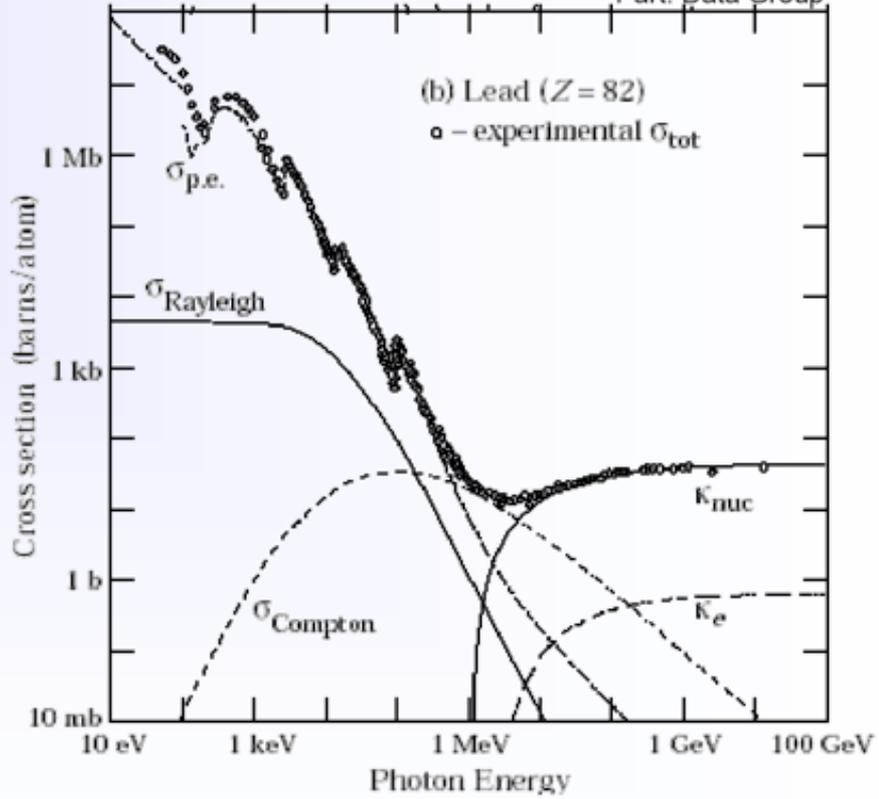
Regions in which photo effect, Compton scattering, and pair creation dominate, as a function of the photon energy and the atomic charge ( $Z$ ) of the absorber [Ref. Grupen]



Rayleigh scattering  
(no energy loss !)

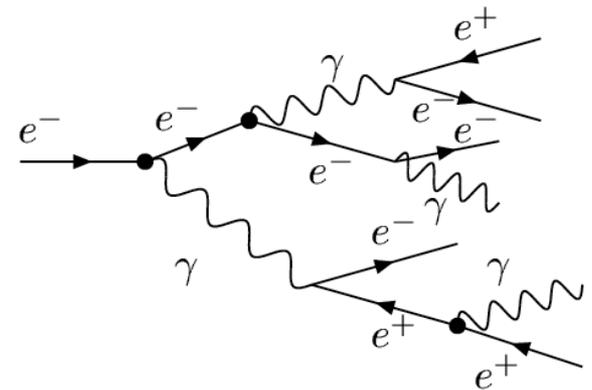
Compton scattering

pair production



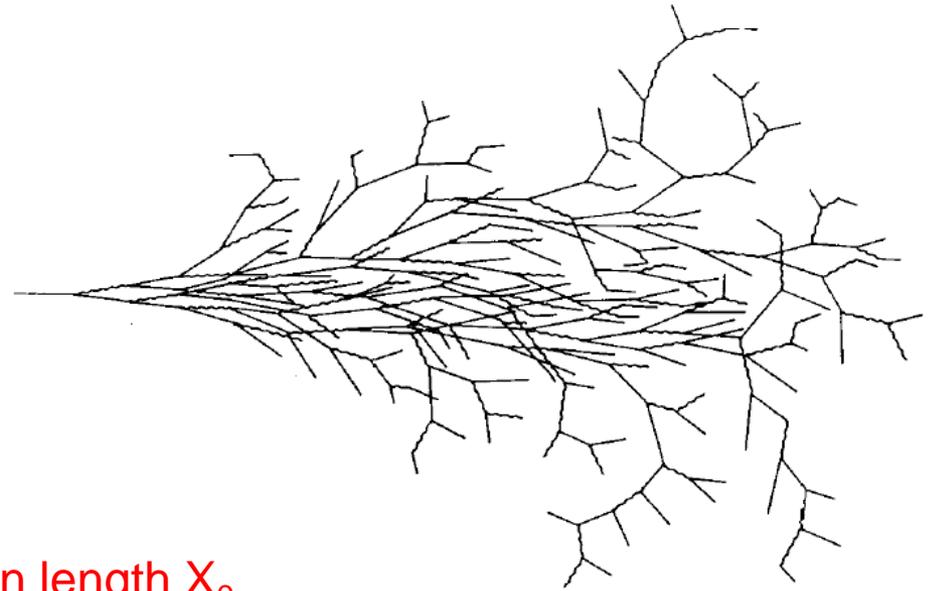
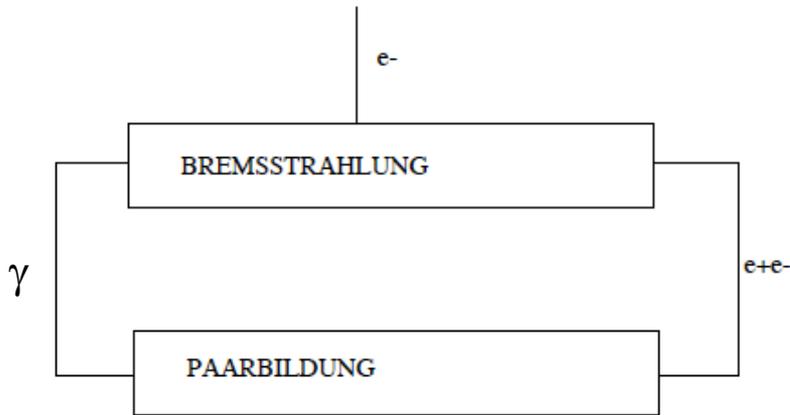
## 3.5 Applications

- Use photon attenuation for investigations of material structure  
(X-ray imaging, ...)
- Tumor therapy  
(with drawback of non-localized energy deposition)
- Energy measurement of electrons and photons in so-called electromagnetic calorimeters  
(Bremsstrahlung, pair-creation processes  
→ formation of so-called electromagnetic showers)



# Concept of electromagnetic calorimeters

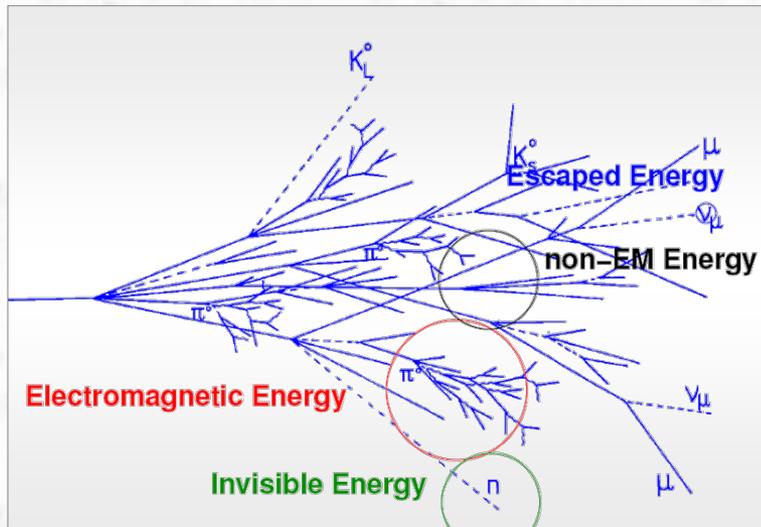
- Particle showers created by electrons/positrons or photons are called electromagnetic showers (only electromagnetic interaction involved)
- Basic processes for particle creation: bremsstrahlung and pair creation



- Characteristic interaction length: radiation length  $X_0$
- Number of particles in the shower increases, until the critical energy  $E_c$  is reached; For  $E < E_c$  the energy loss due to ionization and excitation dominates, the number of particles decreases, due to stopping in material

# Hadronic showers / calorimeters

- Hadrons initiate their energy shower by inelastic hadronic interactions; (strong interaction responsible, showers are called **hadronic showers**)
- Hadronic showers are much more complex than electromagnetic showers



- Several secondary particles, meson production, multiplicity  $\sim \ln(E)$
- $\pi^0$  components,  $\pi^0 \rightarrow \gamma\gamma$ , **electromagnetic sub-showers**;  
The fraction of the electromagnetic component grows with energy,  
 $f_{EM} = 0.1 \ln E$  (E in GeV, in the range  $10 \text{ GeV} < E < 100 \text{ GeV}$ )

- During the hadronic interactions atomic nuclei are broken up or remain in excited states

Corresponding energy (excitation energy, binding energy) comes from original particle energy

→ no or only partial contribution to the visible energy

- In addition, there is an important **neutron component**

The interaction of neutrons depends strongly on their energy;

Extreme cases:

- Nuclear reaction, e.g. nuclear fission → energy recovered
- Escaping the calorimeter (undergo only elastic scattering, without inelastic interaction)

- Decays of particles (slow particles at the end of the shower)

e.g.  $\pi \rightarrow \mu \nu_\mu$  → escaping particles → missing energy

These energy loss processes have important consequences:  
in general, the response of the calorimeter to electrons/photons and hadrons is different ! The signal for hadrons is non-linear and smaller than the e/ $\gamma$  signal for the same particle energy

# Layers of the ATLAS detector

